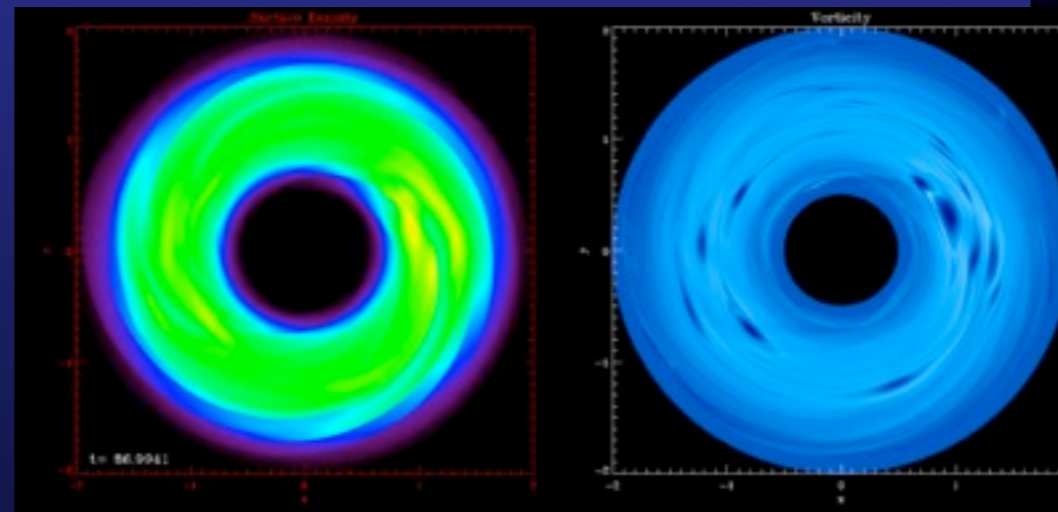
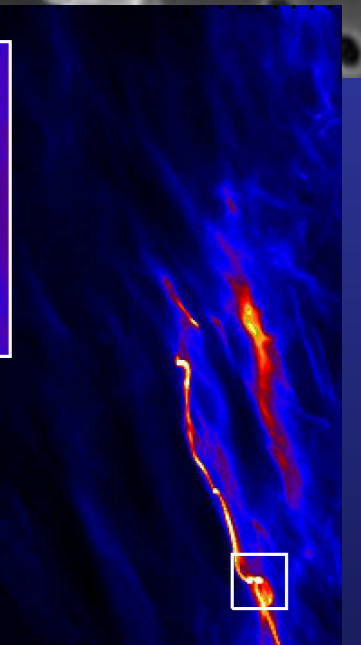
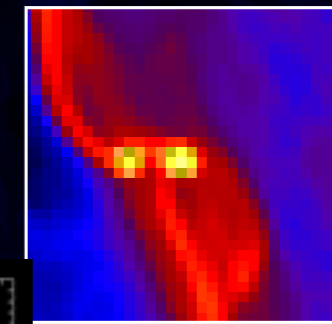
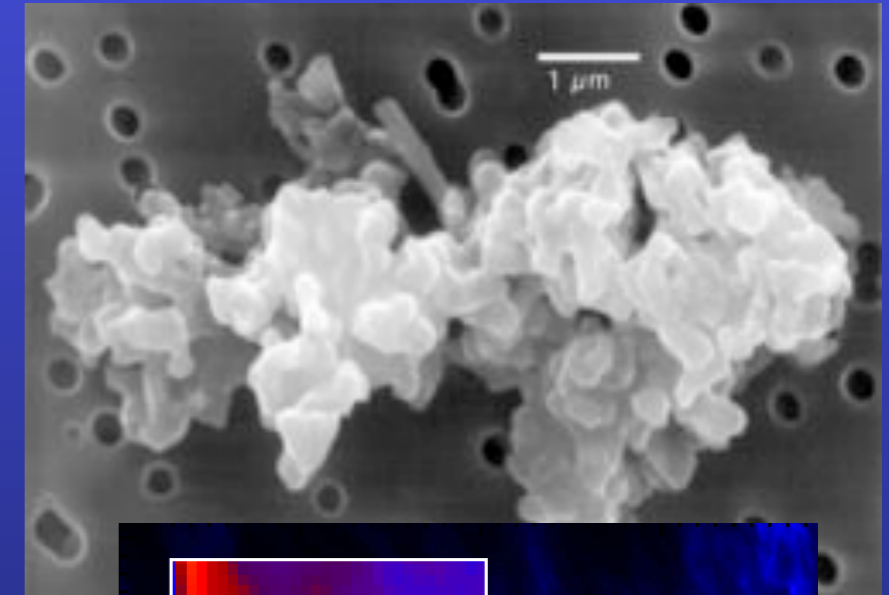
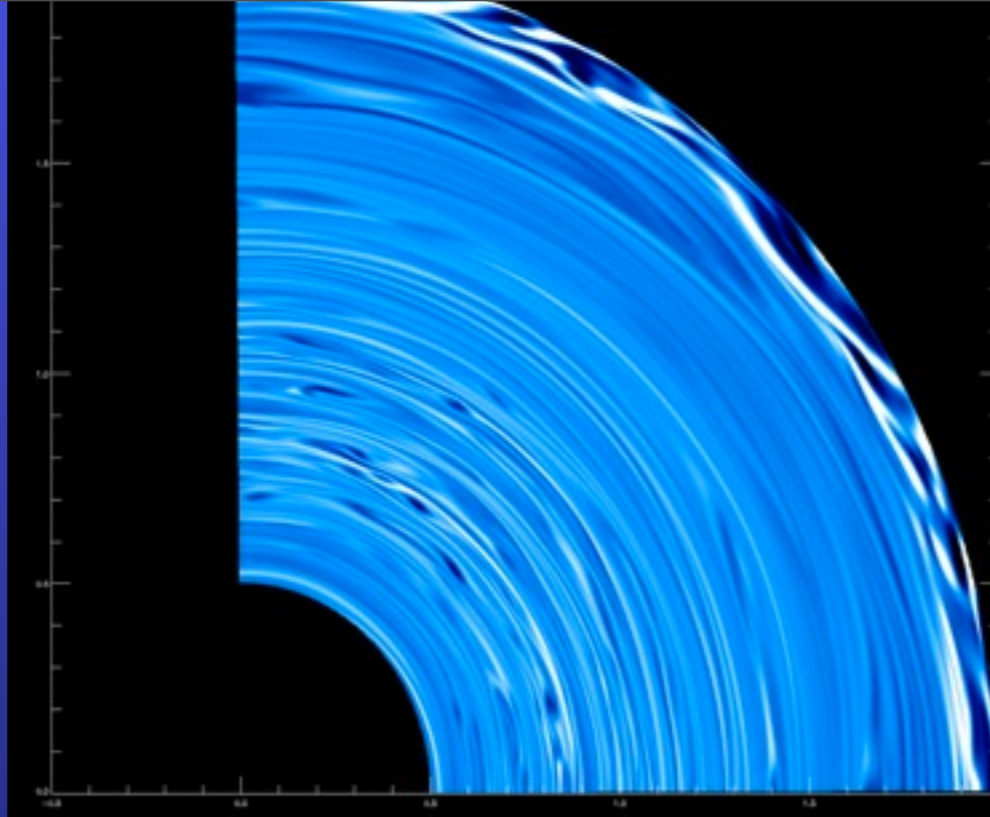




Marseille, Sep. 18<sup>th</sup>, 2012



# Disk Weather: Baroclinic Instability and Vortex Amplification

**Hubert Klahr,**  
Max-Planck-Institut für Astronomie, Heidelberg

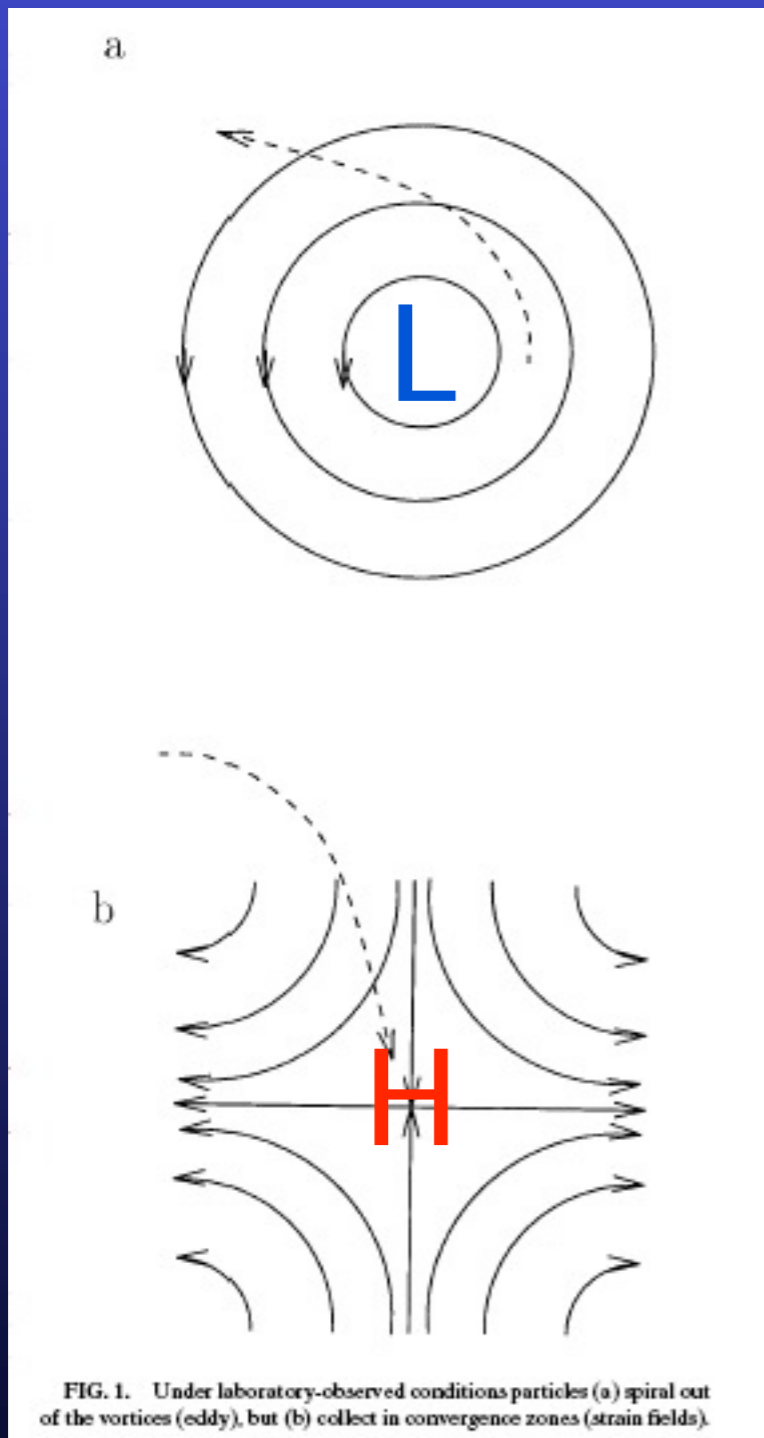
Natalie Raettig, Alex Hubbard, Moritz Beutel, Karsten Dittrich (MPIA), Wlad Lyra (JPL), Peter Bodenheimer (UCSC), Helen Morrison (ITA)

Planetesimal Formation: Rainer Spurzem (NAOC/ARI), Mario Trieloff (HD), Til Birnstiel, Barbara Ercolano (USM), Kees Dullemond (ITA), Chris Ormel (Berkeley), Neal Turner (JPL), Doug Lin (KIAA, UCSC) Anders Johansen (Lund)

# Outline: **Disk Weather: Baroclinic Instability and Vortex Amplification - From physical conditions in disks to 3D global radiation hydro simulations**

- Current state of MHD turbulence and Gravoturbulent Planetesimal formation
- Turbulence and vortices in Baroclinic disks: Disk Weather
- Linear Axisymmetric Instability
- Dust capturing in 3D vortices
- Summary, Conclusions

# Turbulence in a non rotating frame:



Laboratory Conditions  
Dust collects between  
vortices (high pressure)

$$\partial_t v_g = -\frac{1}{\rho} \nabla p + \text{forces}$$

$$\partial_t v_d = -\frac{v_d - v_g}{\tau_f} + \text{forces}$$

$$v_d = v_g + \tau_f \frac{1}{\rho} \nabla p$$

Cuzzi et al.



From dust/gas eq. 1 => Streaming Instability -  
> dust densities larger than local Roche  
density, e.g. Planetesimals can form.

v  
e  
r  
t  
i  
c  
a  
l

h  
o  
r  
i  
z  
o  
n  
t  
a  
l

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Johansen, Henning & Klahr 2006

12/13/2009

From dust/gas eq. 1 => Streaming Instability -  
> dust densities larger than local Roche  
density, e.g. Planetesimals can form.

V  
e  
r  
t  
i  
c  
a  
l

$t = 0.1$

h  
o  
r  
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z  
o  
n  
t  
a  
l

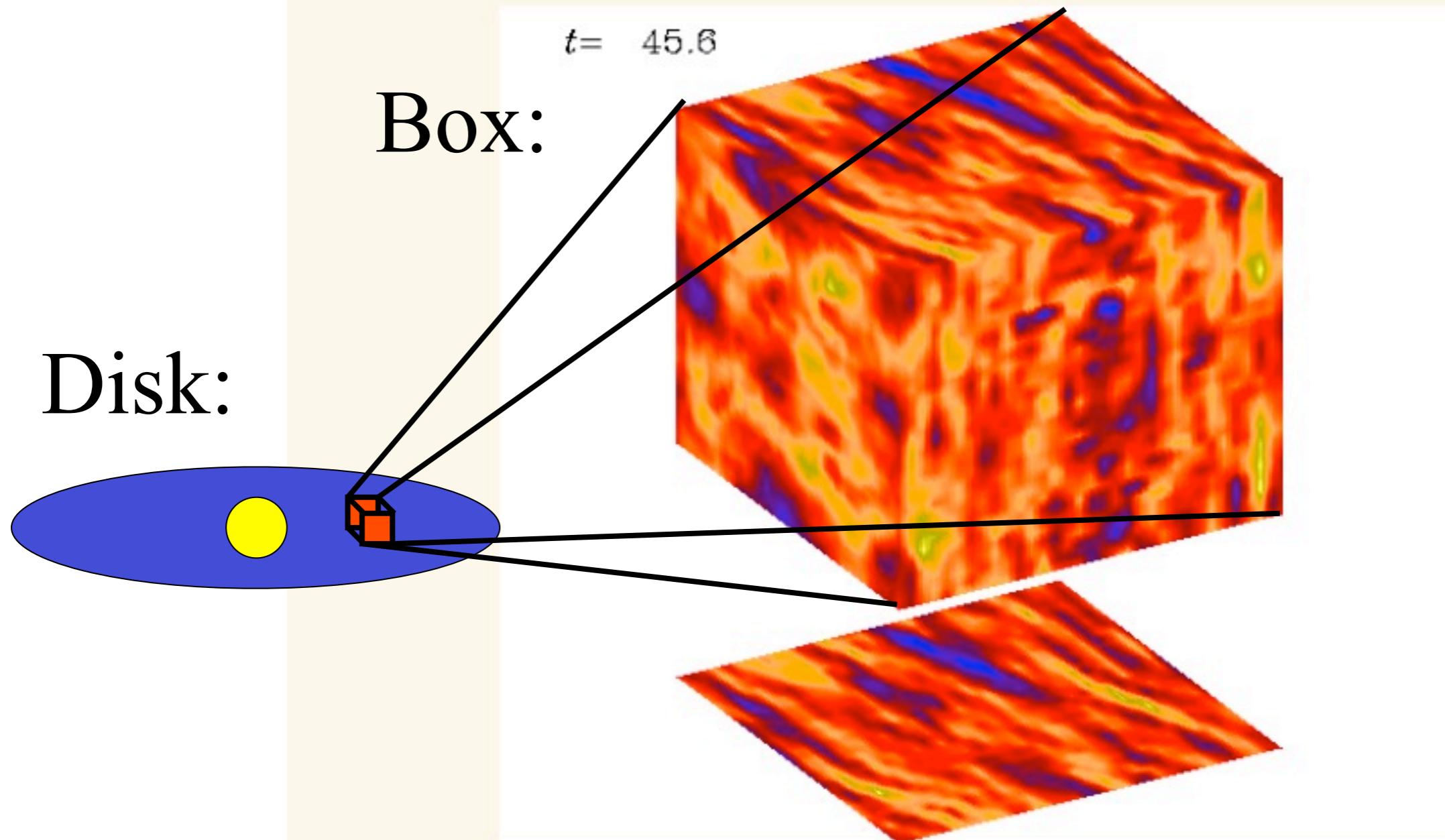
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Johansen, Henning & Klahr 2006

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# MRI turbulence

...because it is a reliable source for turbulence.

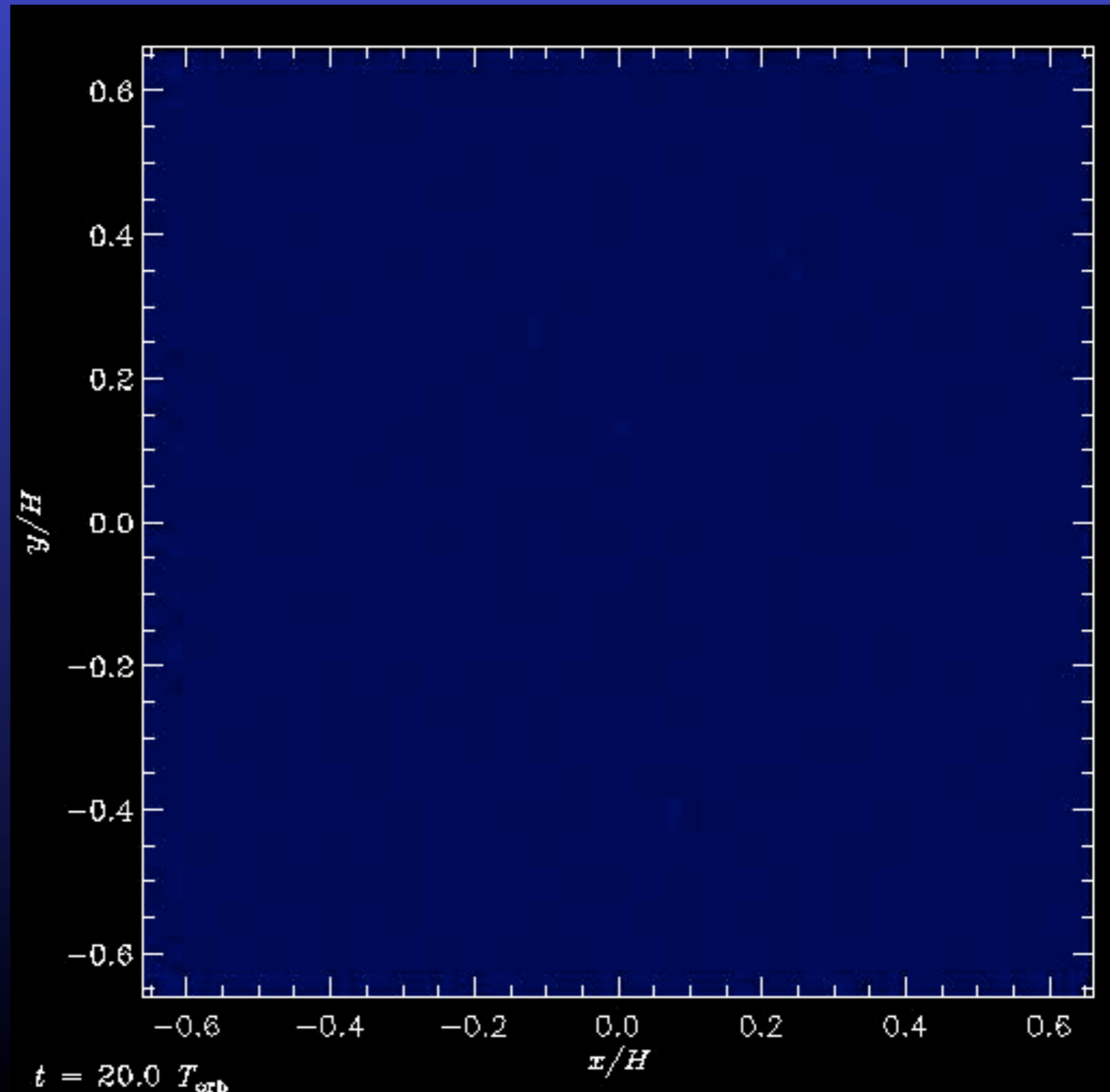


Code: The Pencil-Code [MHD code, finite differences, 6th order in space, 3rd order in time, Brandenburg (2003)]

# Concentration in Zonal Flows:

Formation Of  
Planetesimals  
From pressure  
trapped /  
gravitational  
Bound heaps of  
gravel - here  
magnetic  
turbulence:  
Johansen, Klahr &  
Henning 2011.  
Next: Raettig &  
Lyra

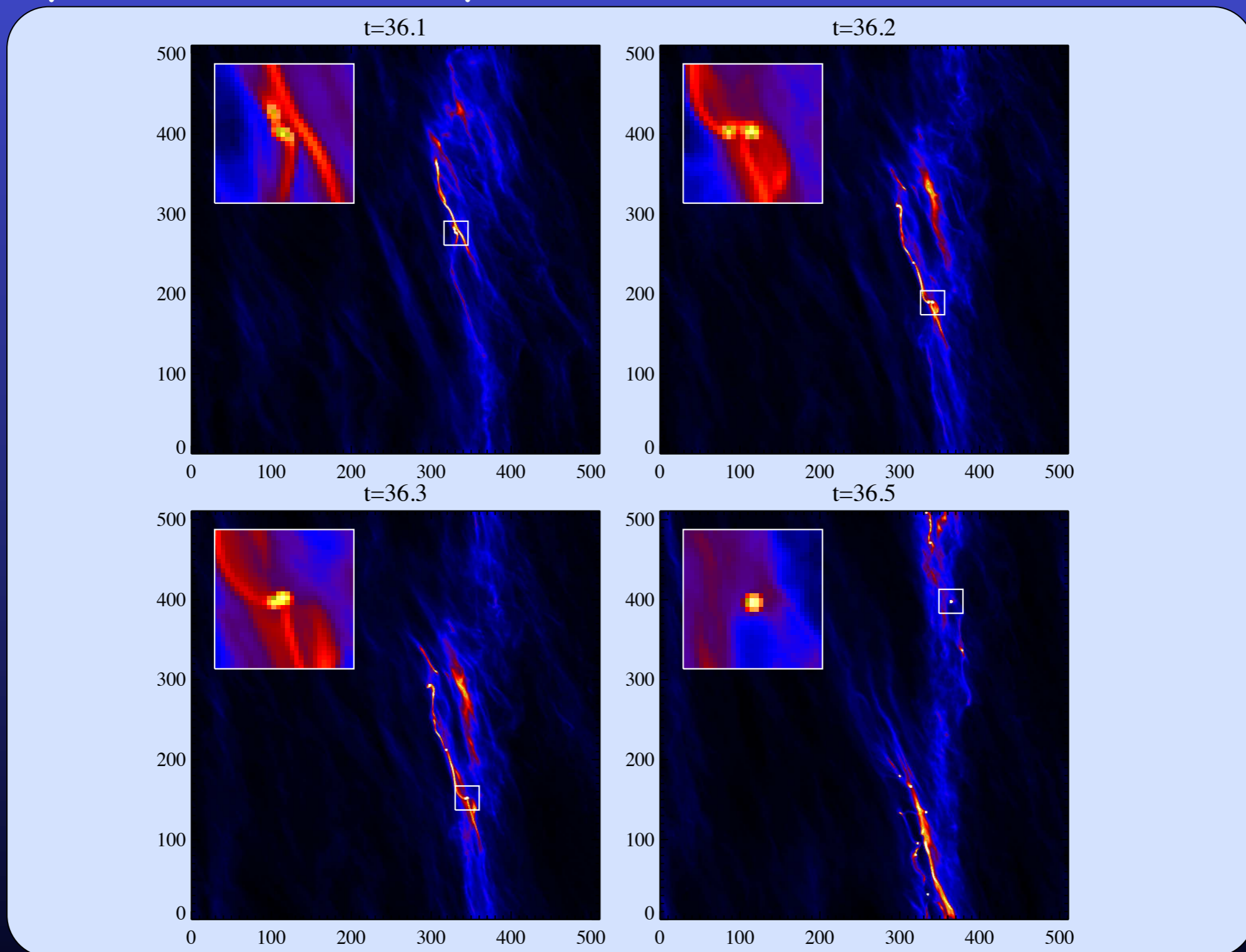
512  $\Lambda^2$  simulation  
64 Mio particles  
Entire project  
used 15 Mio. CPU  
hours.



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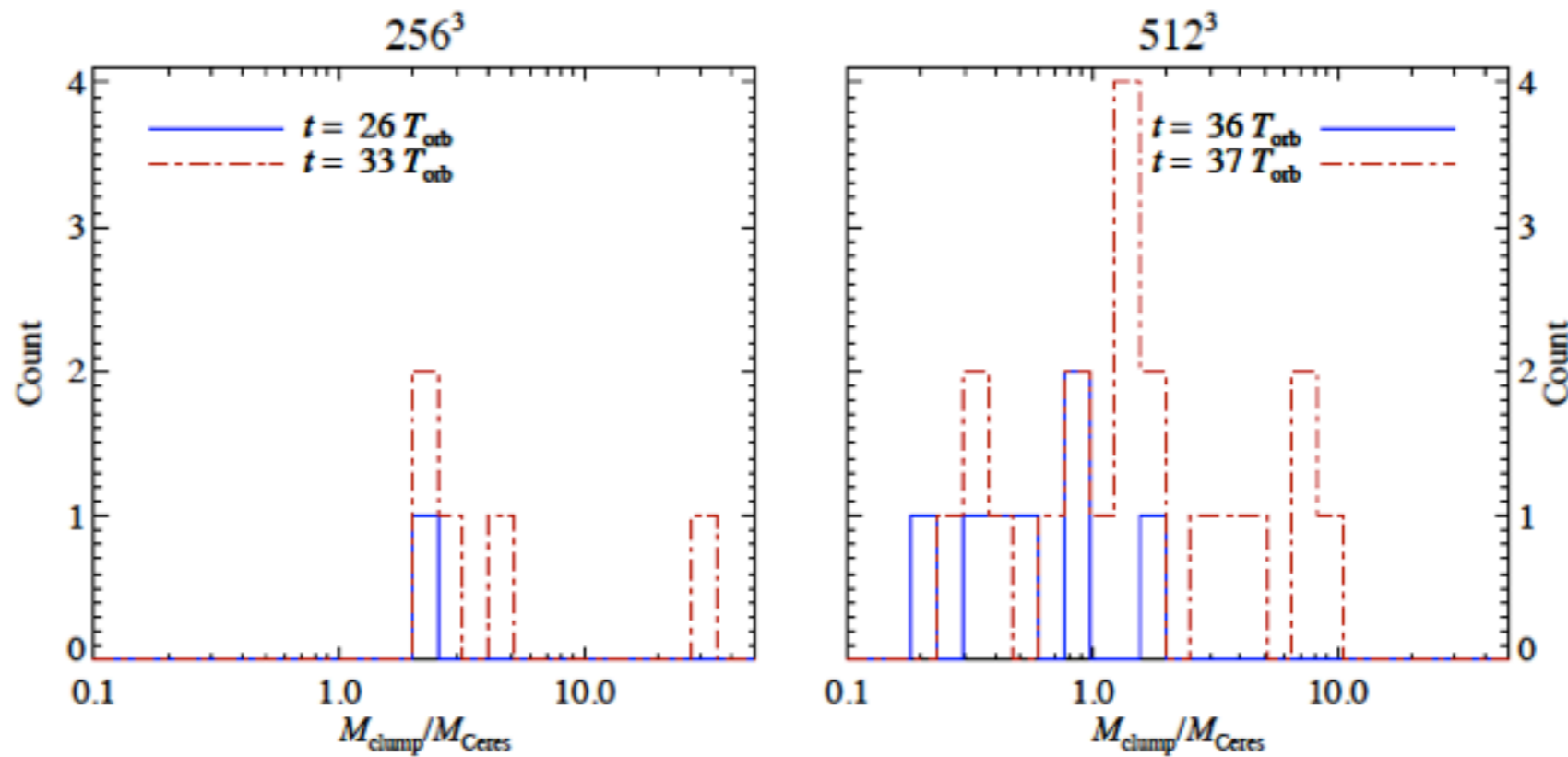


At  $512^3$  and the usual set up: As before mass of planetesimals depends on the available mass.



new: Intermediate Formation of Binaries;  
Johansen, Klahr & Henning (2011)

Not the process but numerical resolution limit the smallest possible planetesimals. Johansen, Klahr & Henning 2011



**Fig. 10.** Histogram of clump masses after first production of bound clumps and at the end of the simulation. At moderate resolution (left panel) only a single clump condenses out initially, but seven orbits later there are five clumps, including the  $30+ M_{\text{Ceres}}$  object formed by merging. At high resolution (right panel) the initial planetesimal burst leads to the formation of many sub-Ceres-mass clumps. The most massive clump is similar to what forms initially in the moderate-resolution run.

To predict an initial mass distribution of Planetesimals one needs A: the proper size distribution of precursors and B: the likely hood for concentrations as a function of particle size.

# Larger Boxes = better concentration Dittrich, Klahr & Johansen submitted

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Dittrich, Klahr, & Johansen

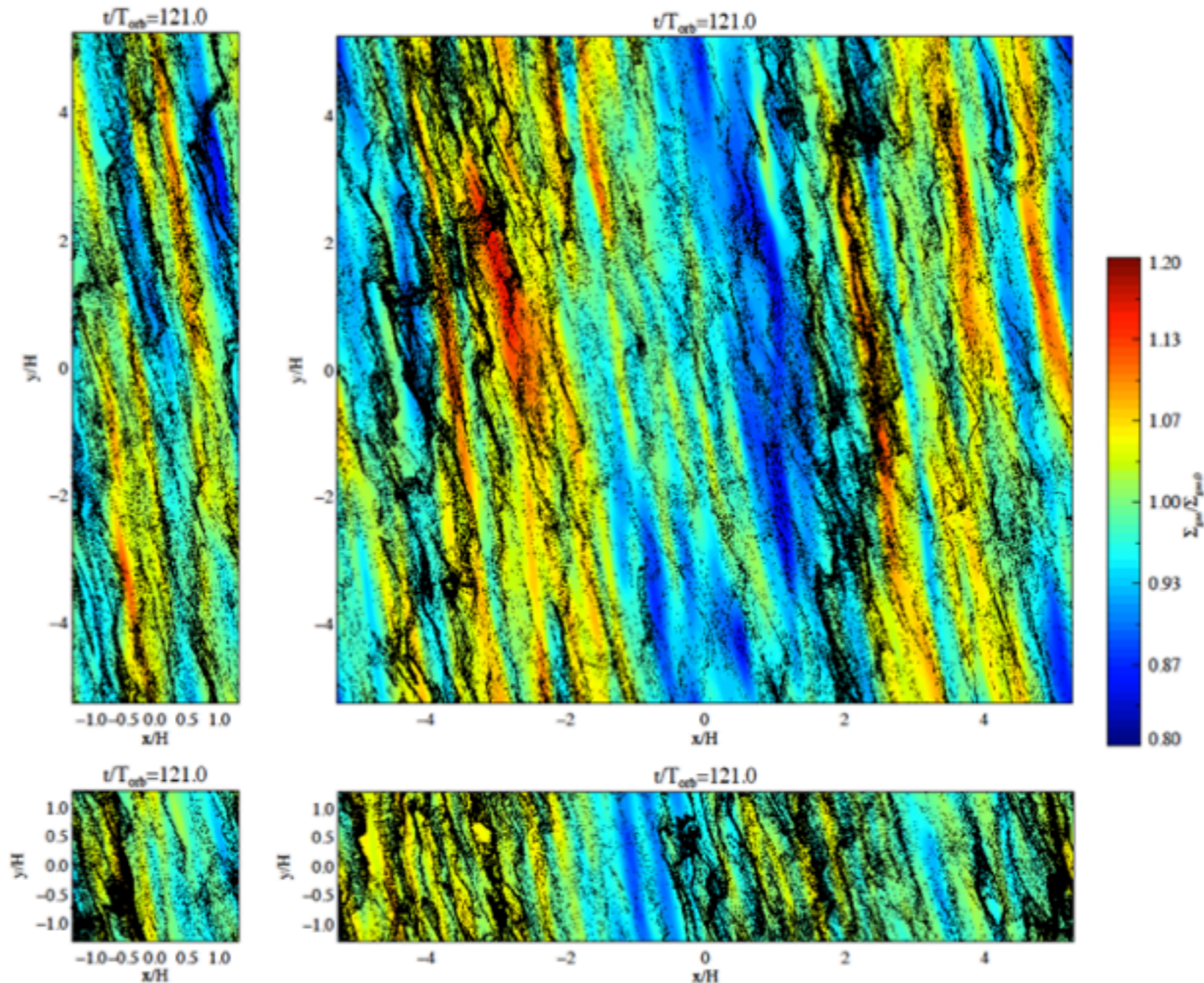
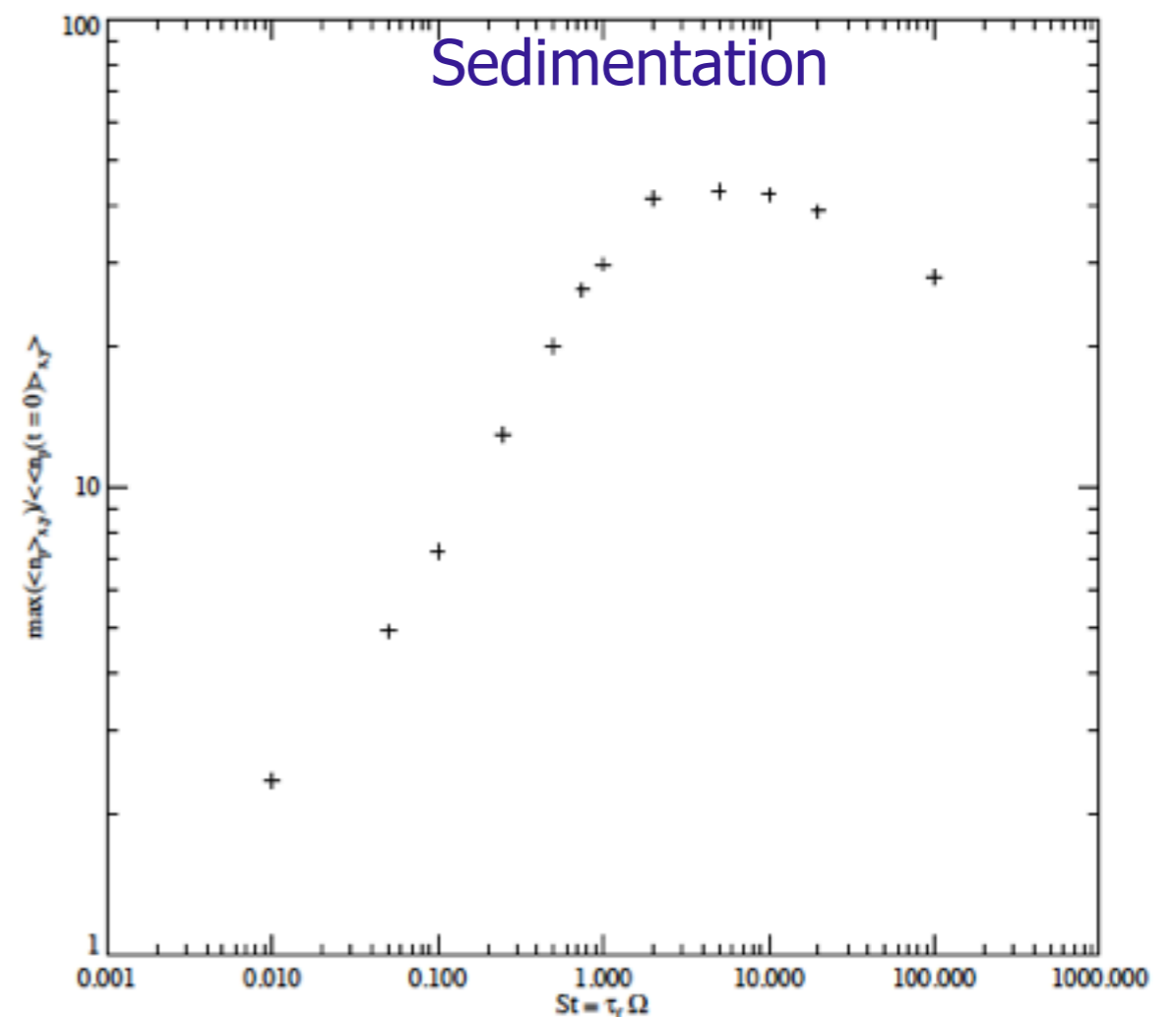
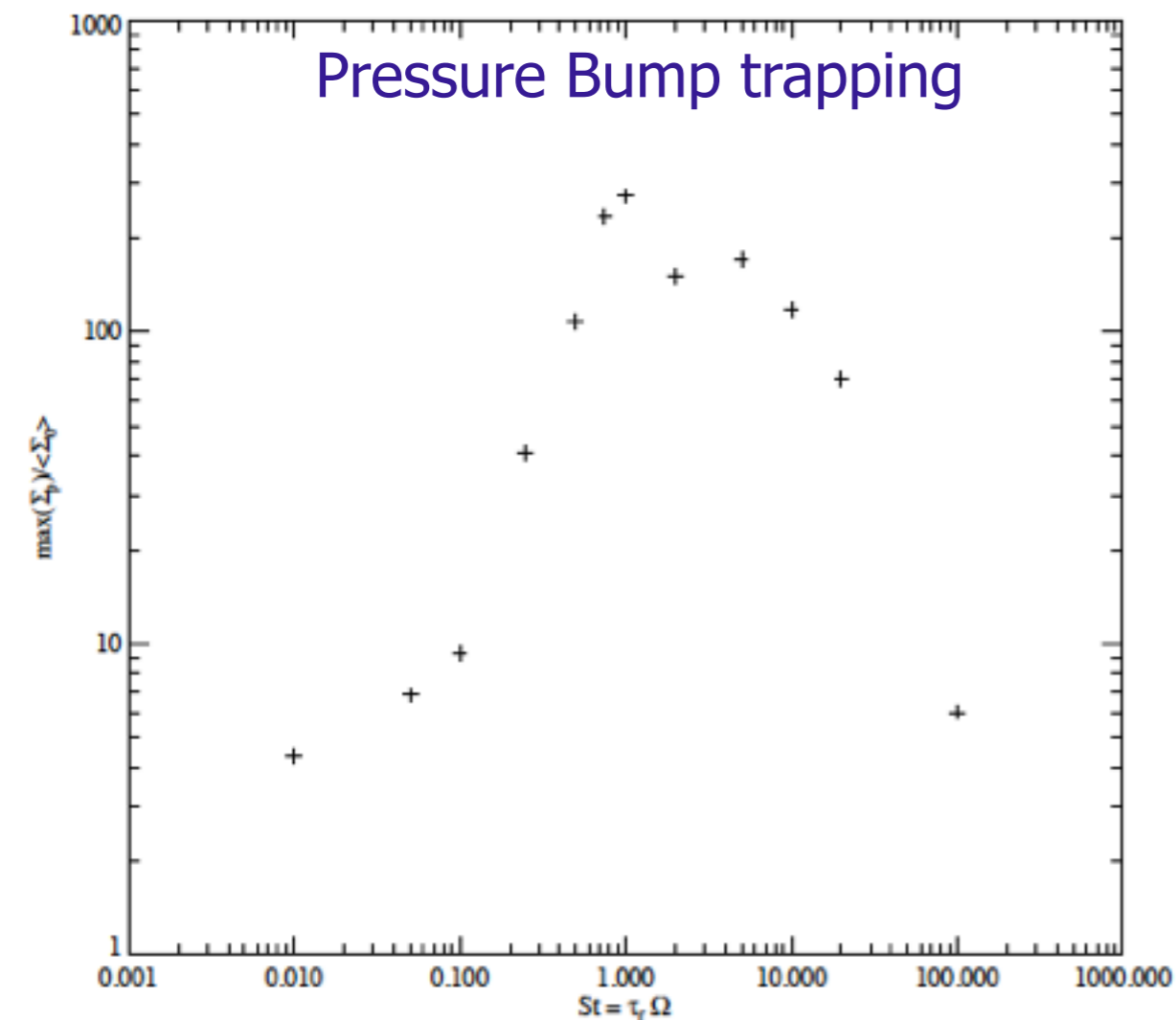


Figure 4. A collage of 4 surface gas density plots of the simulations  $y$ -XL,  $x$ L,  $M$ , and  $x$ -XL (from the upper left to the lower right panel respectively), each at the end of the simulation. The color bar at the right-hand side is valid for all 4 panels. This collage shows that over-densities are most pronounced in the biggest box. The azimuthally extended box shows regions of similar high density as simulation  $x$ L. However, the structure seems to be thinner for simulation  $y$ -XL. The black dots represent the position of every 20<sup>th</sup> particle, integrated in vertical direction.

Larger Boxes = better concentration  
 Even for much smaller particles!  
 Even  $St = 0.1$  can ignite streaming inst.  
 and grav.-turb. planetesimal formation

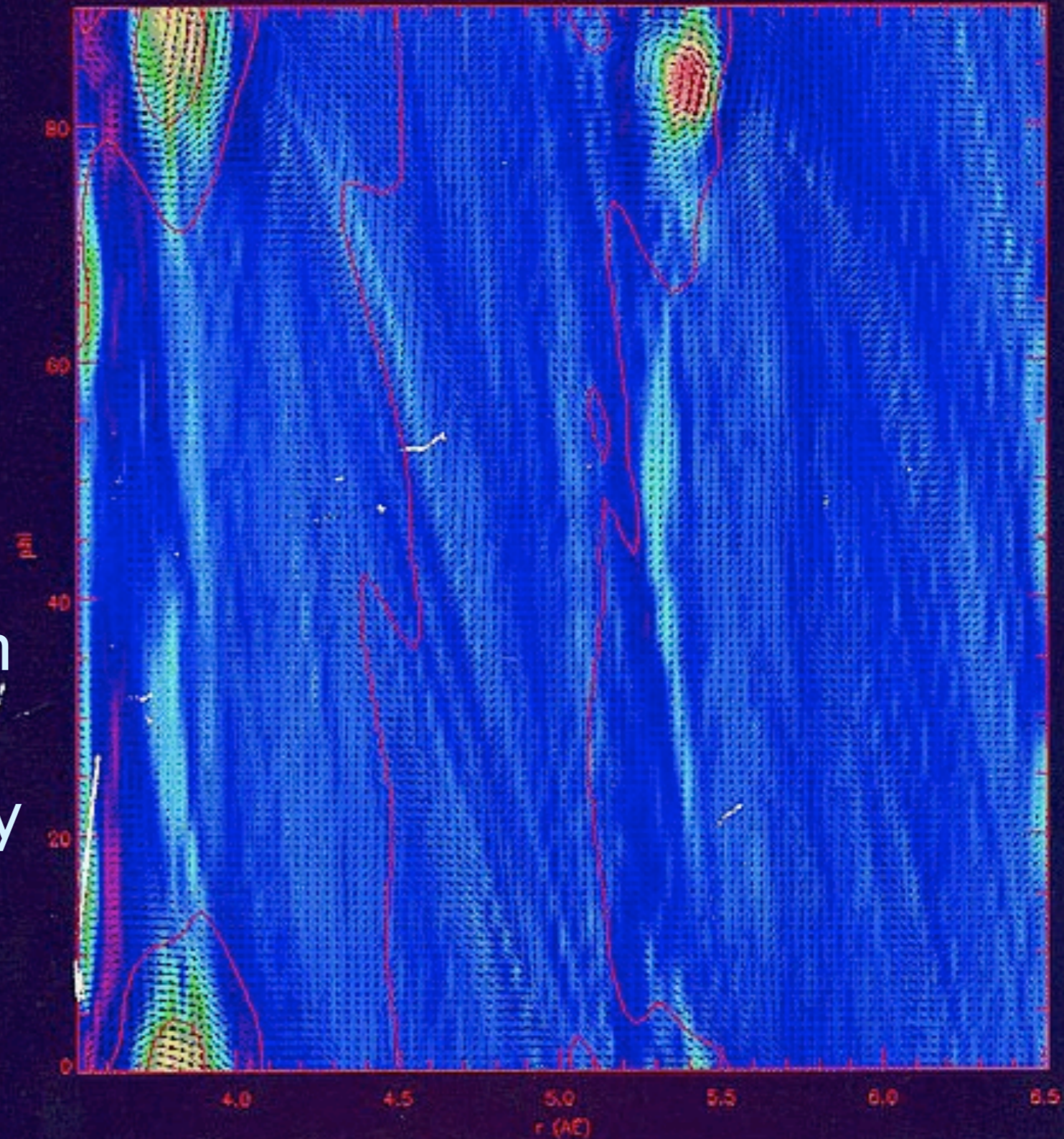


14 years back...  
...before I was doing this  
with MHD we tried Thermal  
Convection because my  
advisors said so... ;)

My very first vortices:  
The historic Simulation:  
About 1999  
Large Scale  
3D - Simulation  
90 degree  
3.5 - 6.5 AU  
102 X 40 X 120 cells  
=> Vortices

3D Global Disk Simulation  
flux limited Diffusion  
temperatur maintained by  
artificial heating

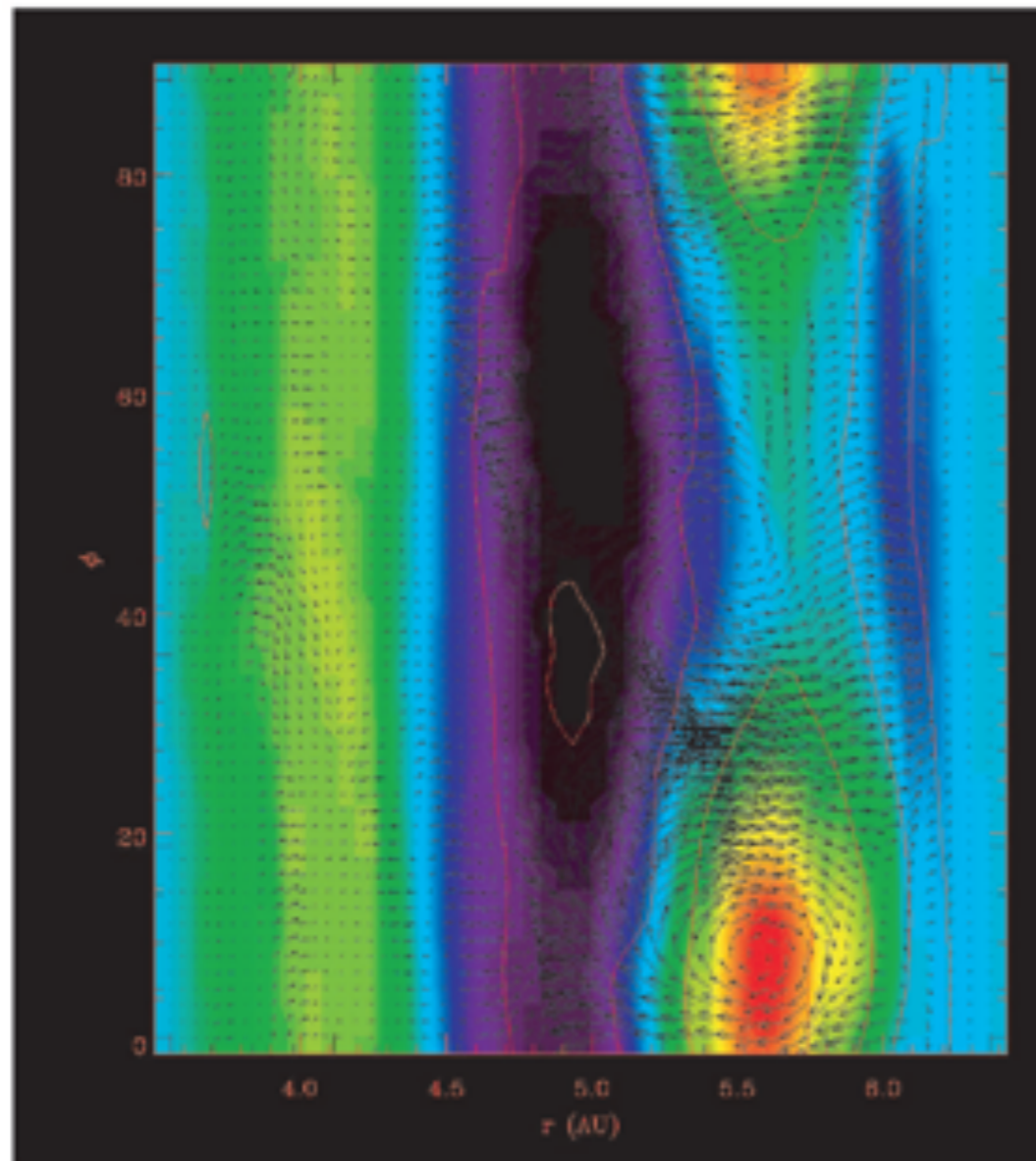
VORTICITY  $\Rightarrow$  ANTI-  
CYCLOVES  
&  
ROSSBY WAVES  
102 x 40 x 120



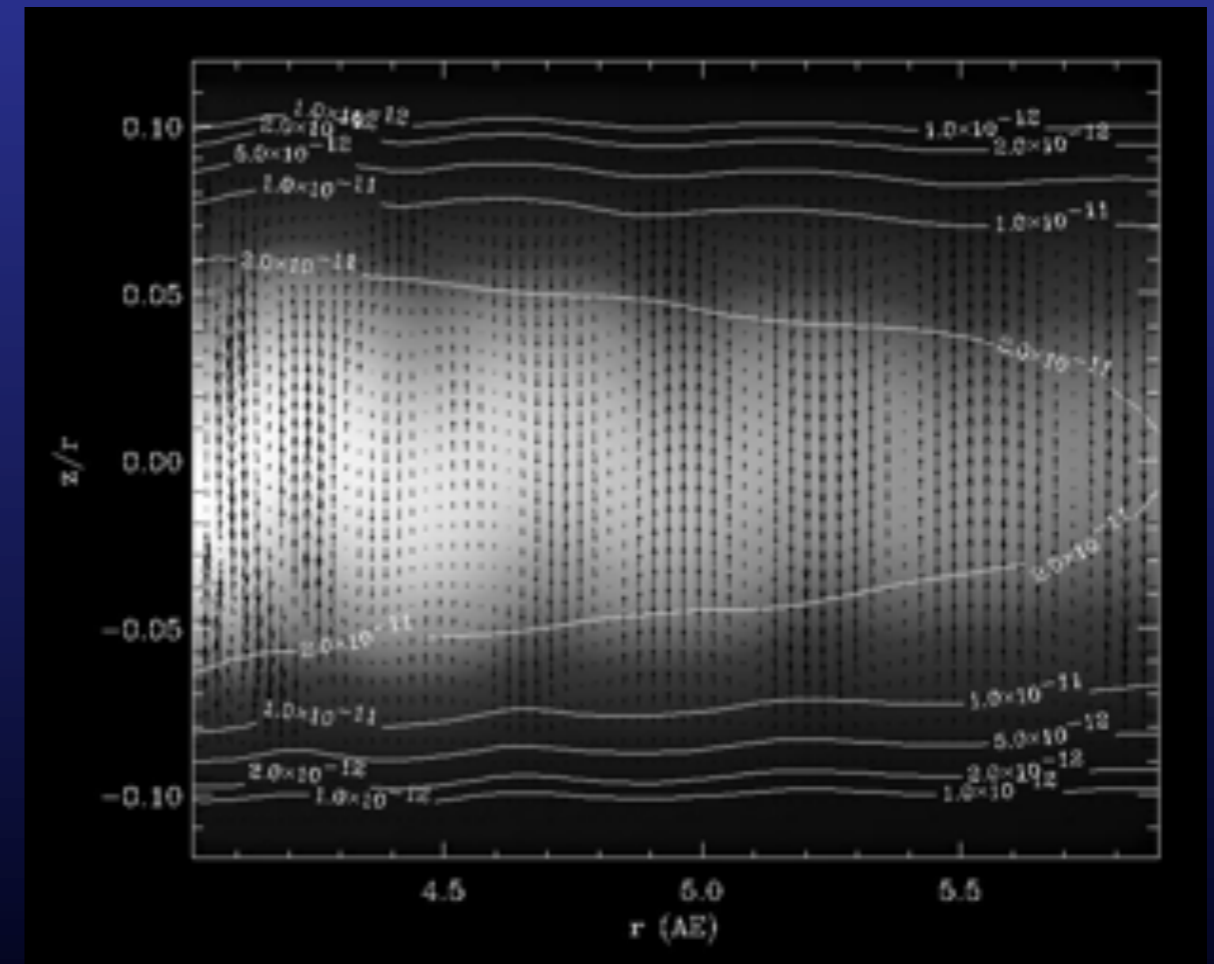
# TURBULENCE IN ACCRETION DISKS: VORTICITY GENERATION AND ANGULAR MOMENTUM TRANSPORT VIA THE GLOBAL BAROCLINIC INSTABILITY

H. H. KLAHR<sup>1</sup> AND P. BODENHEIMER  
UCO/Lick Observatory, University of California, Santa Cruz, CA 95064  
*Received 2000 June 7; accepted 2002 September 17*

KLAHR & BODENHEIMER



## 3D - Radiation Hydro of Convection in Disks. Klahr, Henning & Kley 1999



Klahr & Bodenheimer 2003

Radial Entropy Gradient leads to vortices  
somehow... Klahr 2004?

Then a lot of discussion started...

...but 4 years later:

Petersen et al. 2007

Yeah!

Now with thermal relaxation

Simulations are robust!

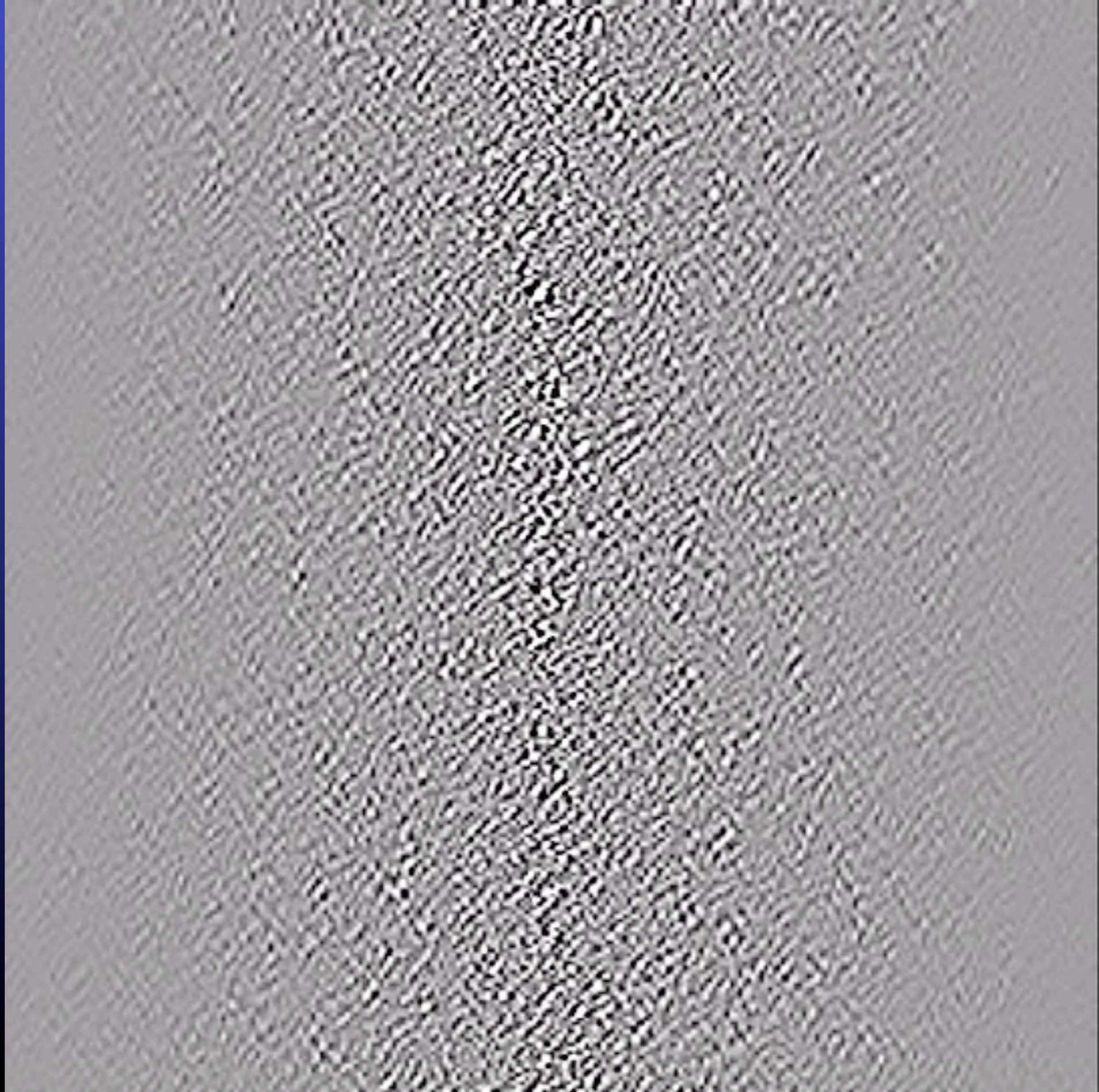


# Vorticity

Vorticity: Pencil Code: Lyra and  
Klahr 2011;  $\beta = 2$ ;  $N = 256$ ;  $\tau_c = 1$

# Vorticity

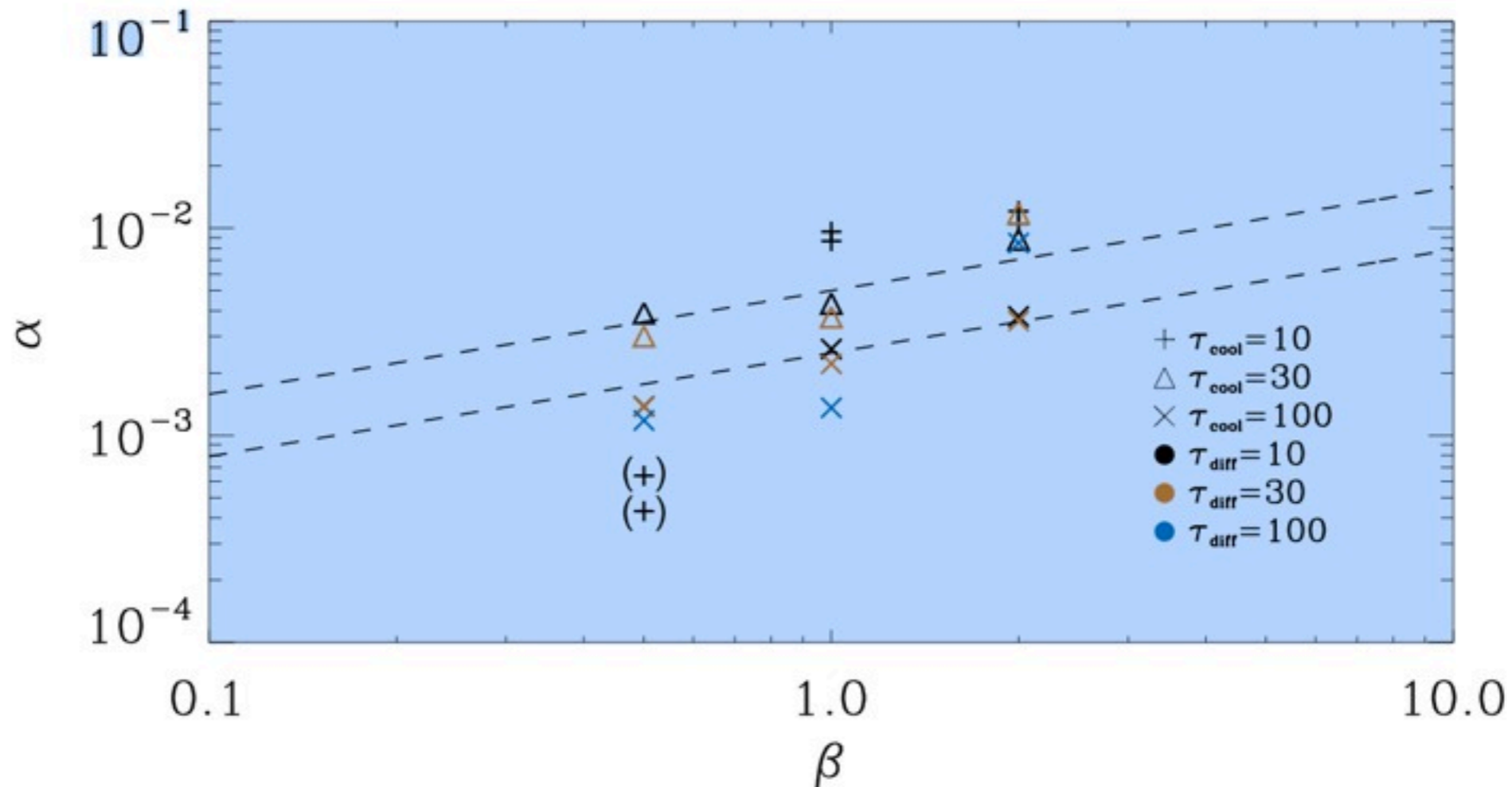
Vorticity: Pencil Code: Lyra and  
Klahr 2011;  $\beta = 2$ ;  $N = 256$ ;  $\tau_c = 1$



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# Strength of alpha (local)

Raettig, Klahr and Lyra, submitted



Growth rates

$$\Gamma = -\tau N_R^2 = -\frac{a^2}{D} N_R^2$$

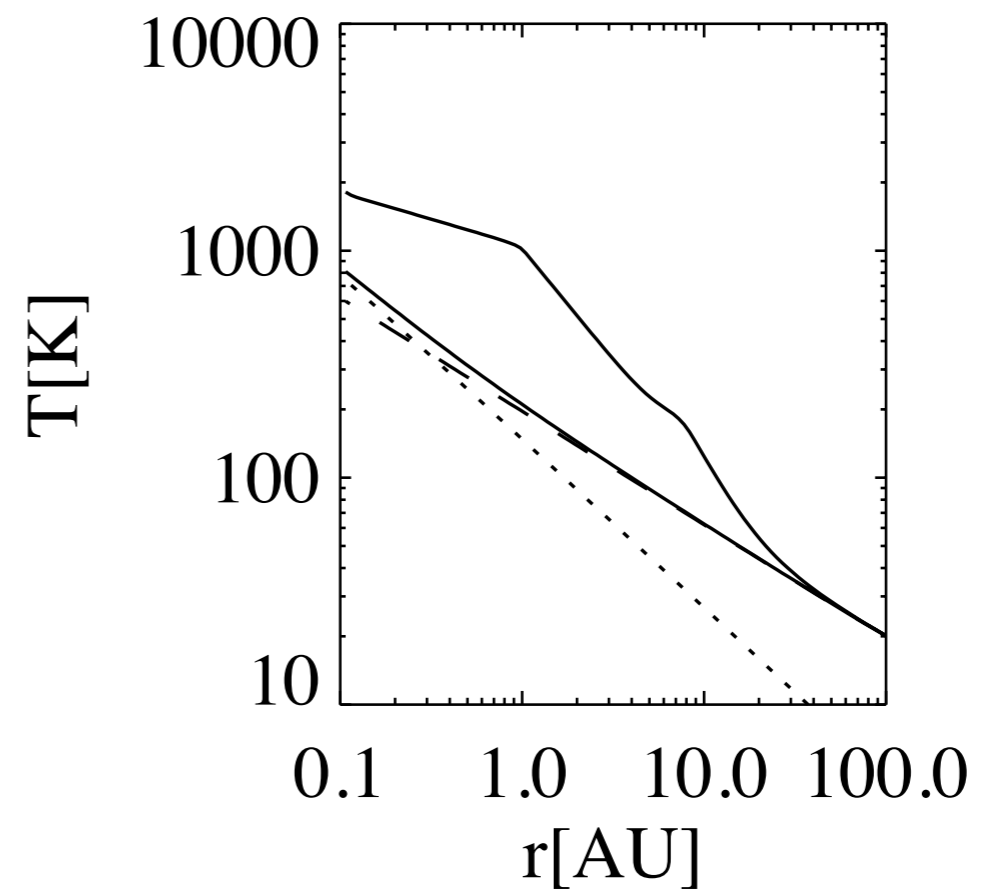
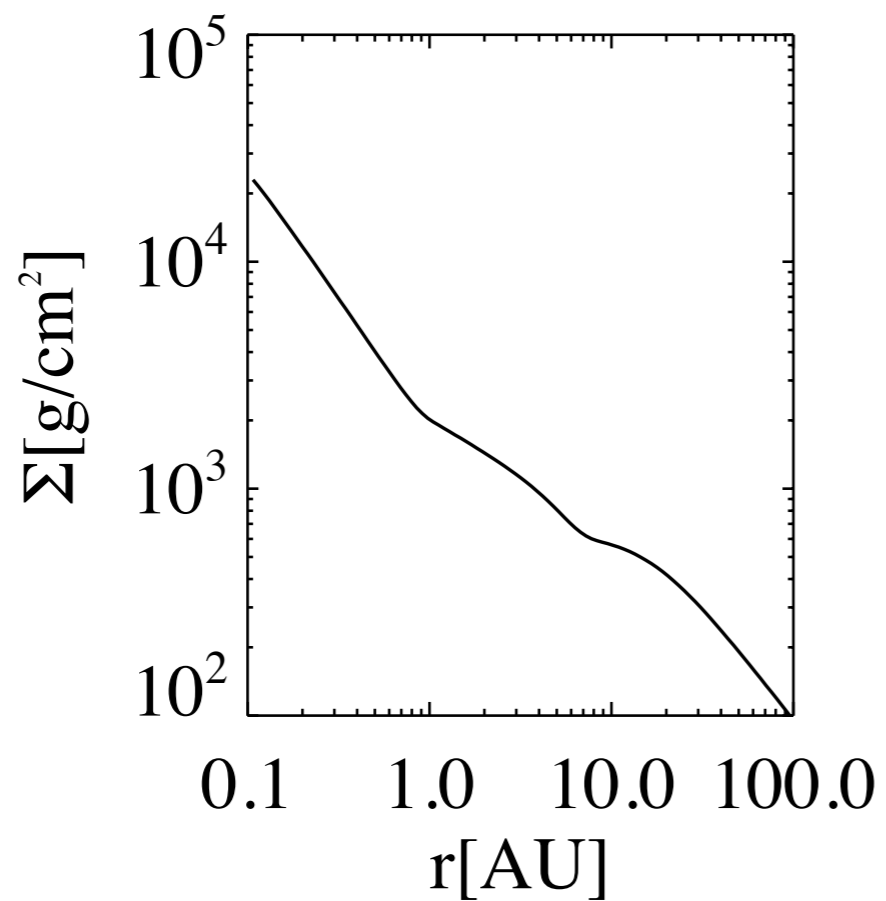
see also John's talk

Klahr & Raettig in prep.

Using data from Sean Andrews: Accretion disks in Ophiuchus.

$\alpha = 0.001$ ;  $\dot{M} = 1E-7 \text{ Msol/yr}$ ;

Plus irradiation:  $T_{\text{star}} = 4300$ ;  $R_{\text{star}} = 2 R_{\text{sol}}$

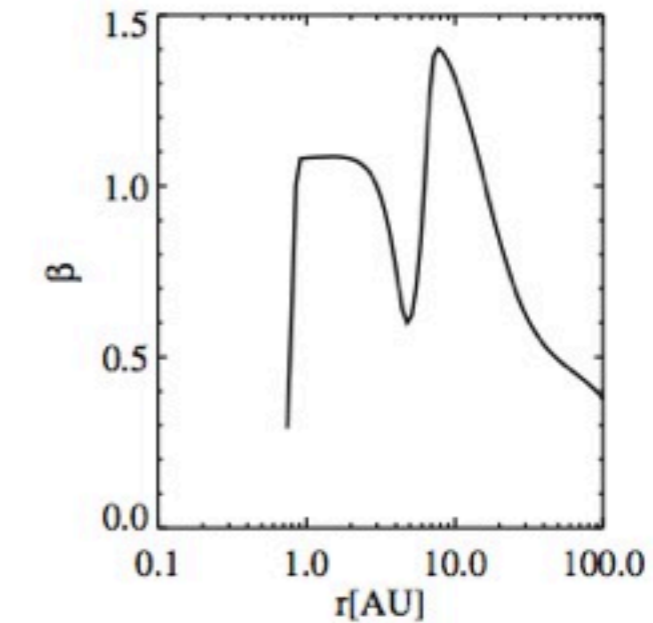
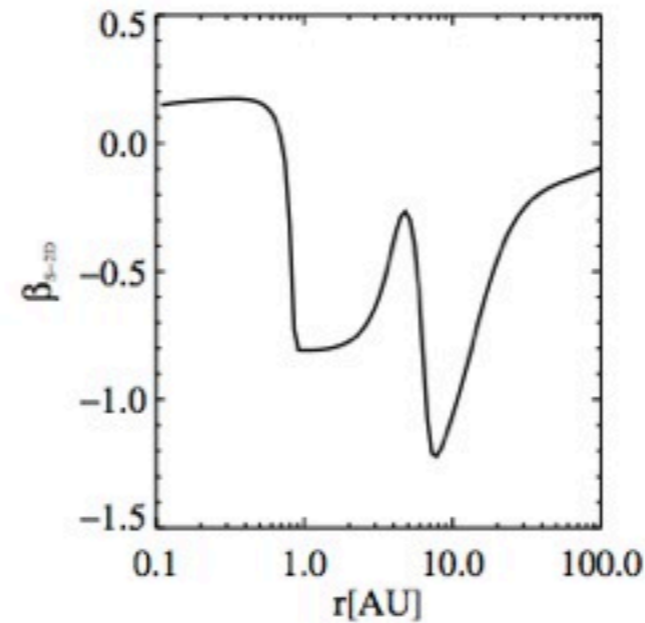
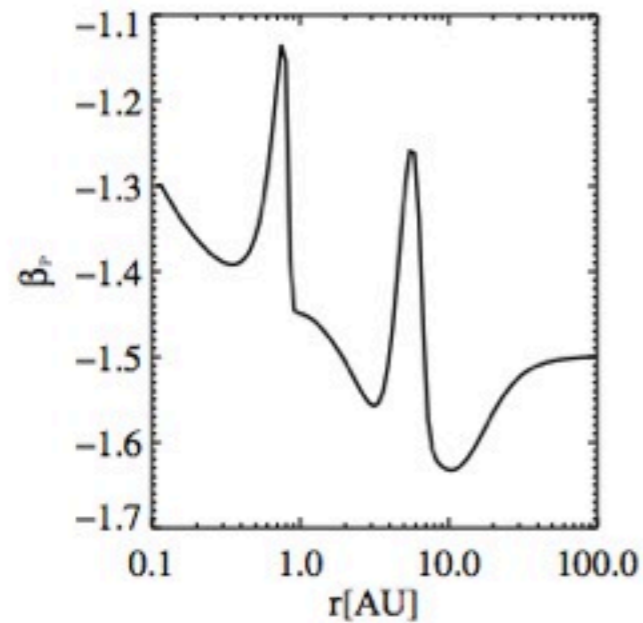
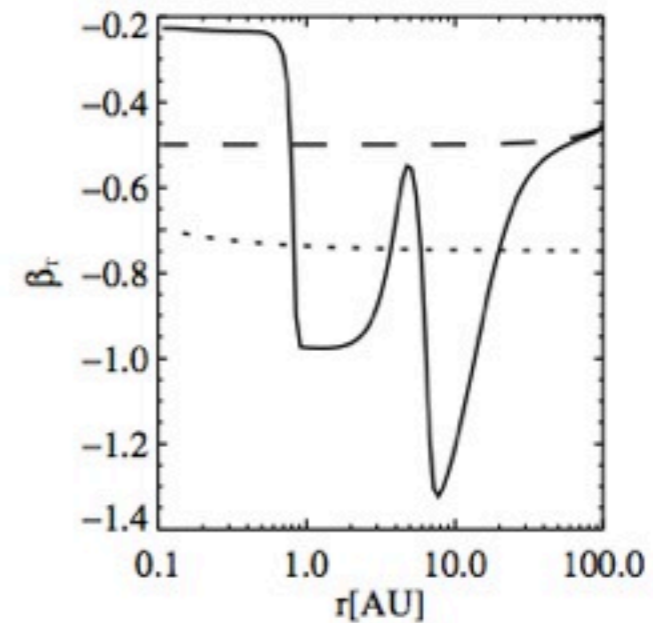
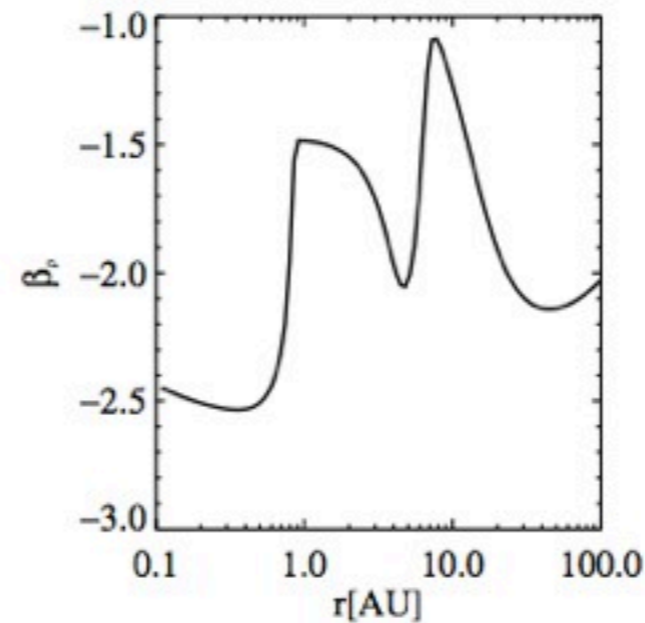
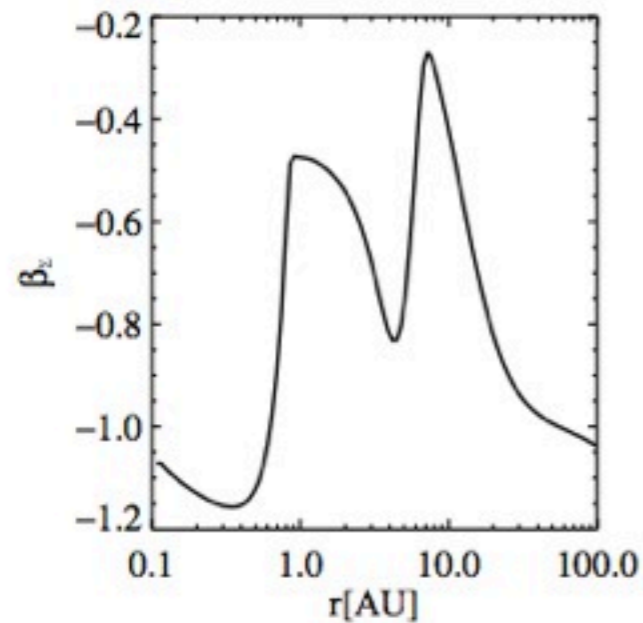


Klahr & Raettig in prep.

Using data from Sean Andrews

$\alpha = 0.001$ ;  $\dot{M} = 1E-7 \text{ Msol/yr}$ ;

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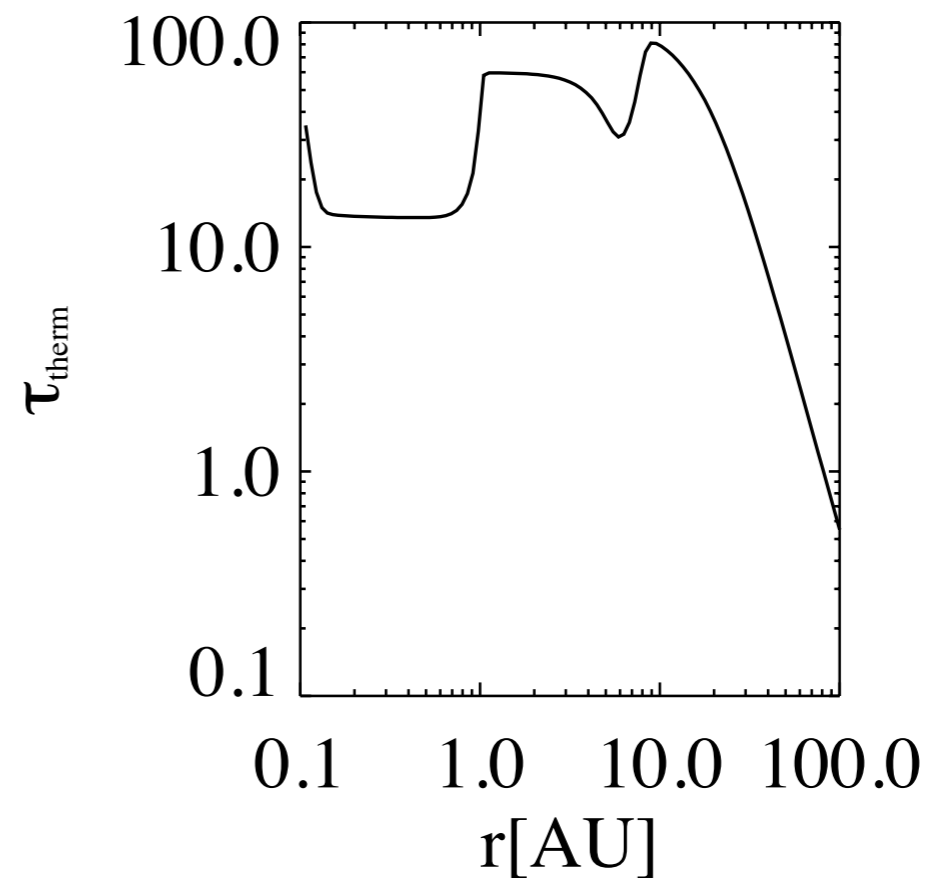
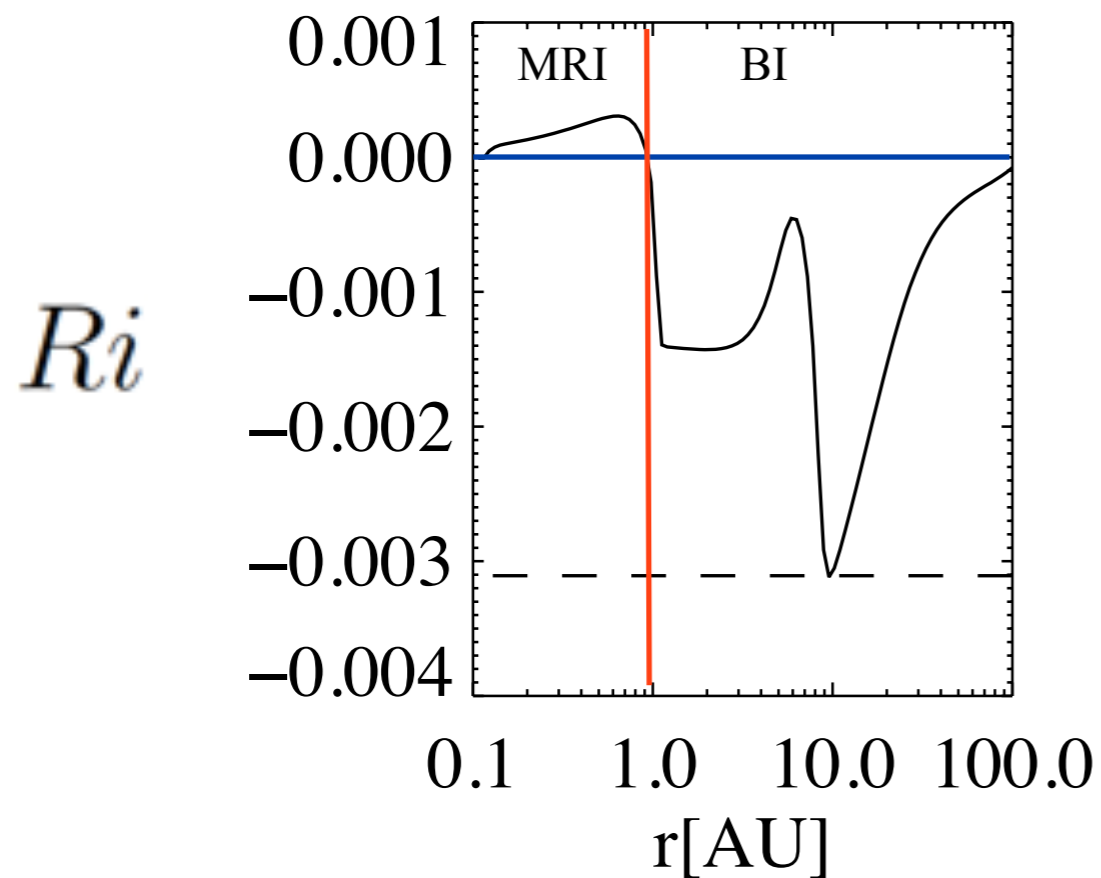
# Richardson number & thermal diffusion time

$$N^2 = -\frac{1}{\gamma} \left(\frac{H}{R}\right)^2 \beta_s \beta_p \Omega^2$$

$$Ri = -\frac{2}{3\gamma} \left(\frac{H}{R}\right)^2 \beta_p \beta_s$$

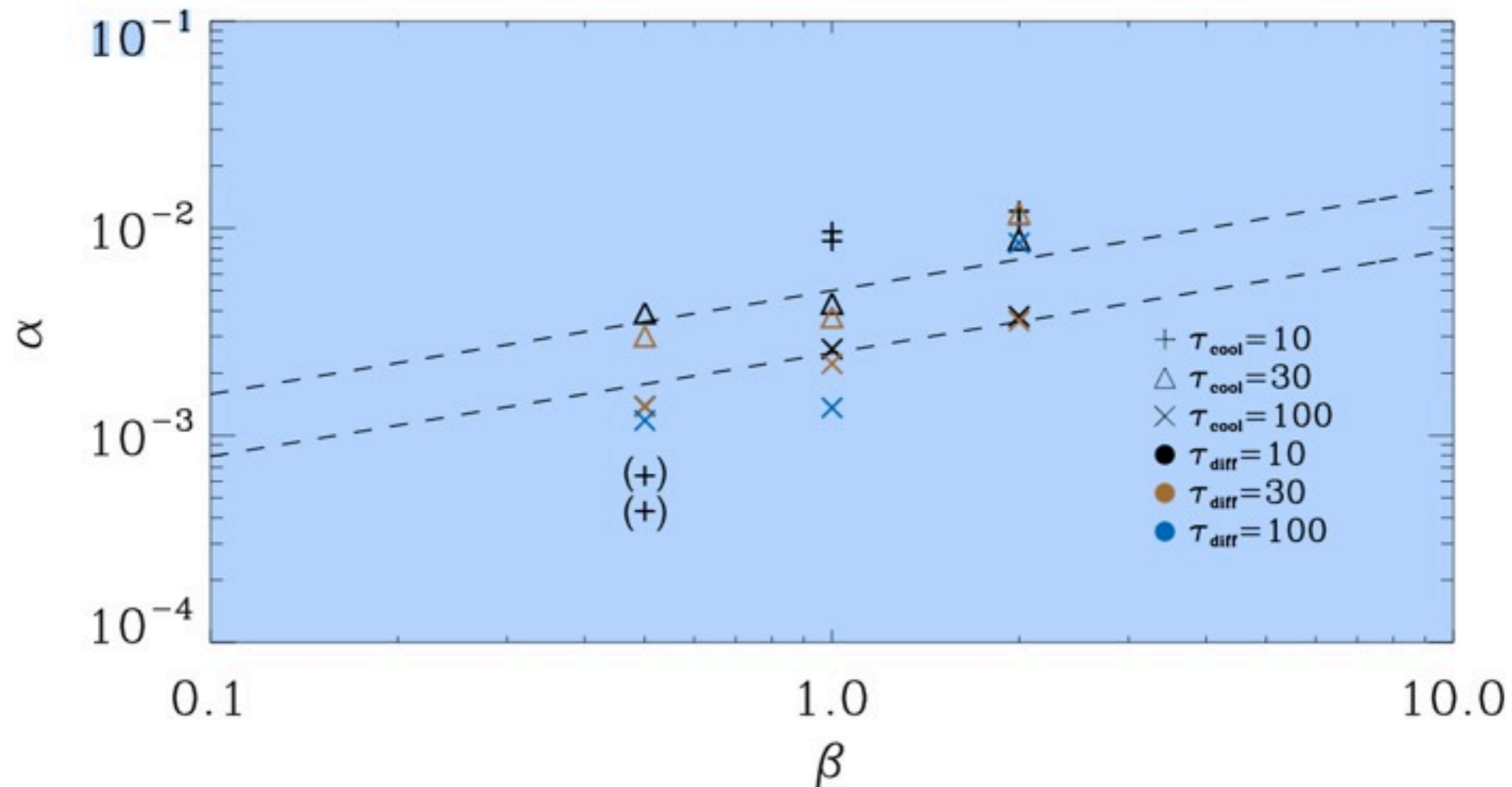
$$D = \frac{\lambda c 4 a_R T^3}{\rho(\kappa + \sigma)},$$

$$\tau_{therm} = H^2 / \frac{D}{\rho C_v}$$



# Strength of alpha (local)

Raettig, Klahr and Lyra, submitted



Growth rates

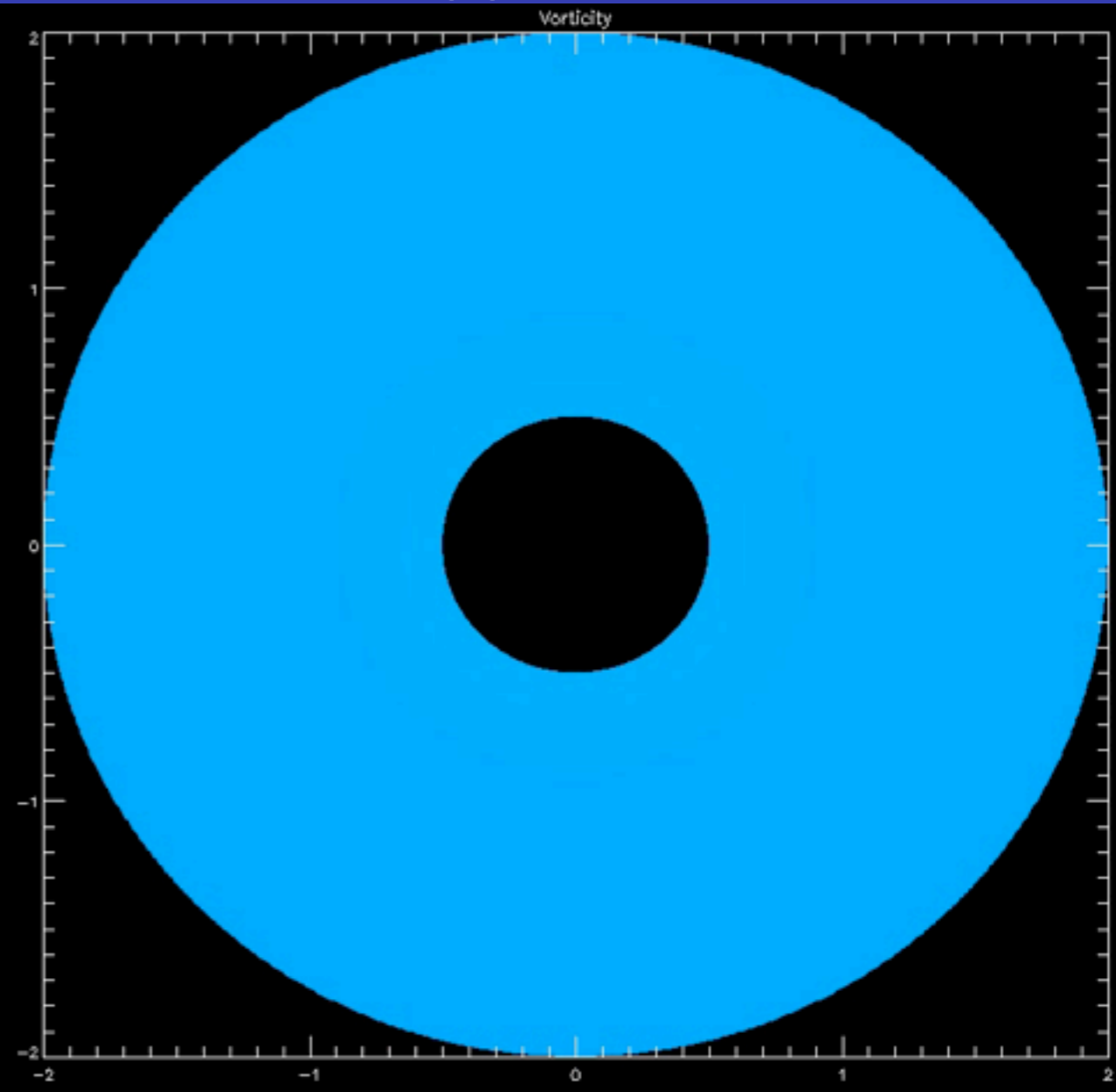
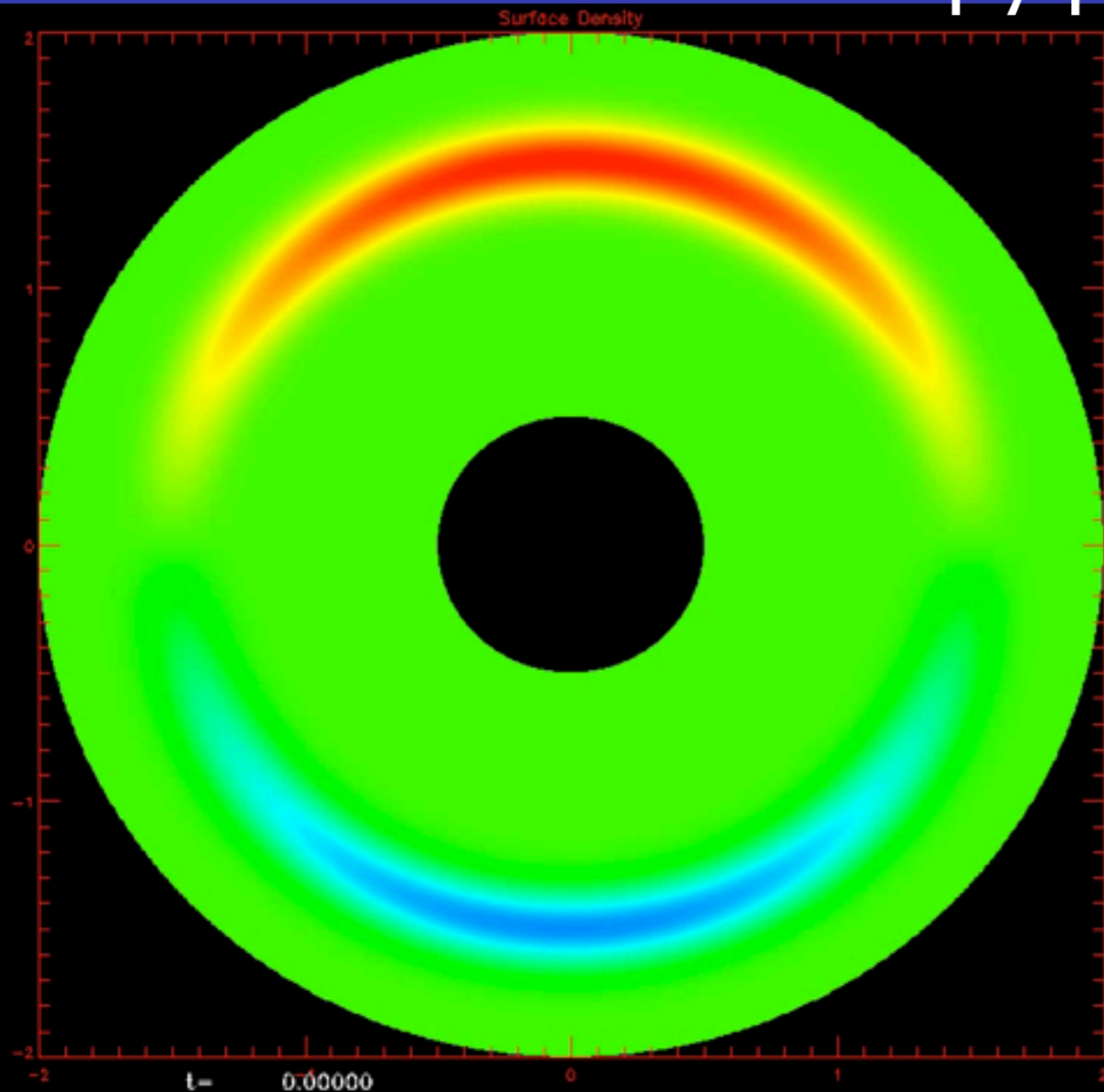
$$\Gamma = -\tau N_R^2 = -\frac{a^2}{D} N_R^2$$

see also John's talk

Pluto Code:  $1024^2$ ; WENO3-RK3; HLLE; FARGO

How to get a critical Vortex?

Here entropy pert. + Klahr 2004!



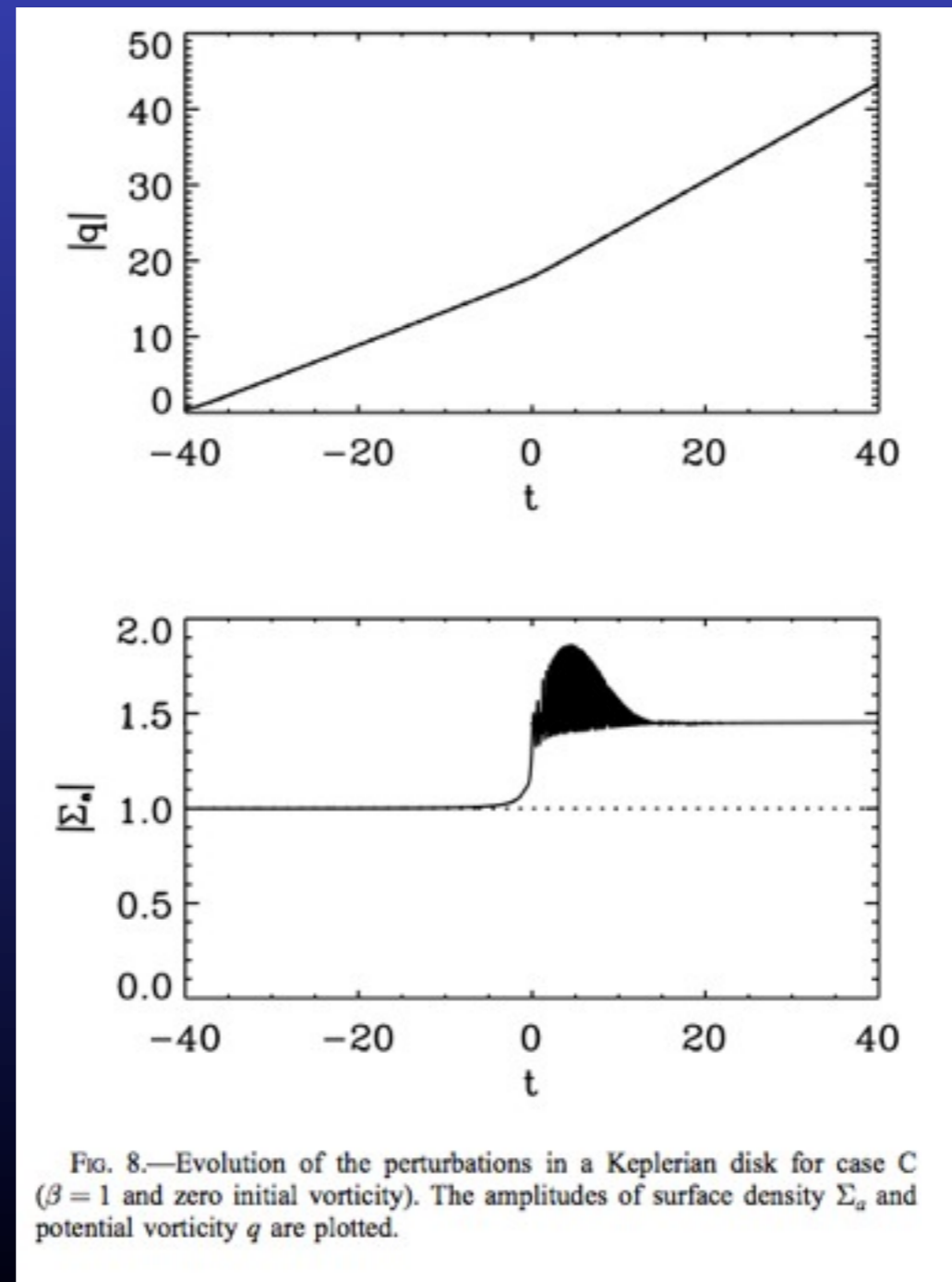
vortices migrate inward, but are recreated by waves from other vortices.



Pluto Code:  $1024^2$ ; WENO3-RK3; HLLE; FARGO

How to get a critical Vortex?

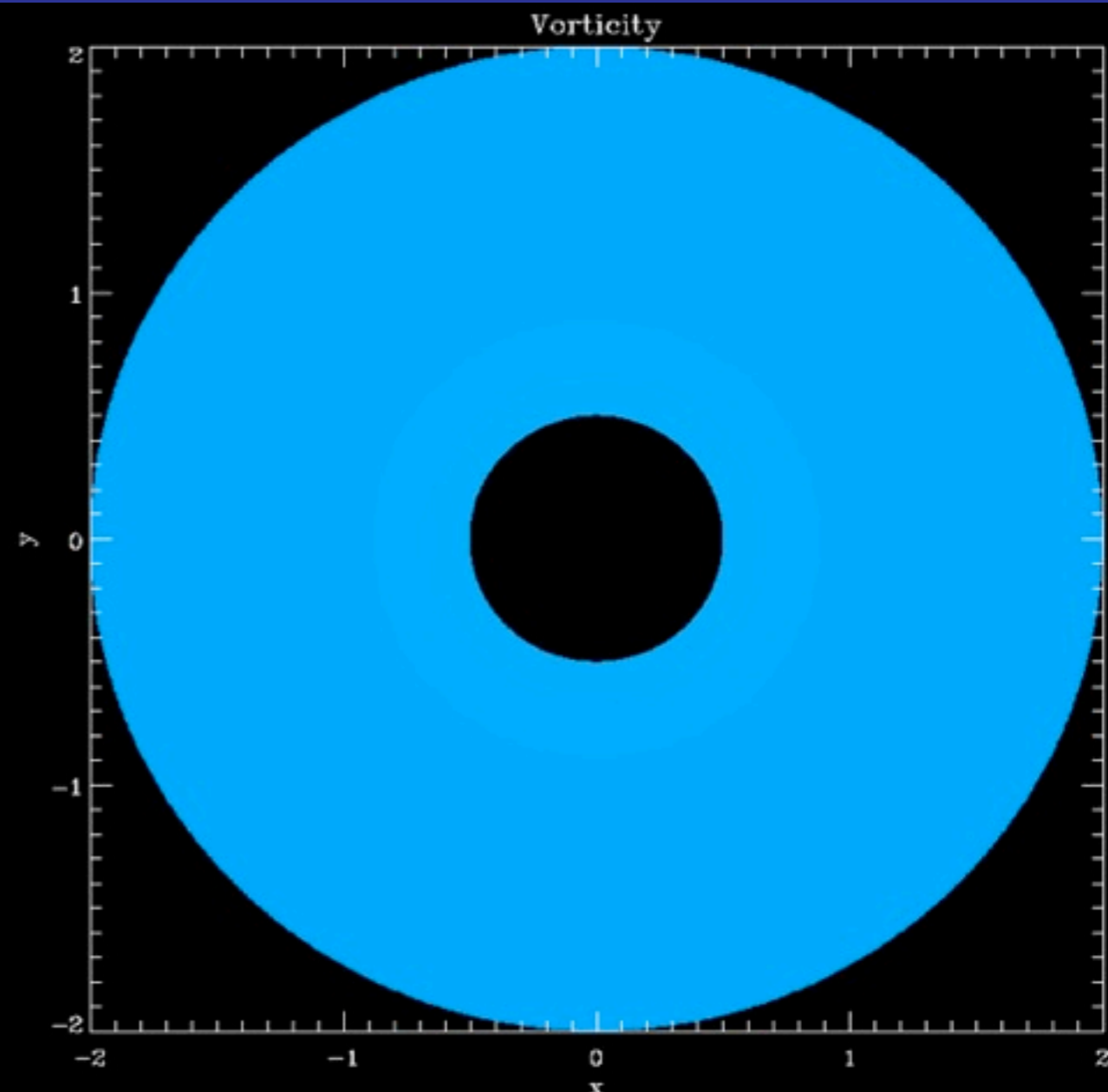
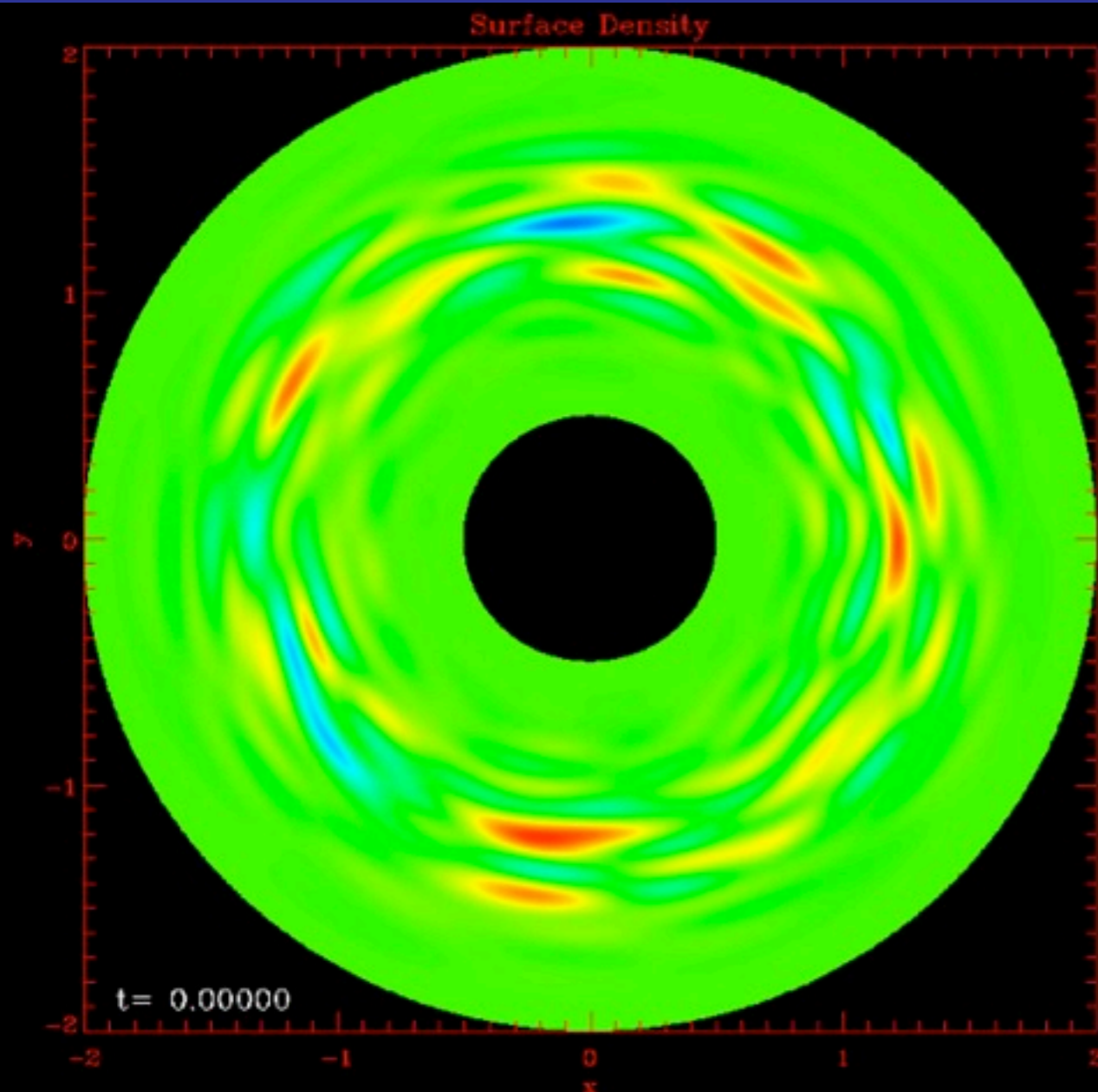
Here entropy pert. + Klahr 2004!



Pluto Code:  $1024^2$ ; WENO3-RK3; HLLE; FARGO

## Non linear stage of Klahr 2004

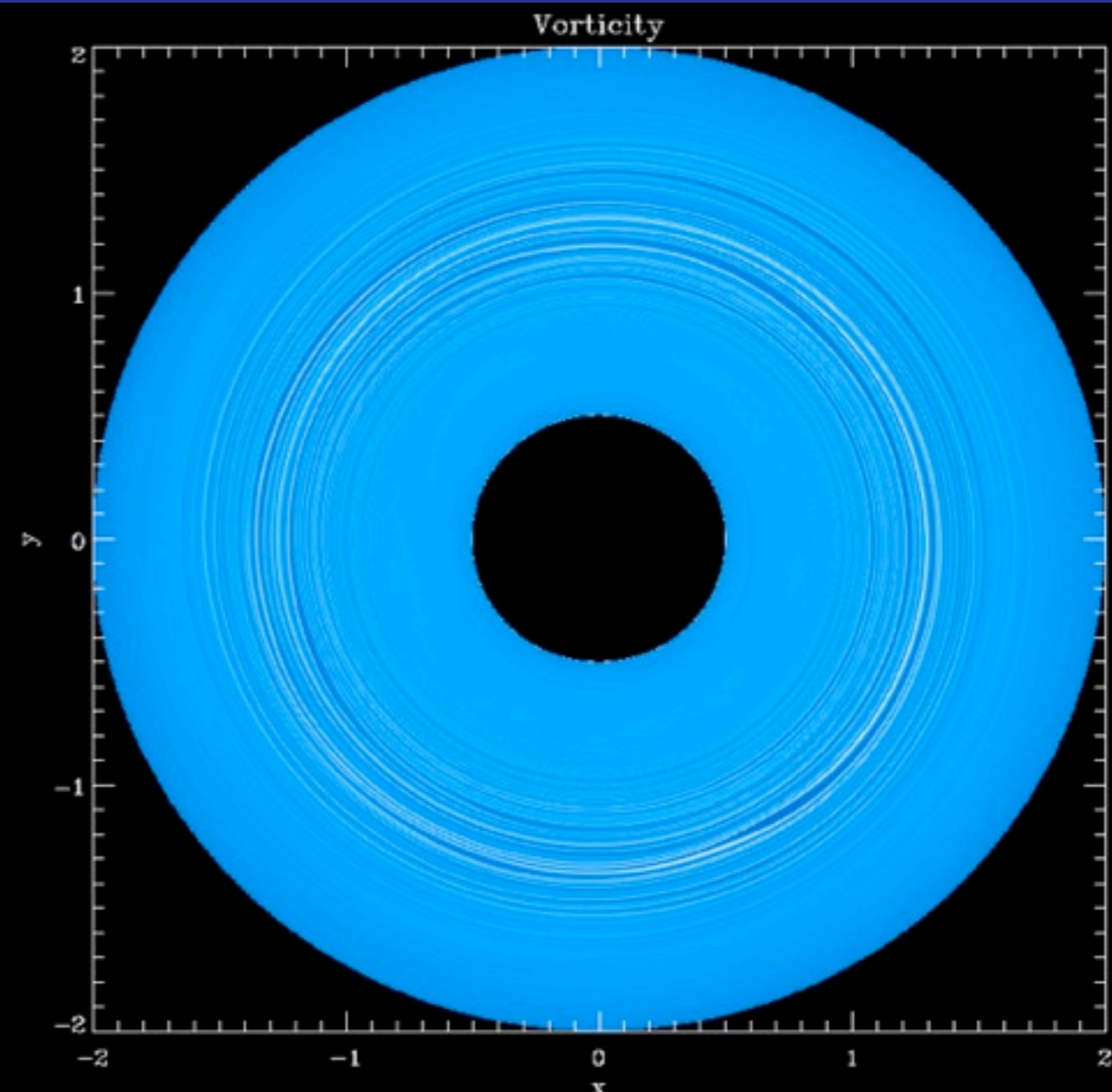
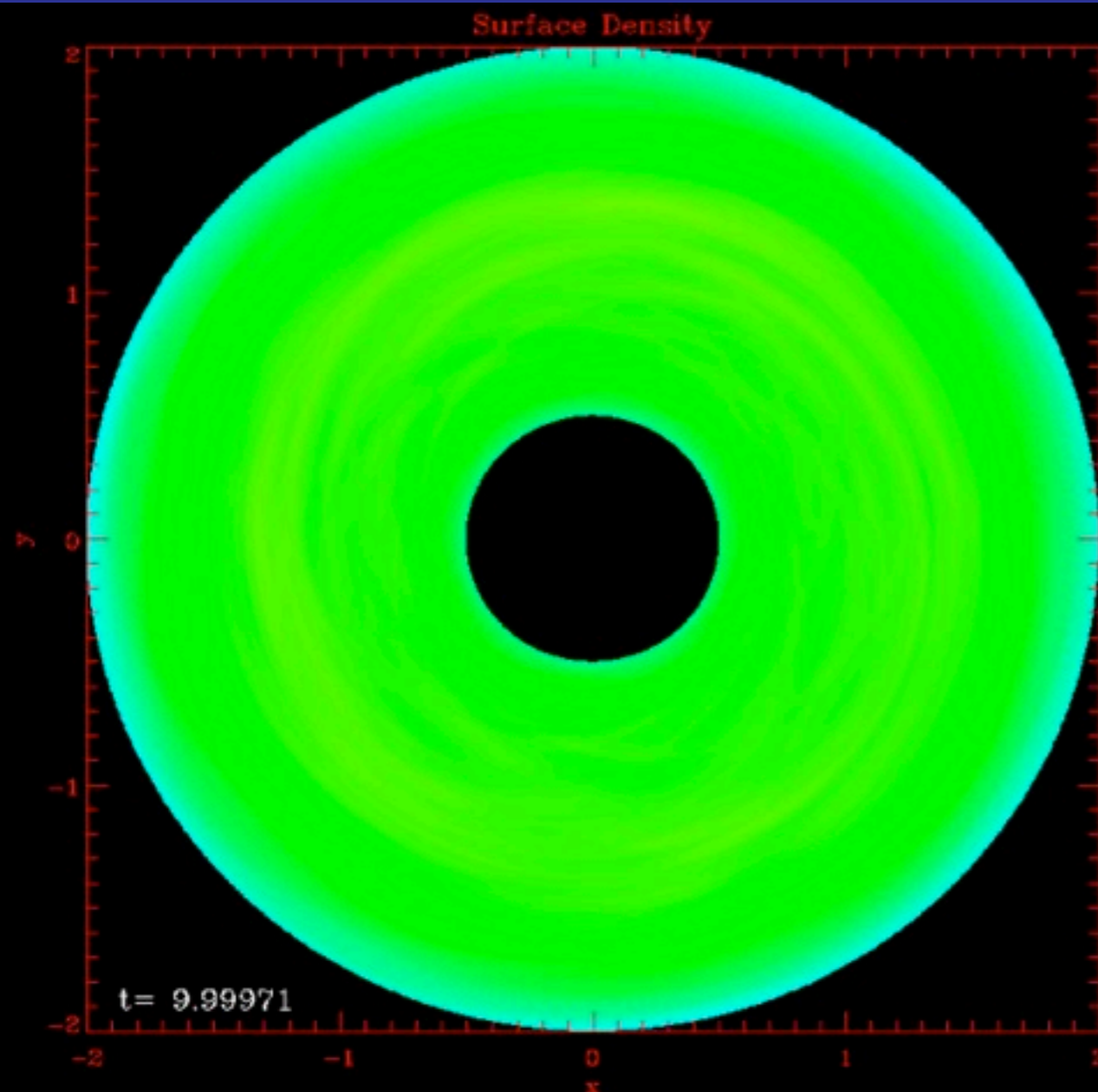
- 1.) entropy perturbation
- 2.) tightly wrapped growing perturbation of vorticity
- 3.) Papaloizou Pringle - RWI Instability
- 4.) SBI - Vortex Amplification
- 5.) Spirals - GOTO 2.)



Pluto Code:  $1024^2$ ; WENO3-RK3; HLLE; FARGO

## Non linear stage of Klahr 2004

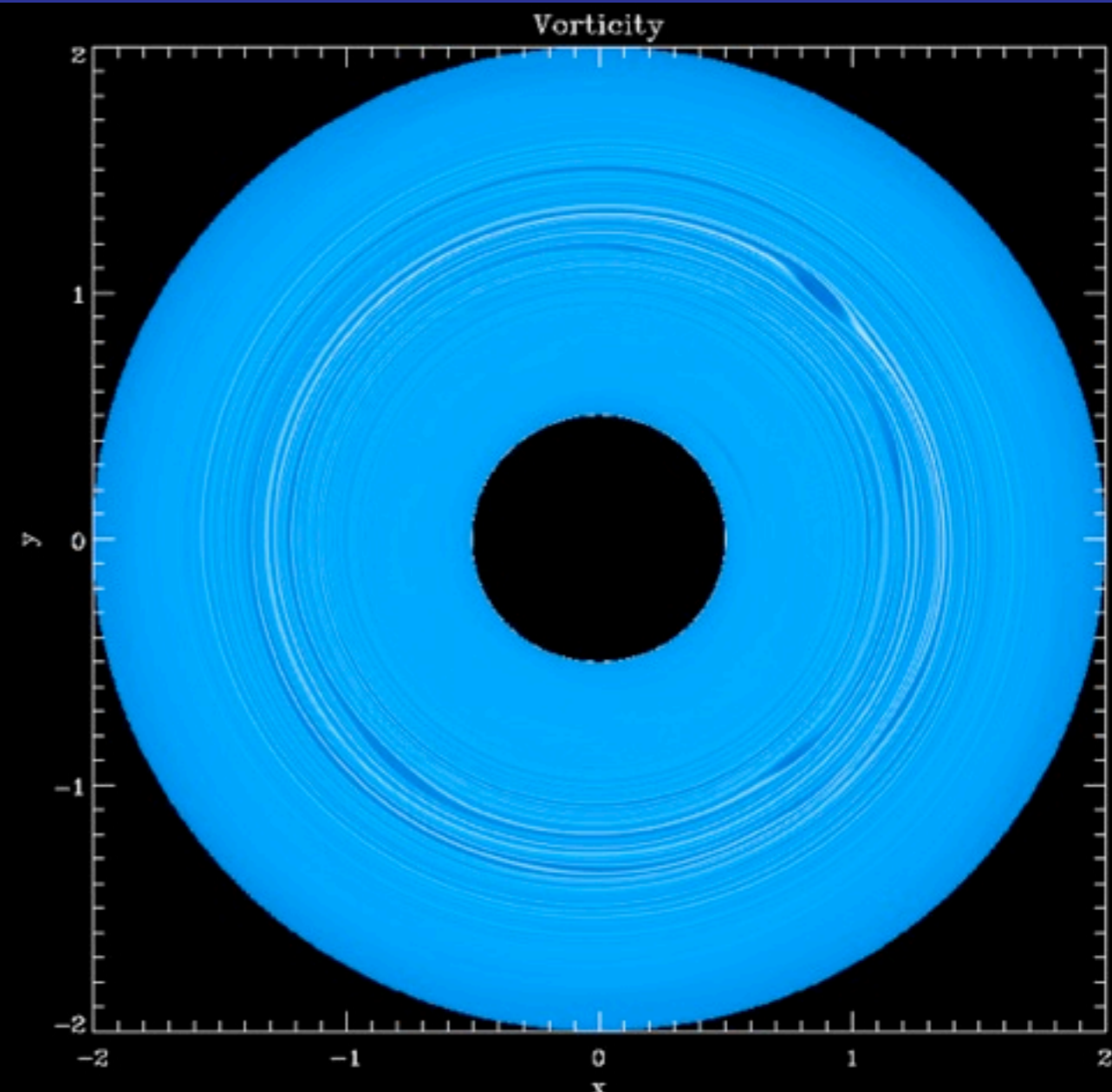
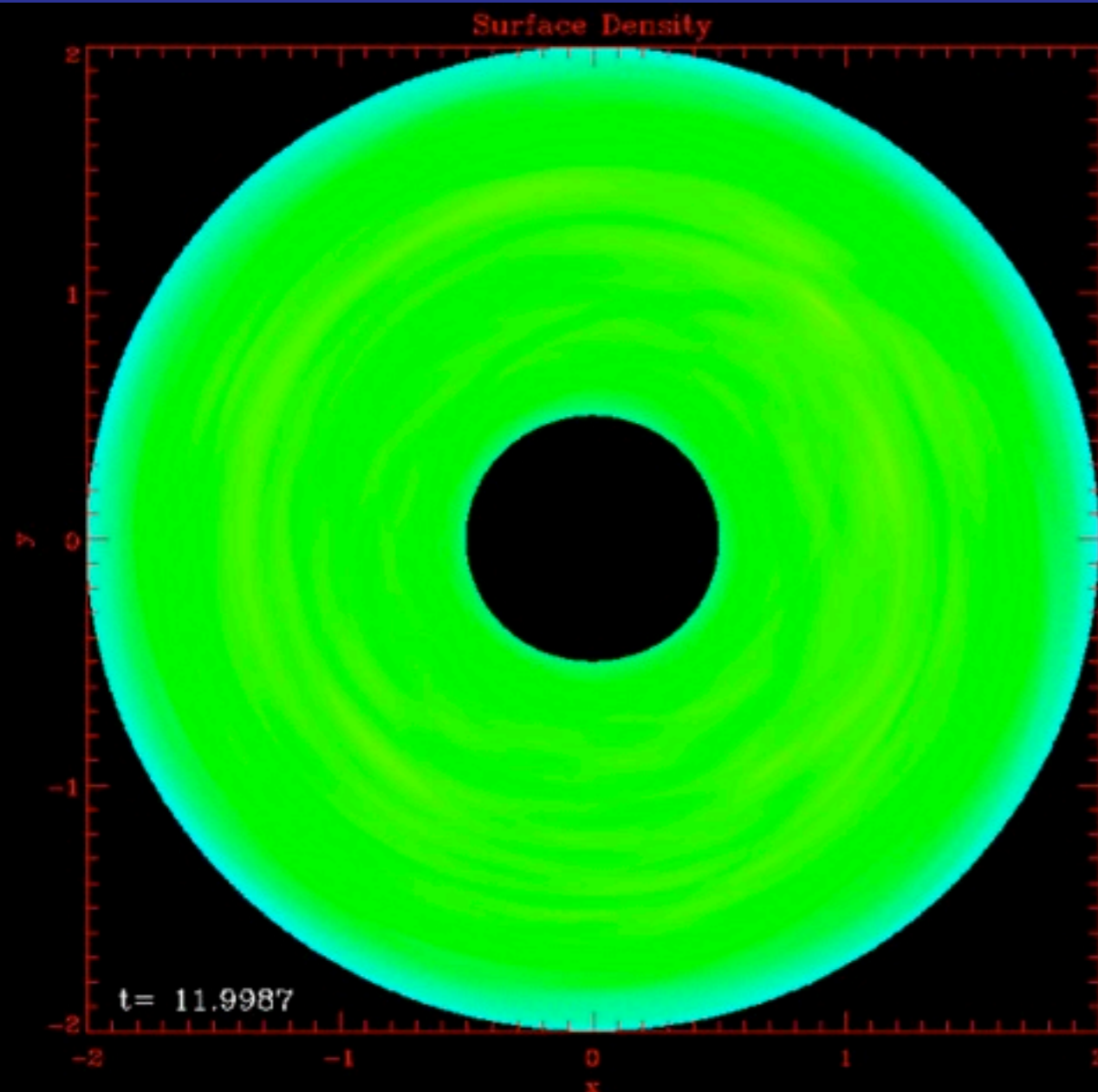
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- 5.) Spirals - GOTO 2.)



Pluto Code: 1024<sup>2</sup>; WENO3-RK3; HLLE; FARGO

## Non linear stage of Klahr 2004

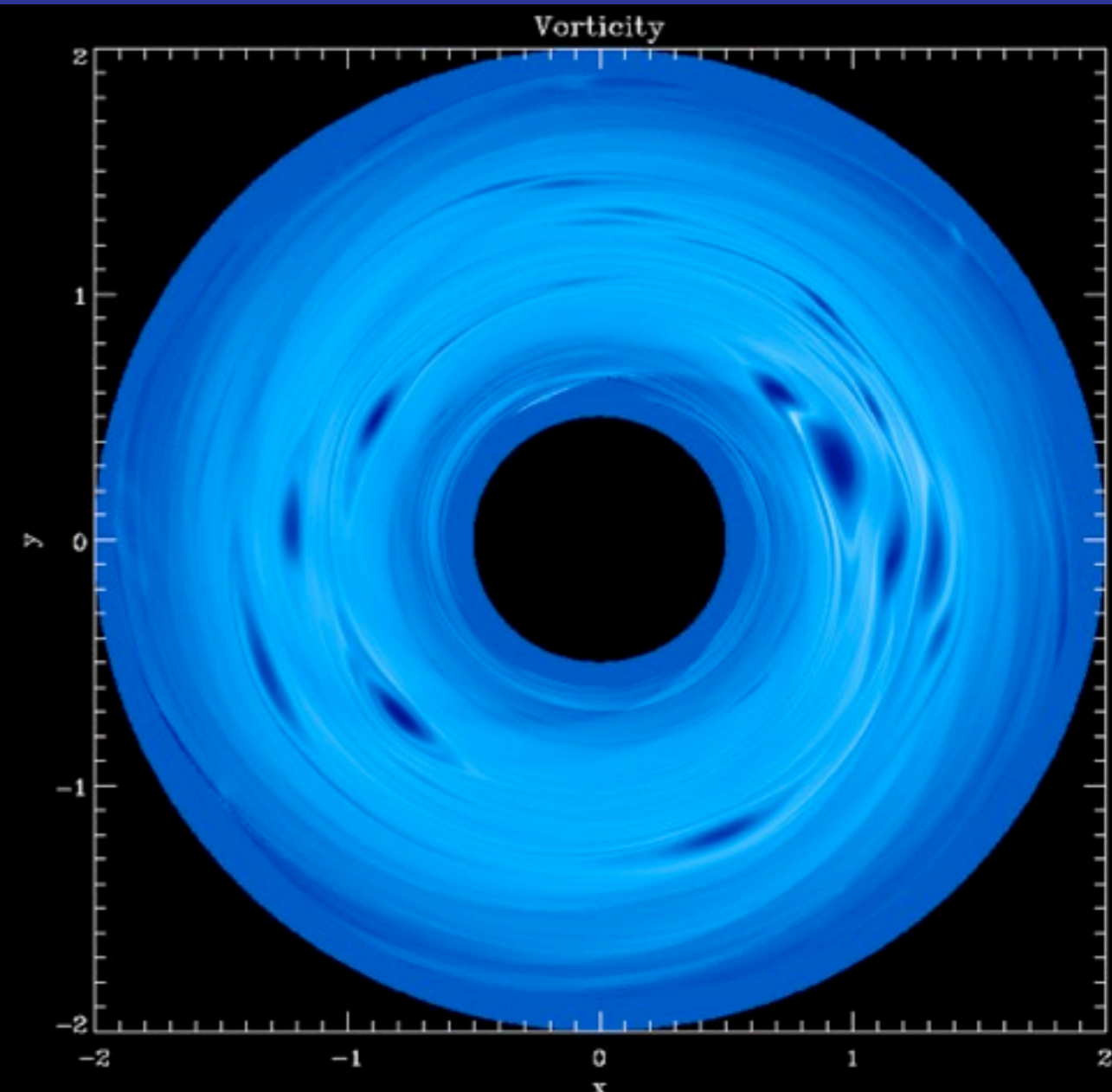
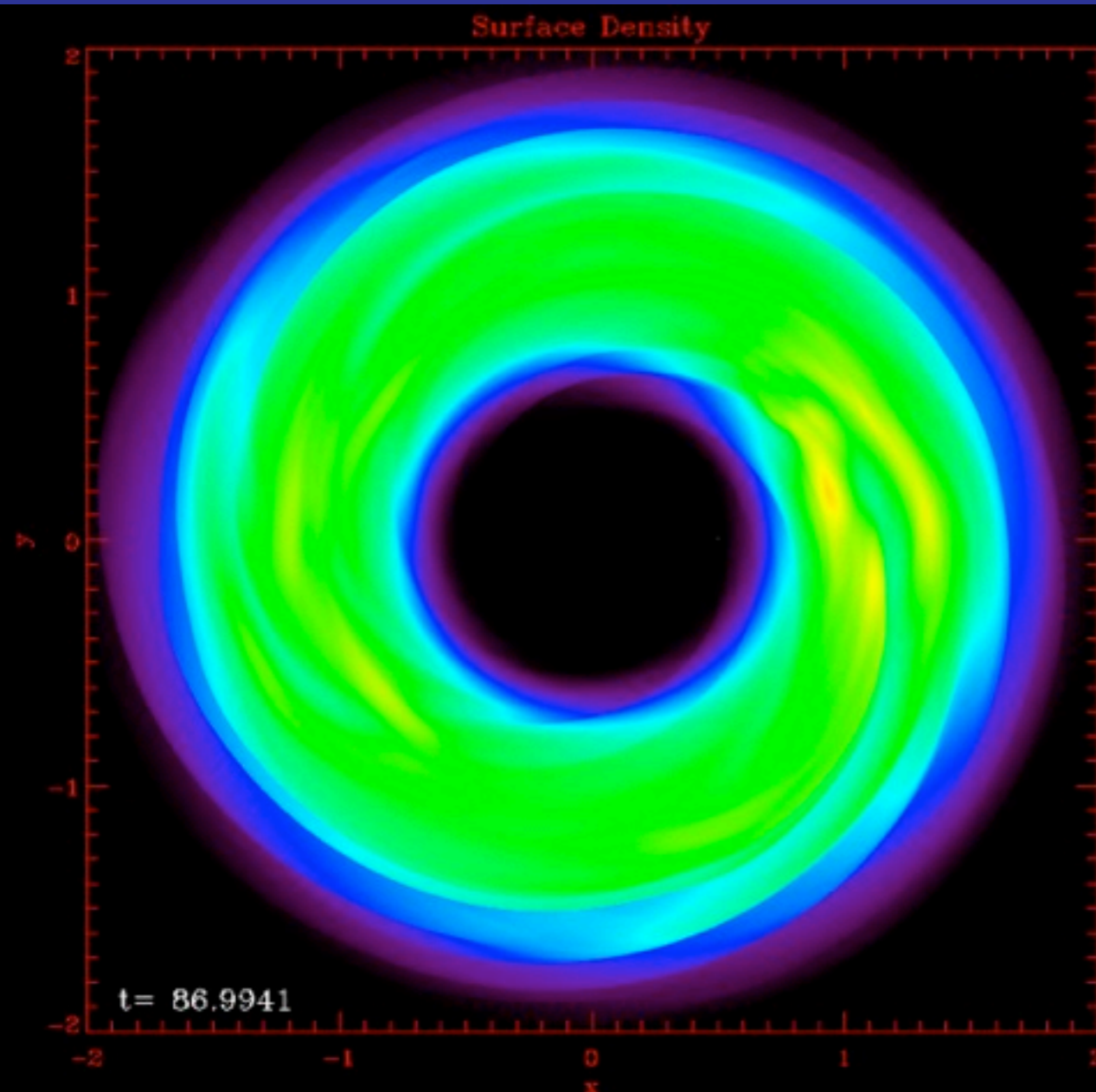
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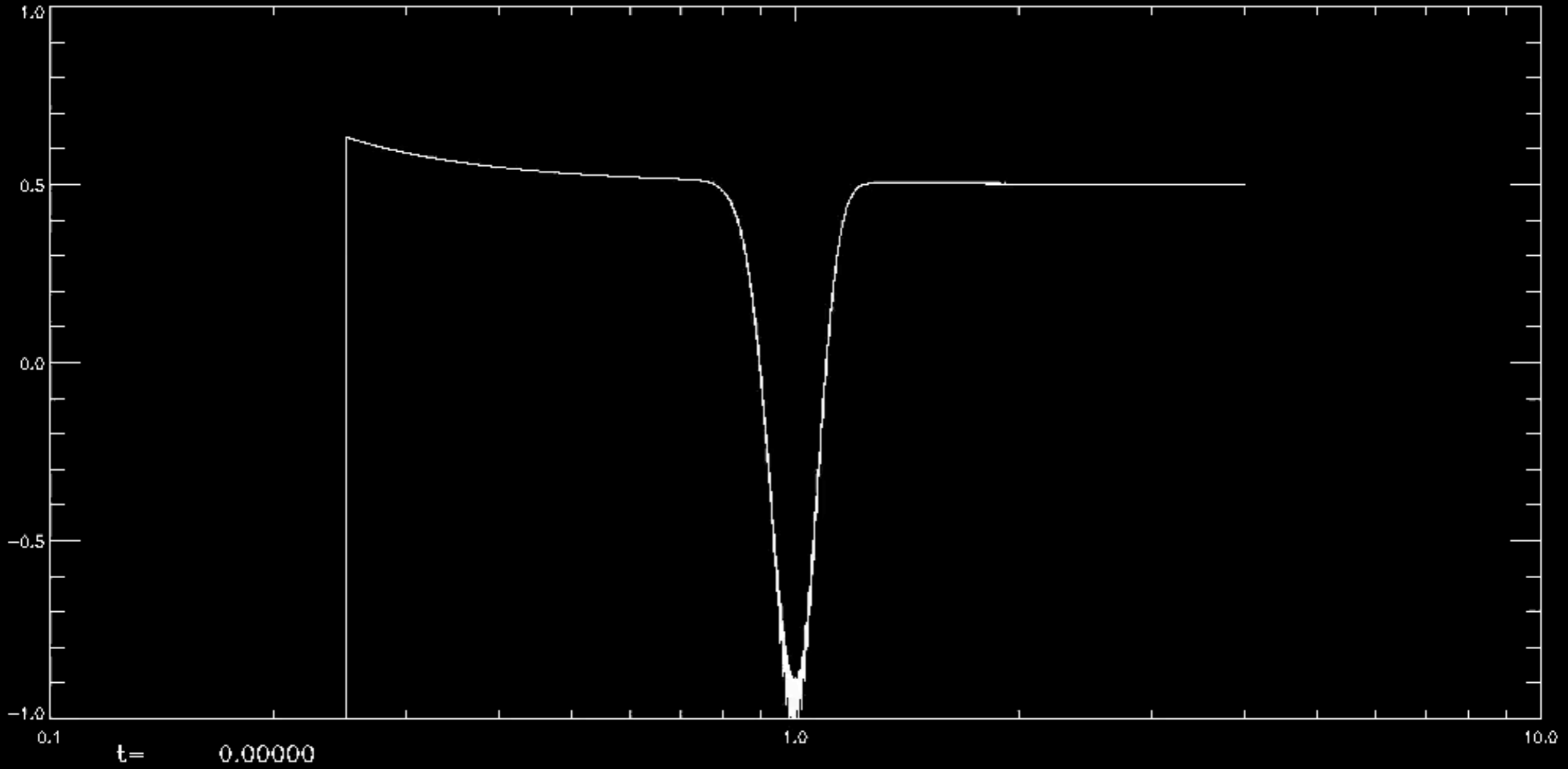
Pluto Code:  $1024^2$ ; WENO3-RK3; HLLE; FARGO

## Non linear stage of Klahr 2004

- 1.) entropy perturbation
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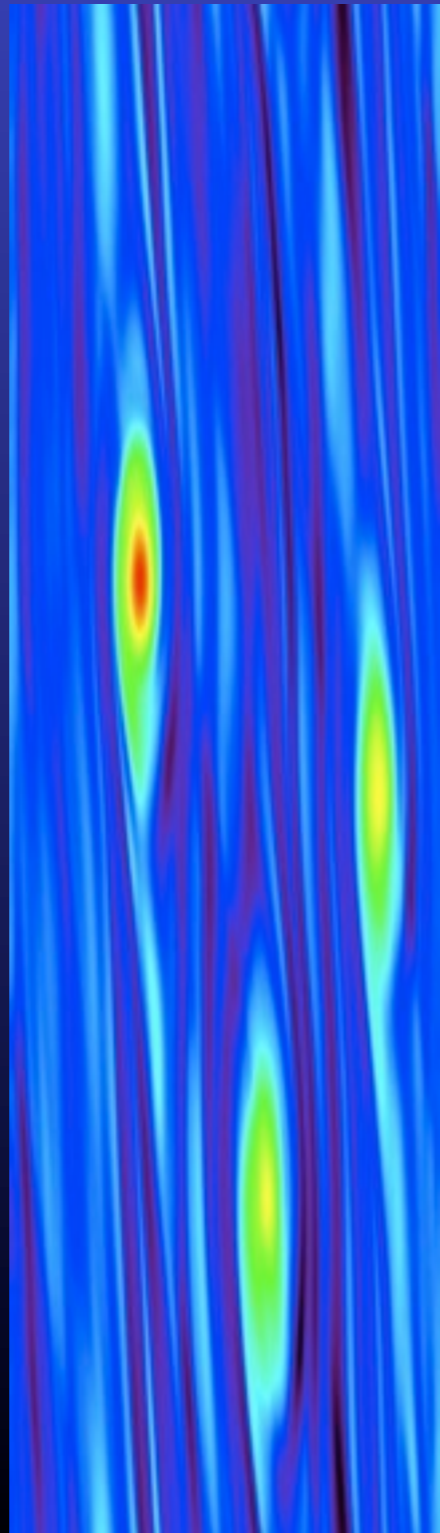


# Vortex Migration: Local minimum of vorticity:



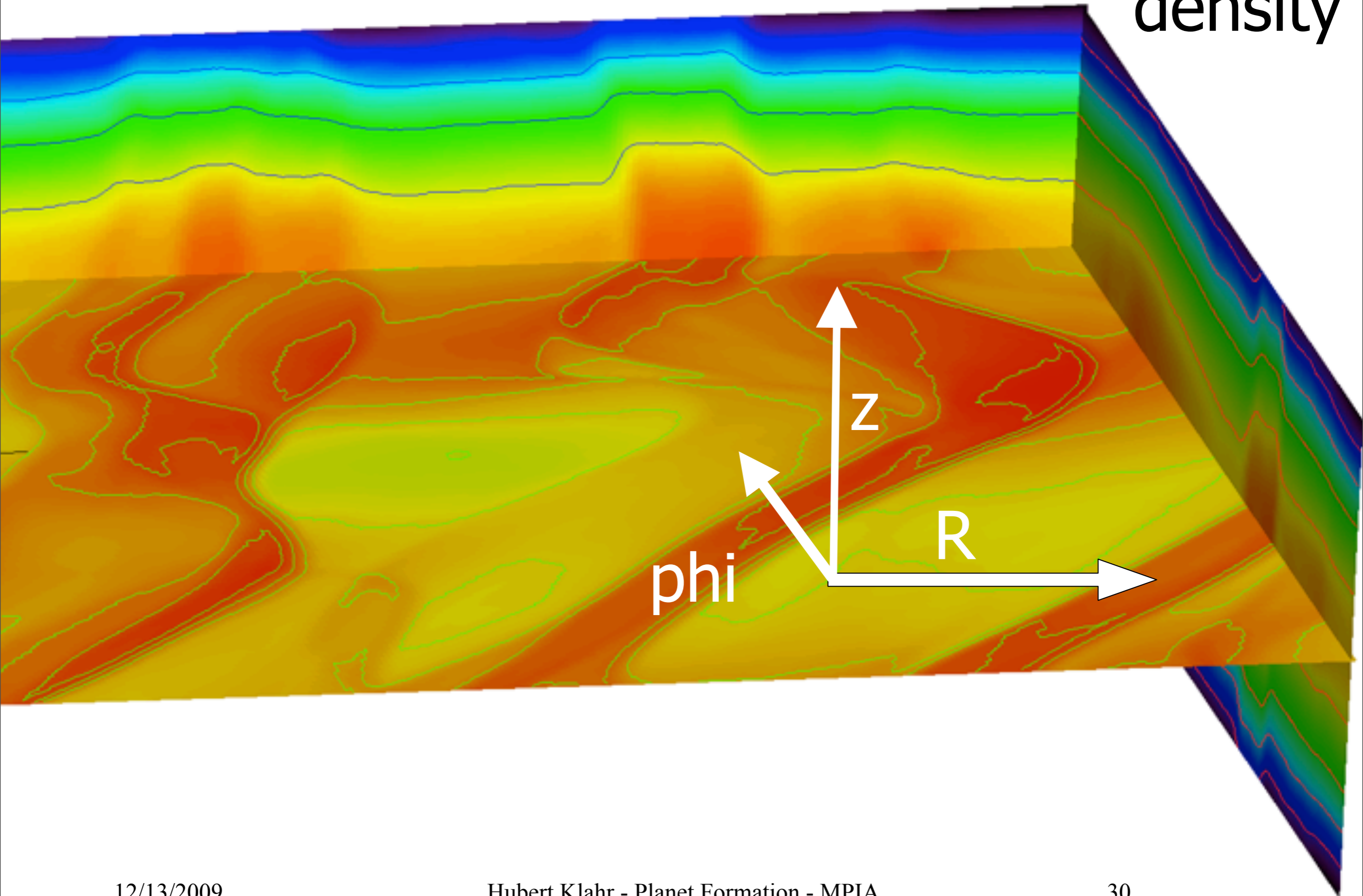
- Parameter: 3D;  $256^2 \times 128$  LX = 4H;
- LY = 16H; LZ = 4H;  $\beta = 2$ ;  $t_{\text{cool}} = 1$ ;  $H/R = 0.1$

azimuthal



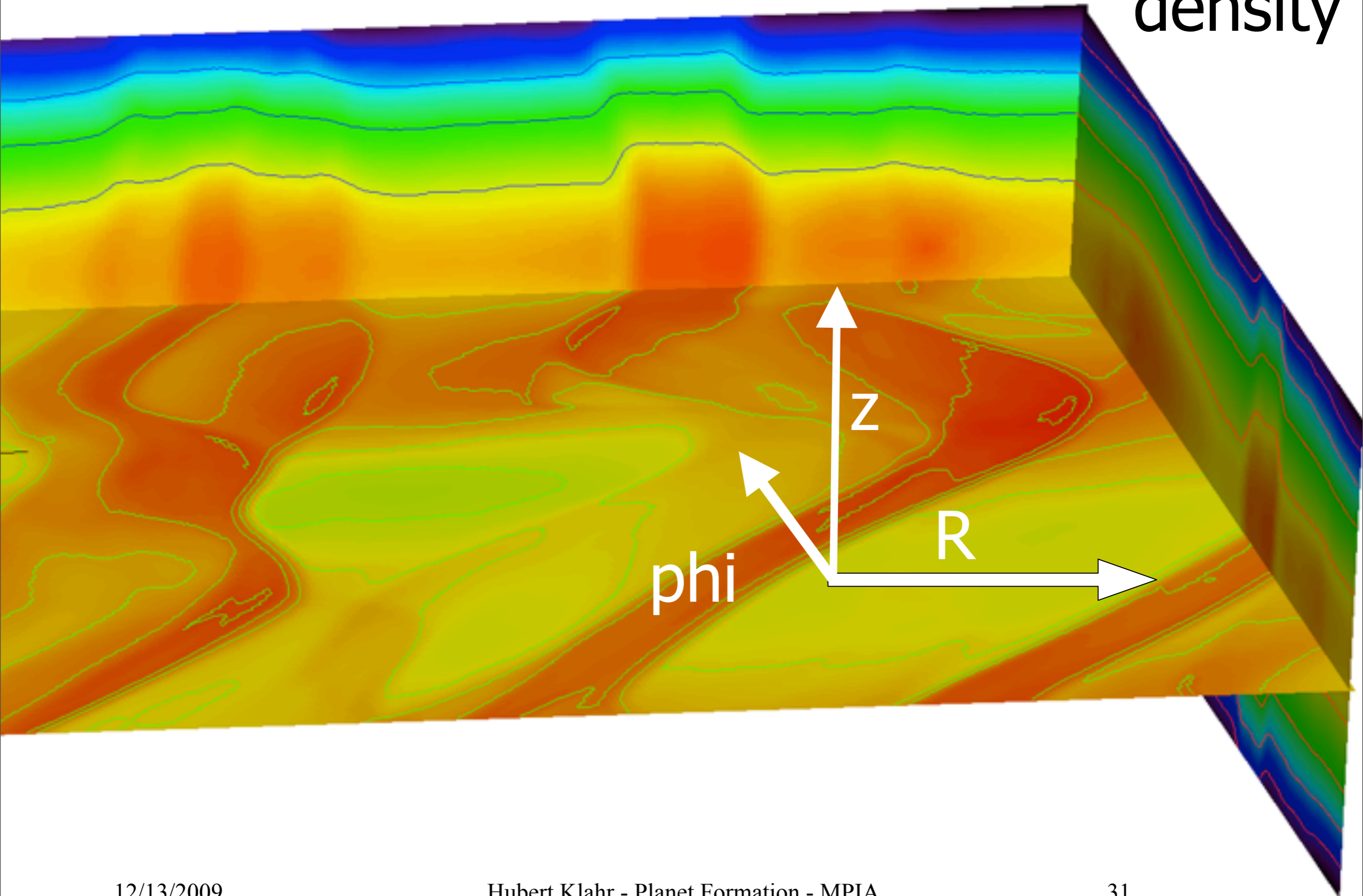
Radial

density

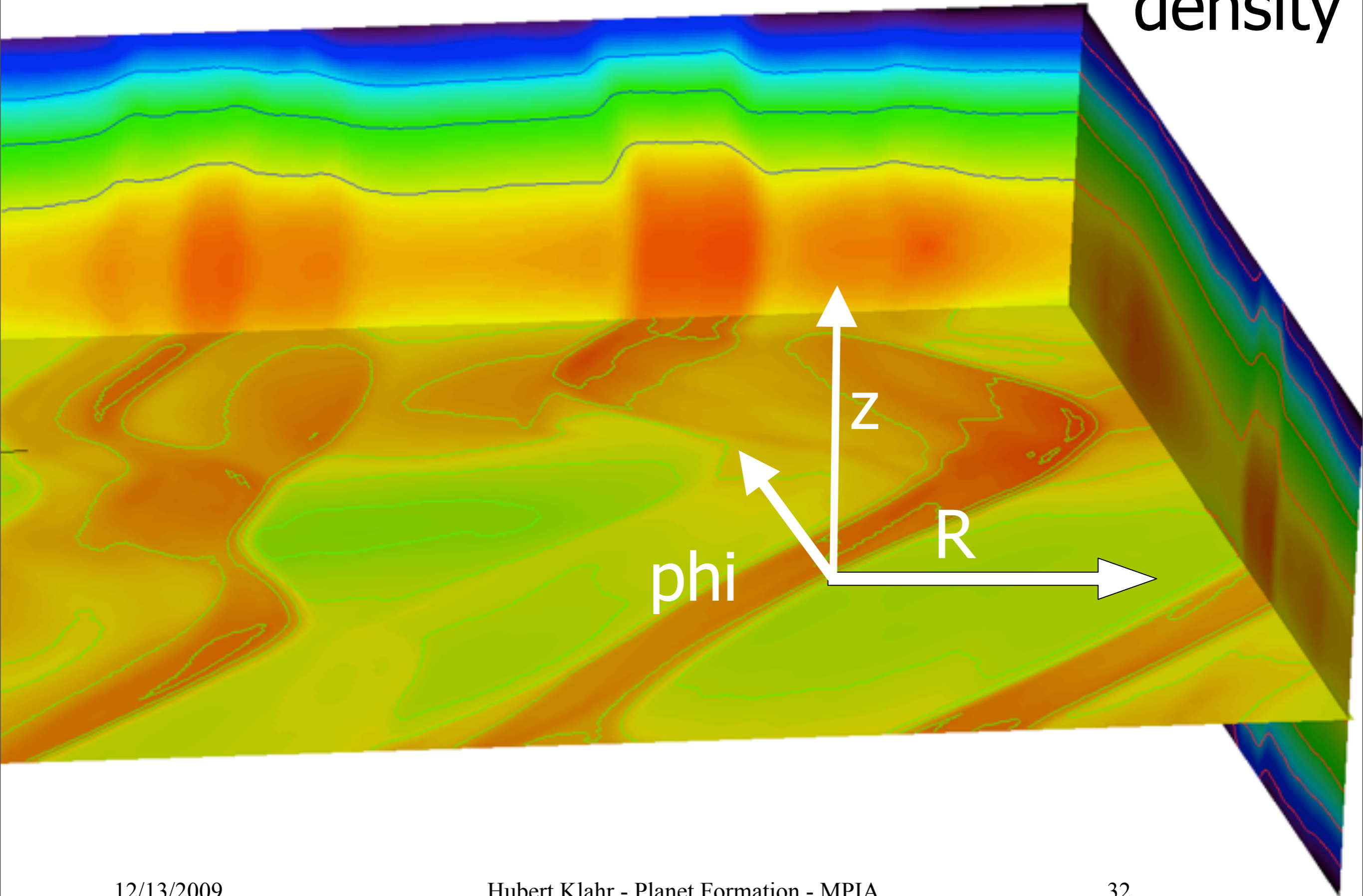




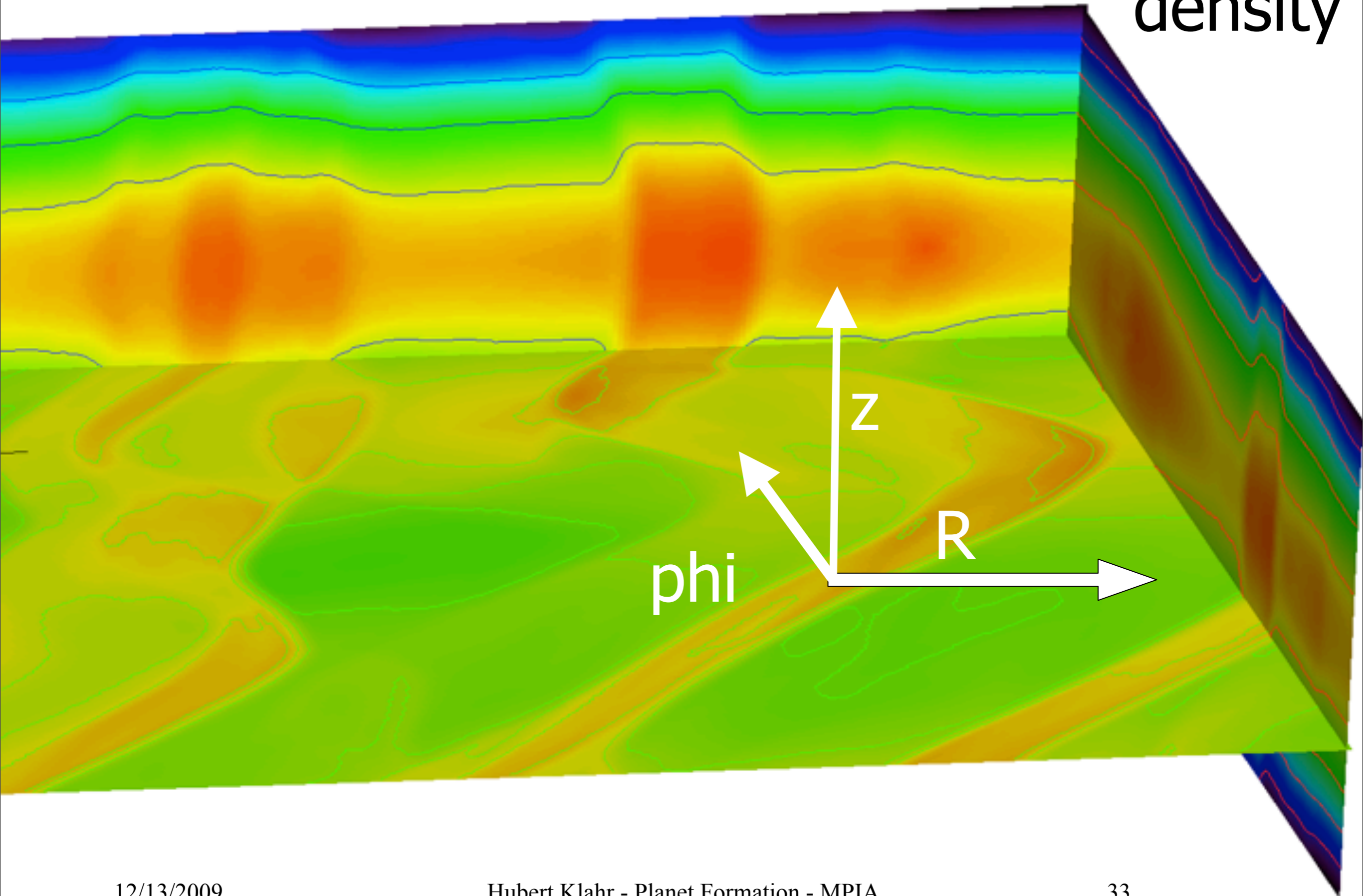
density



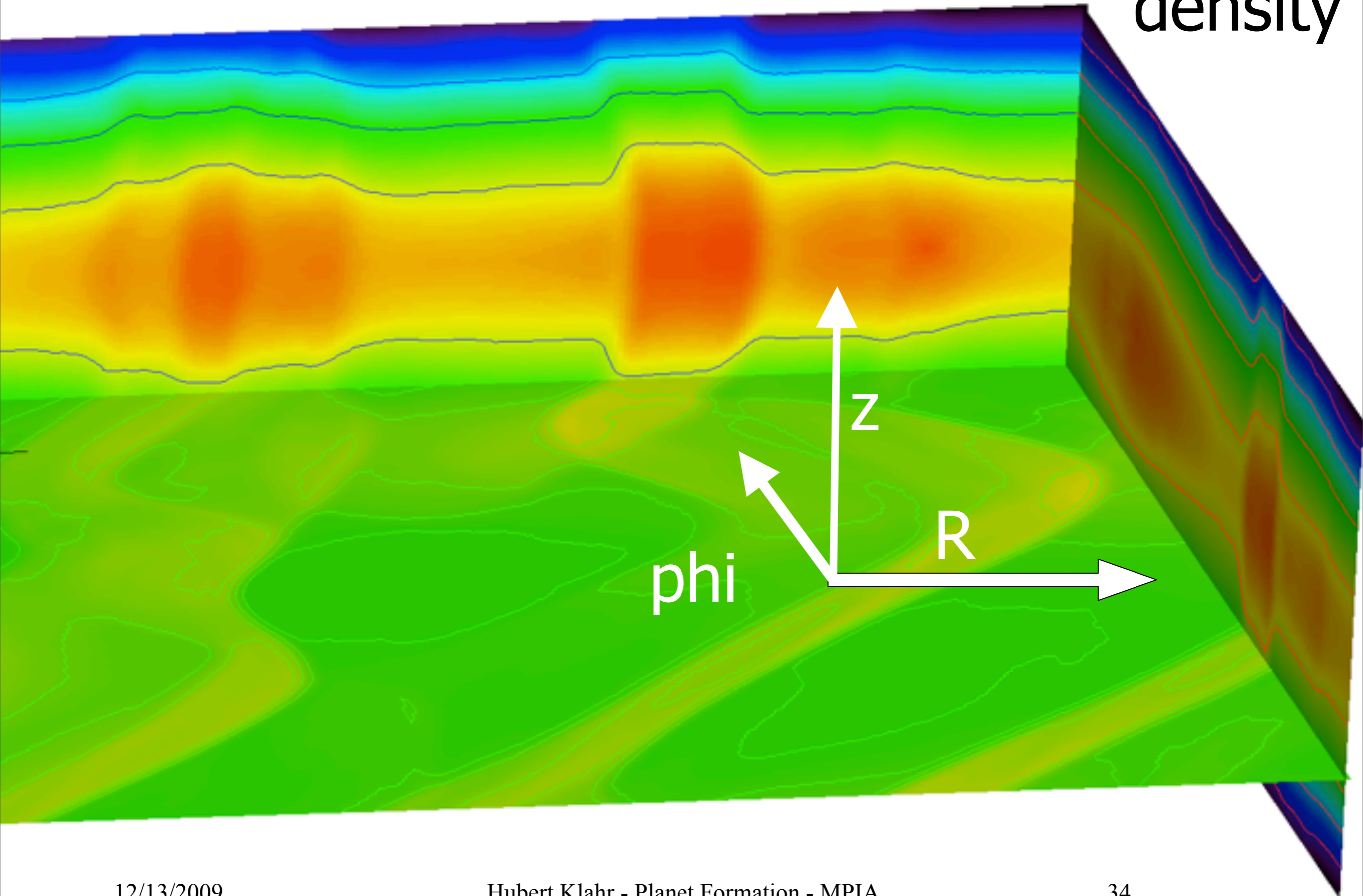
density



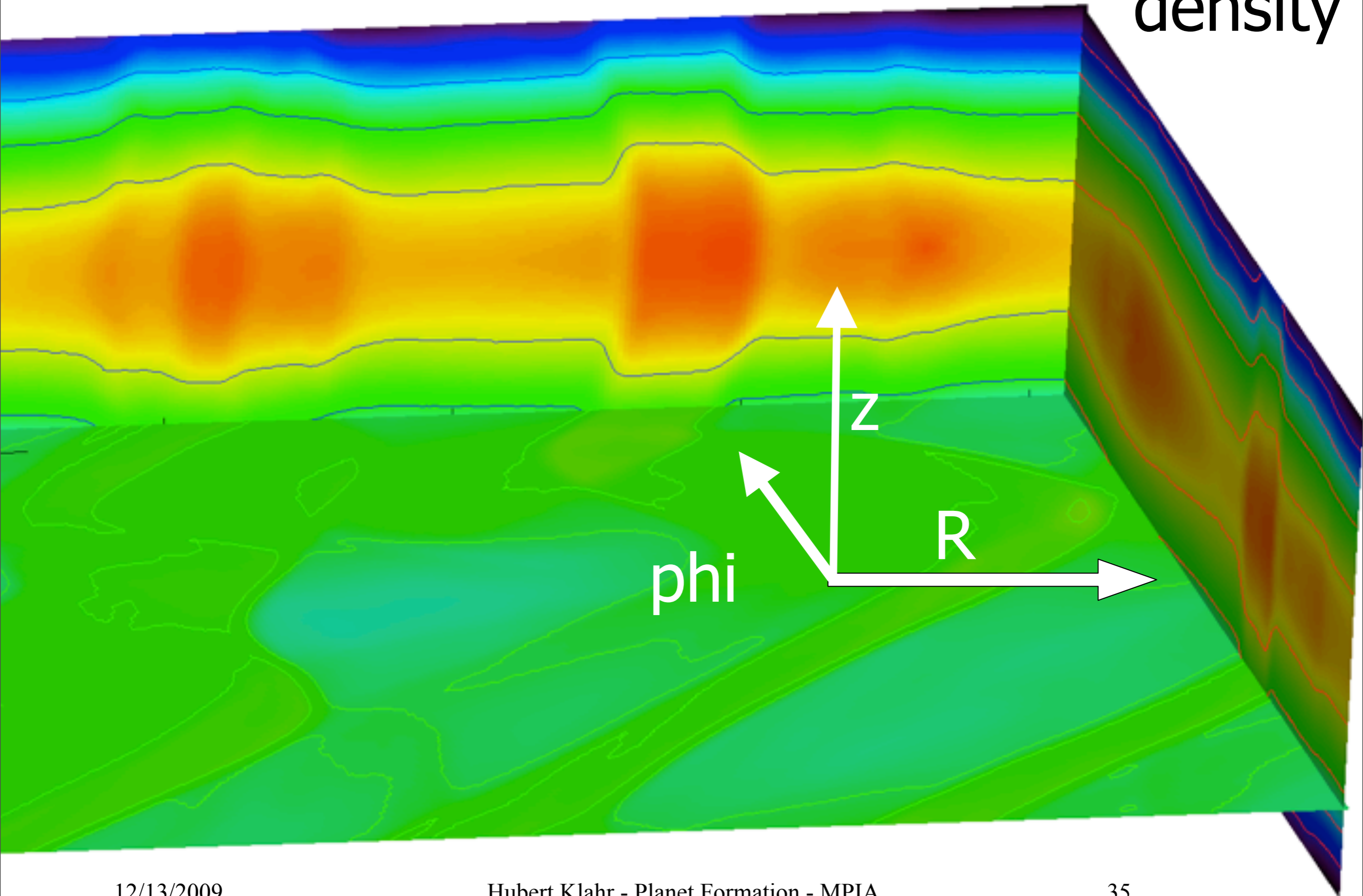
density

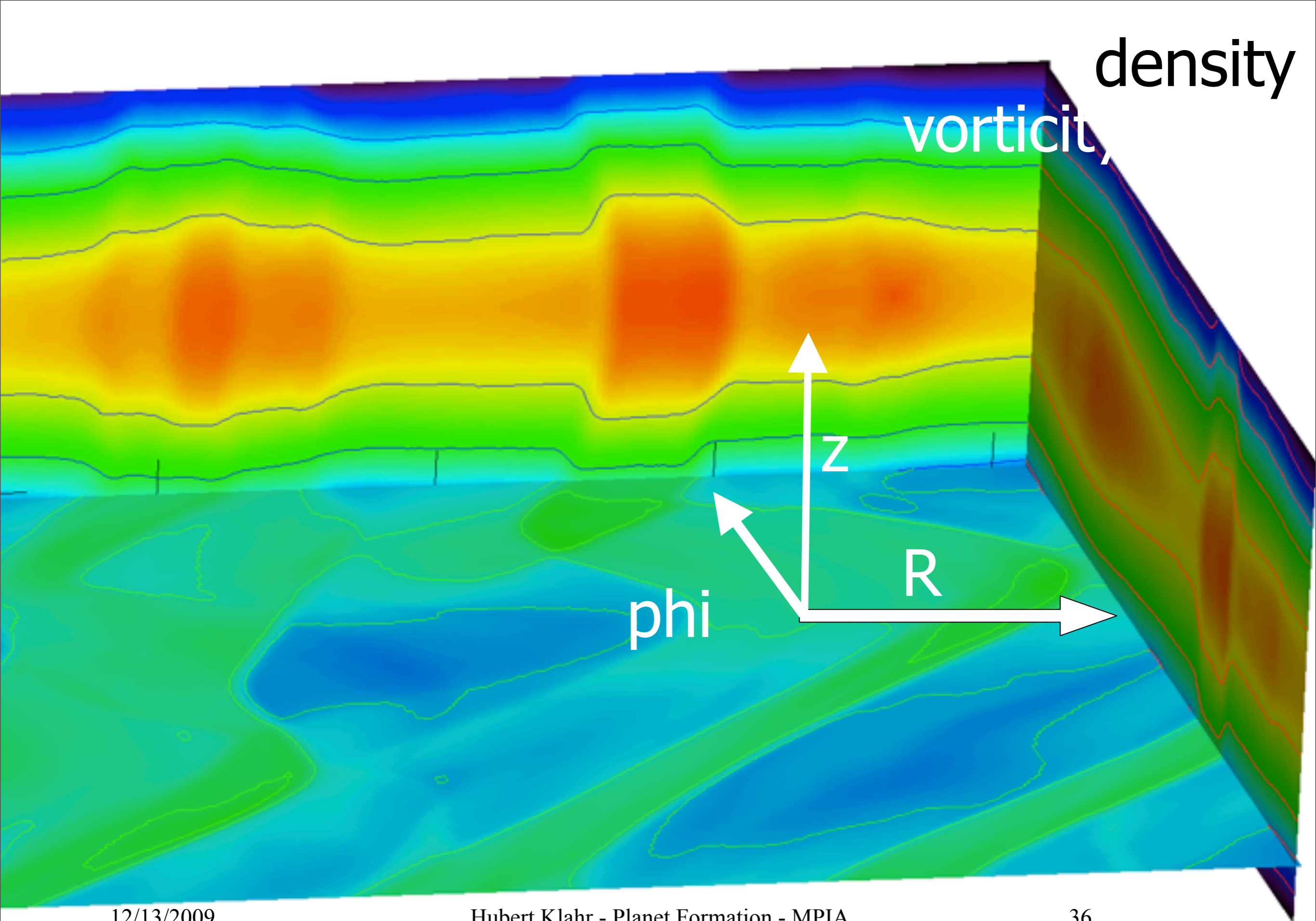


density

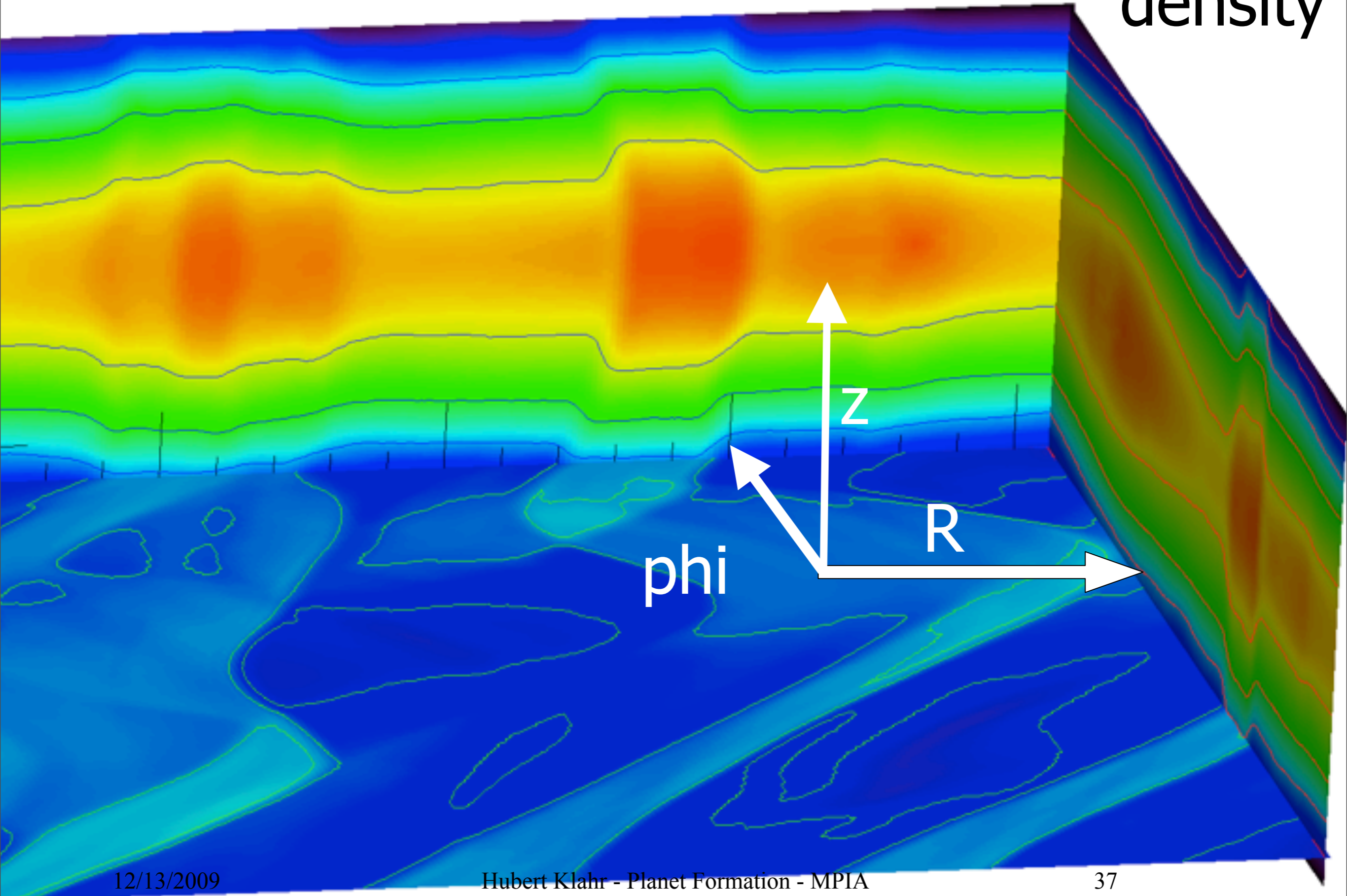


density





density

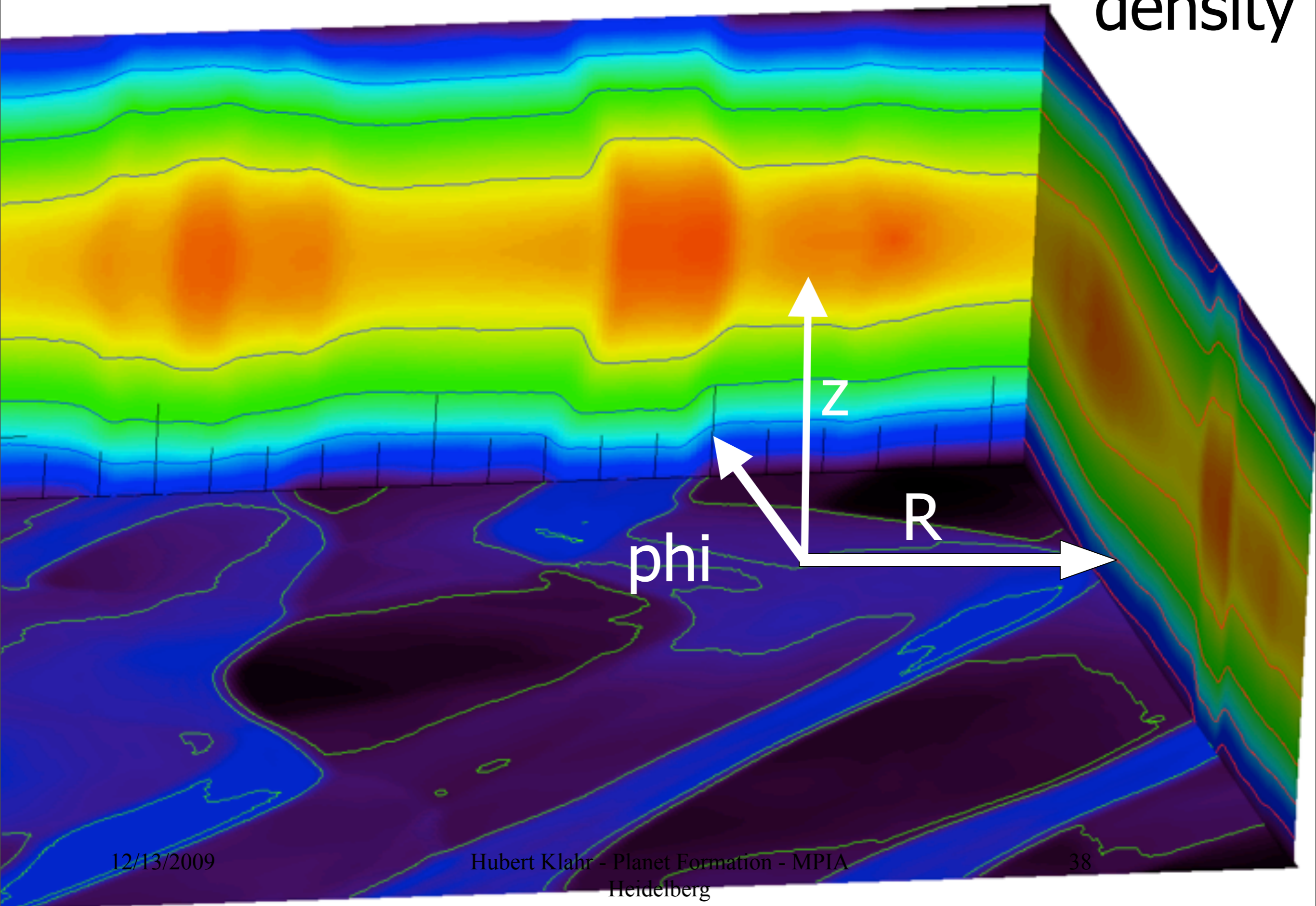


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density



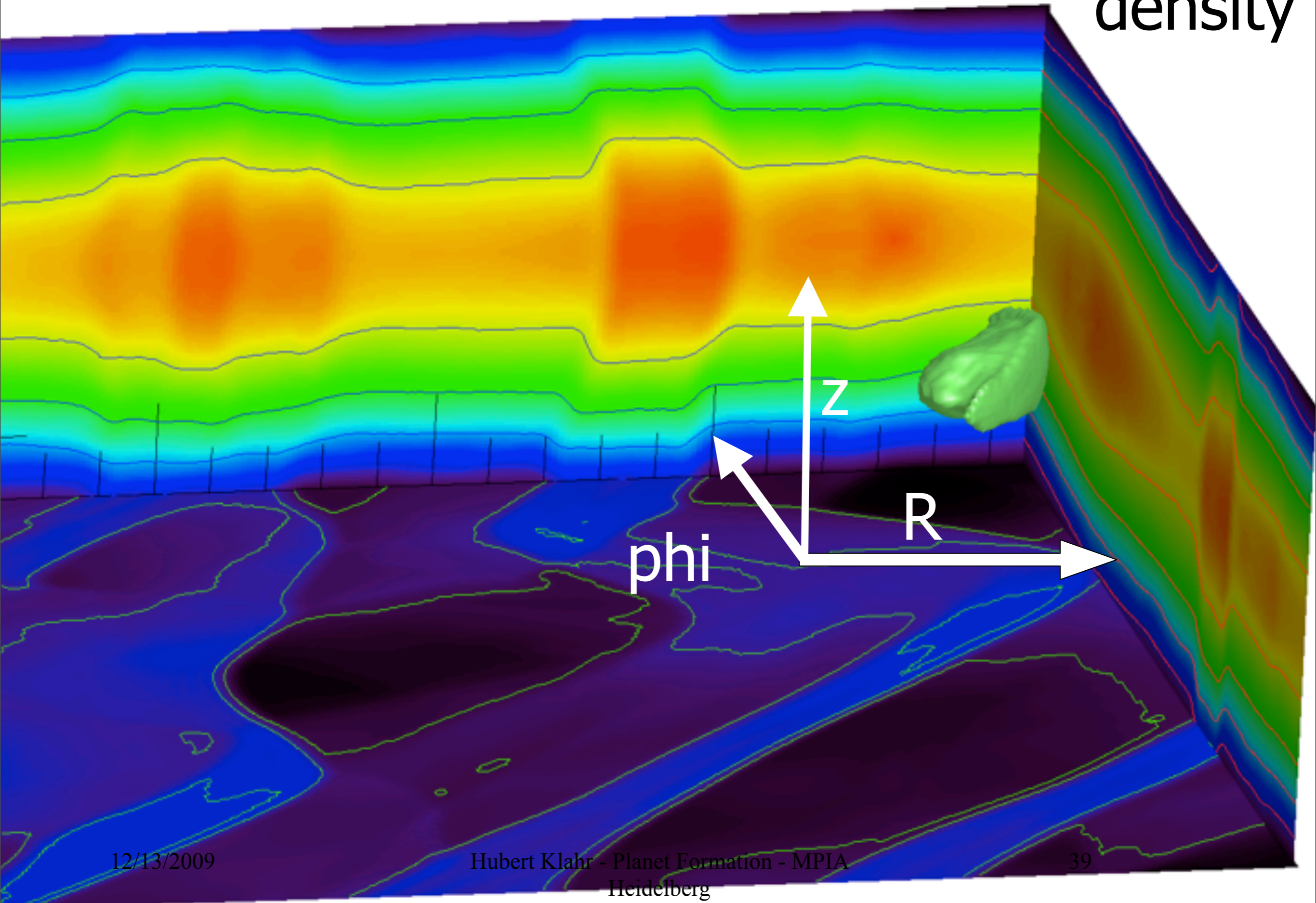
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density

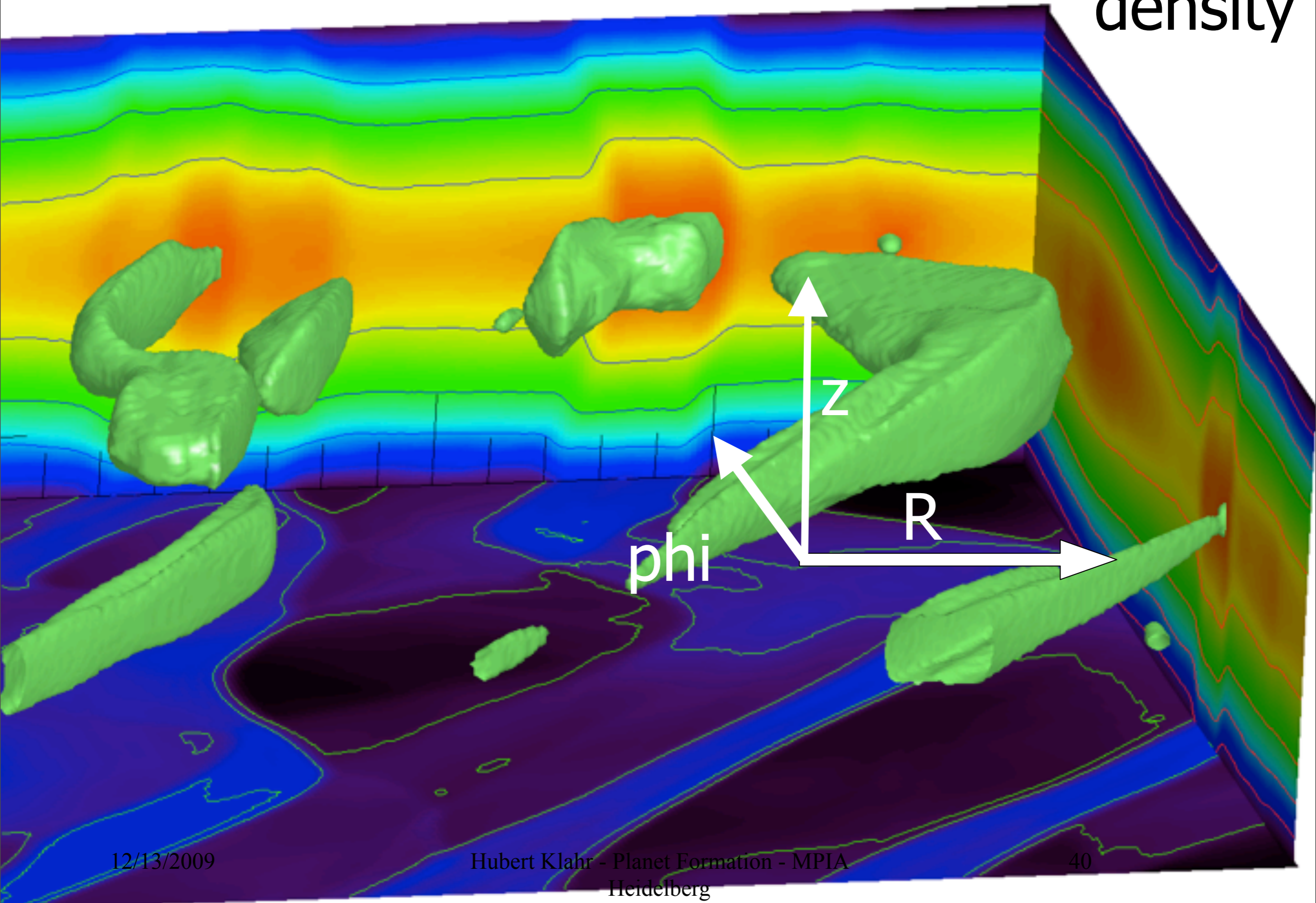


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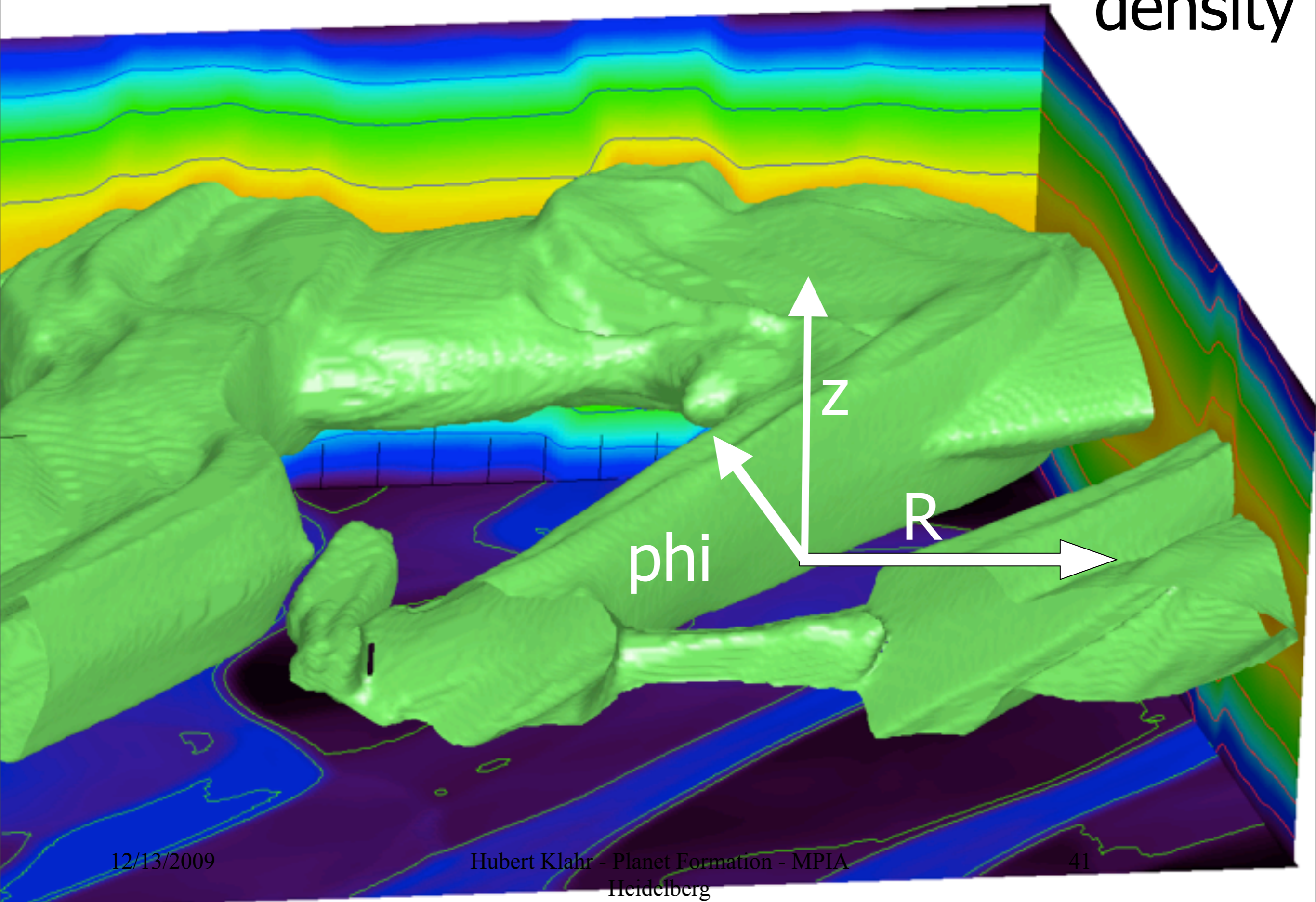


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density

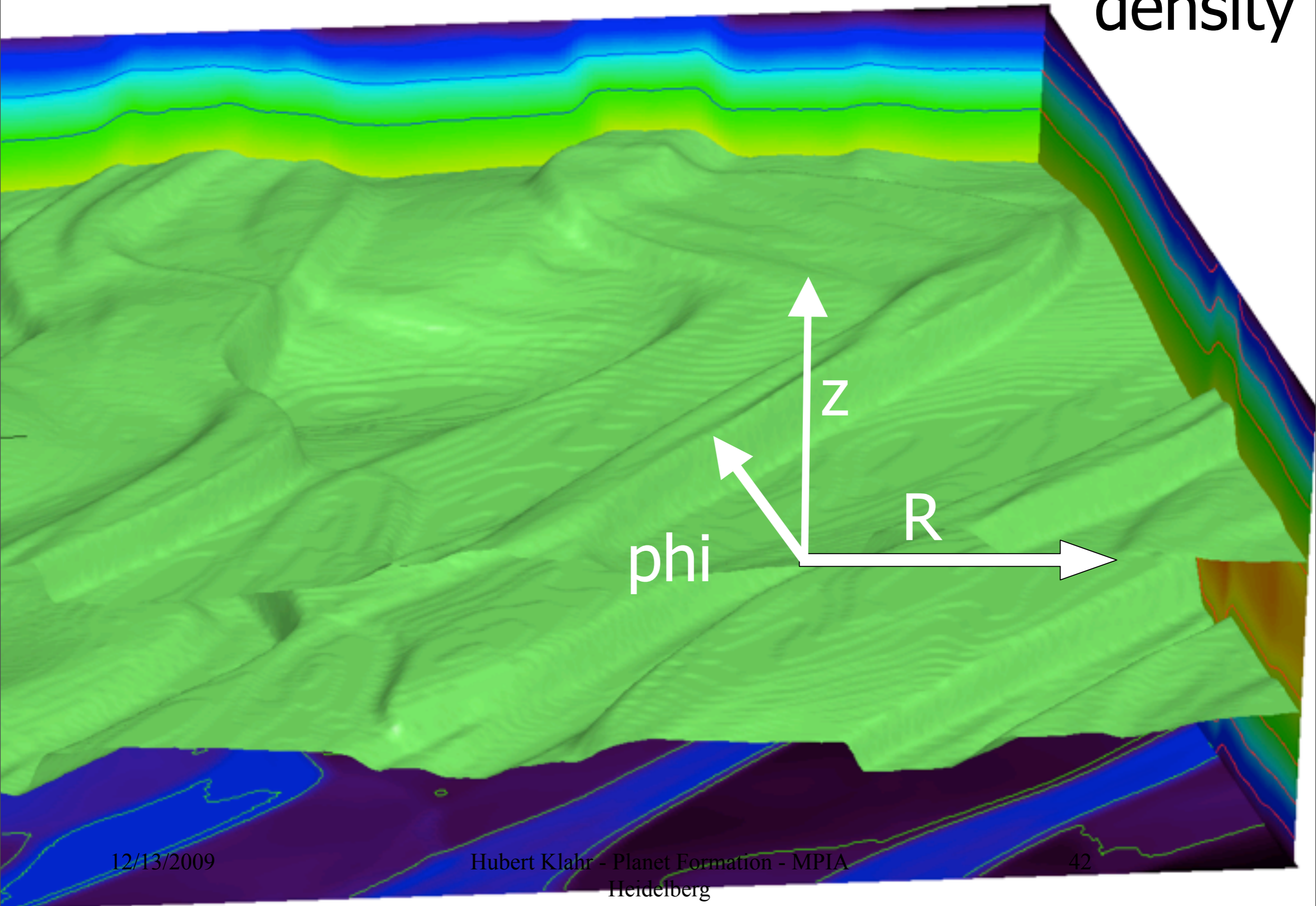


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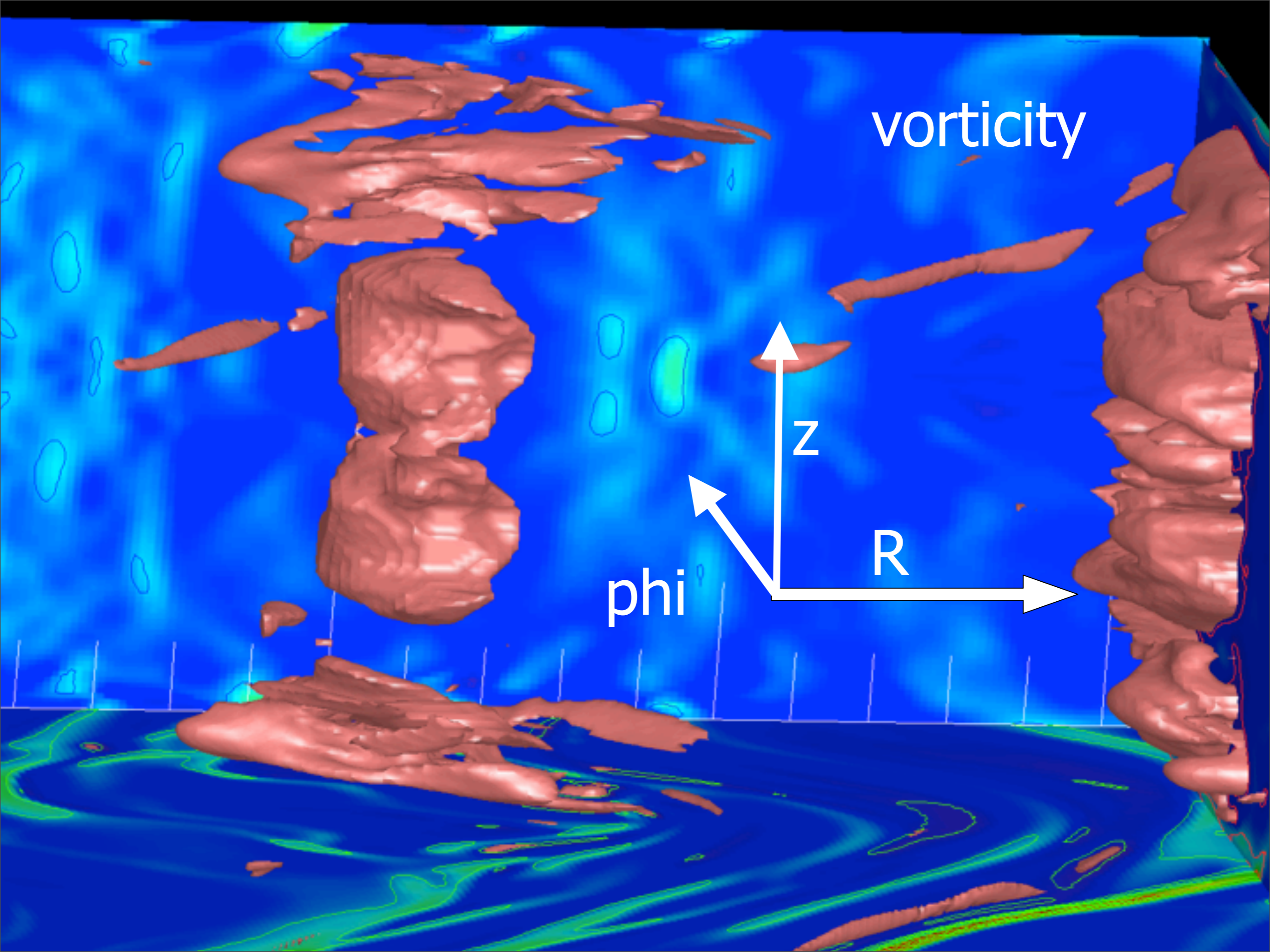
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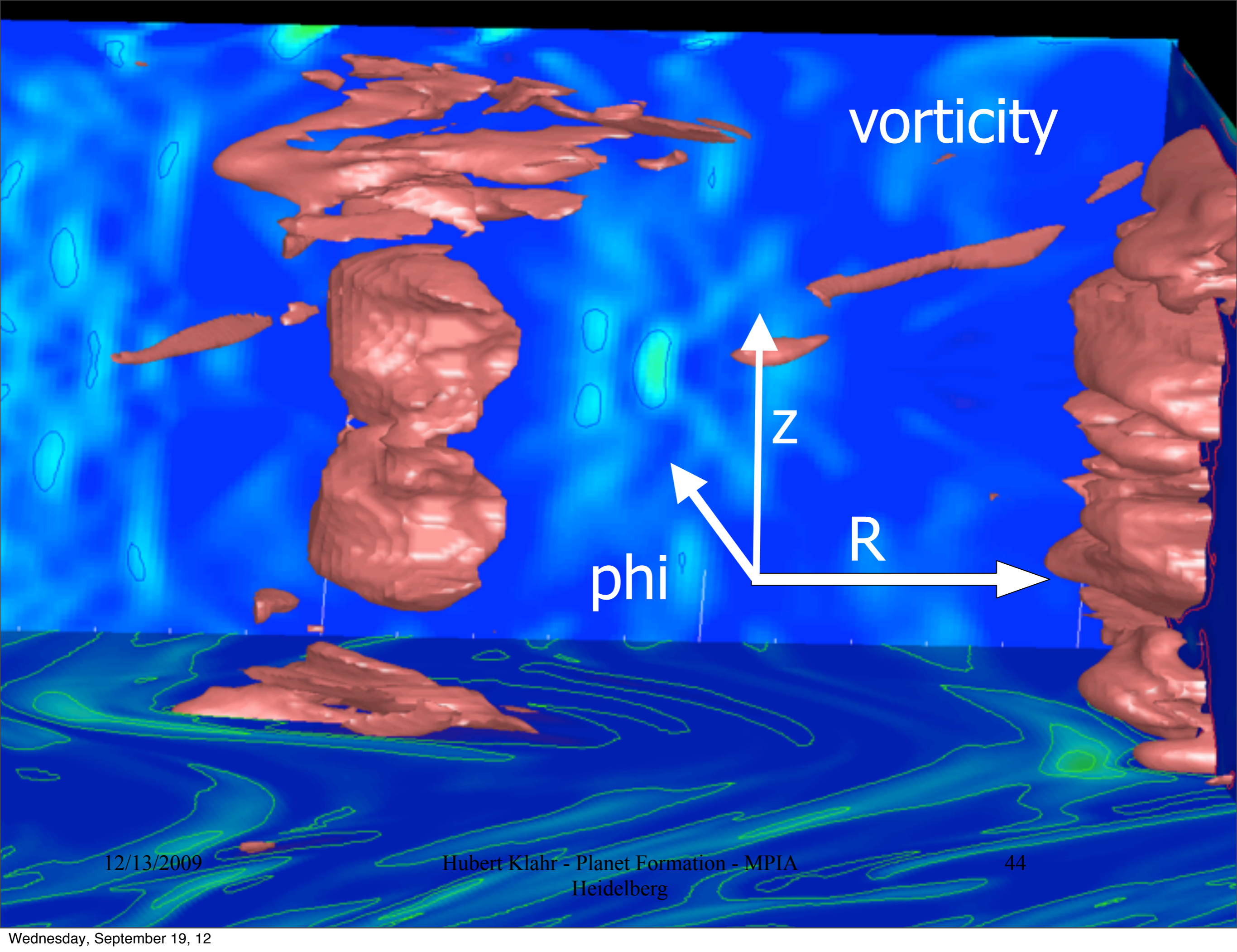


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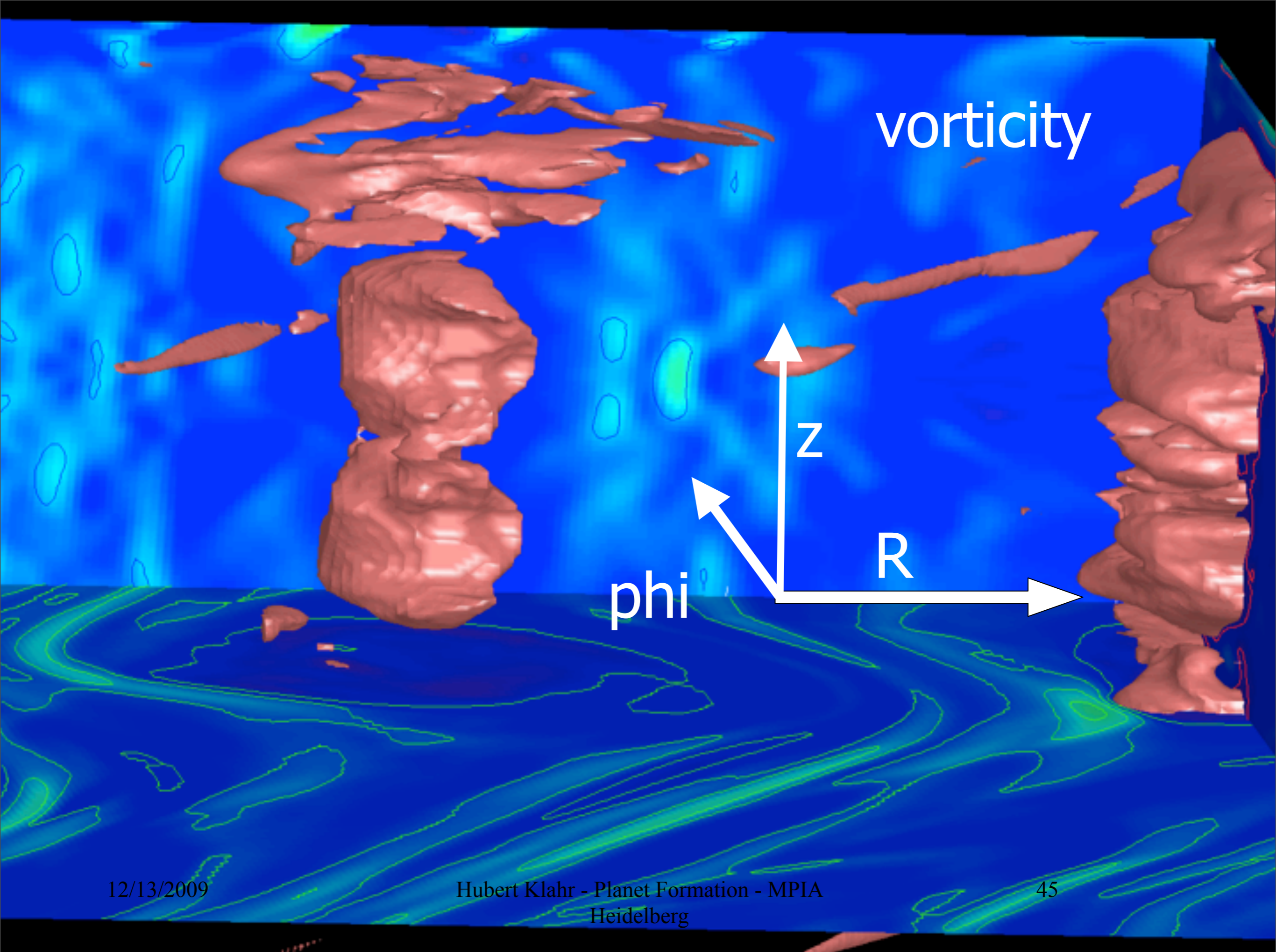


vorticity

phi

z

R

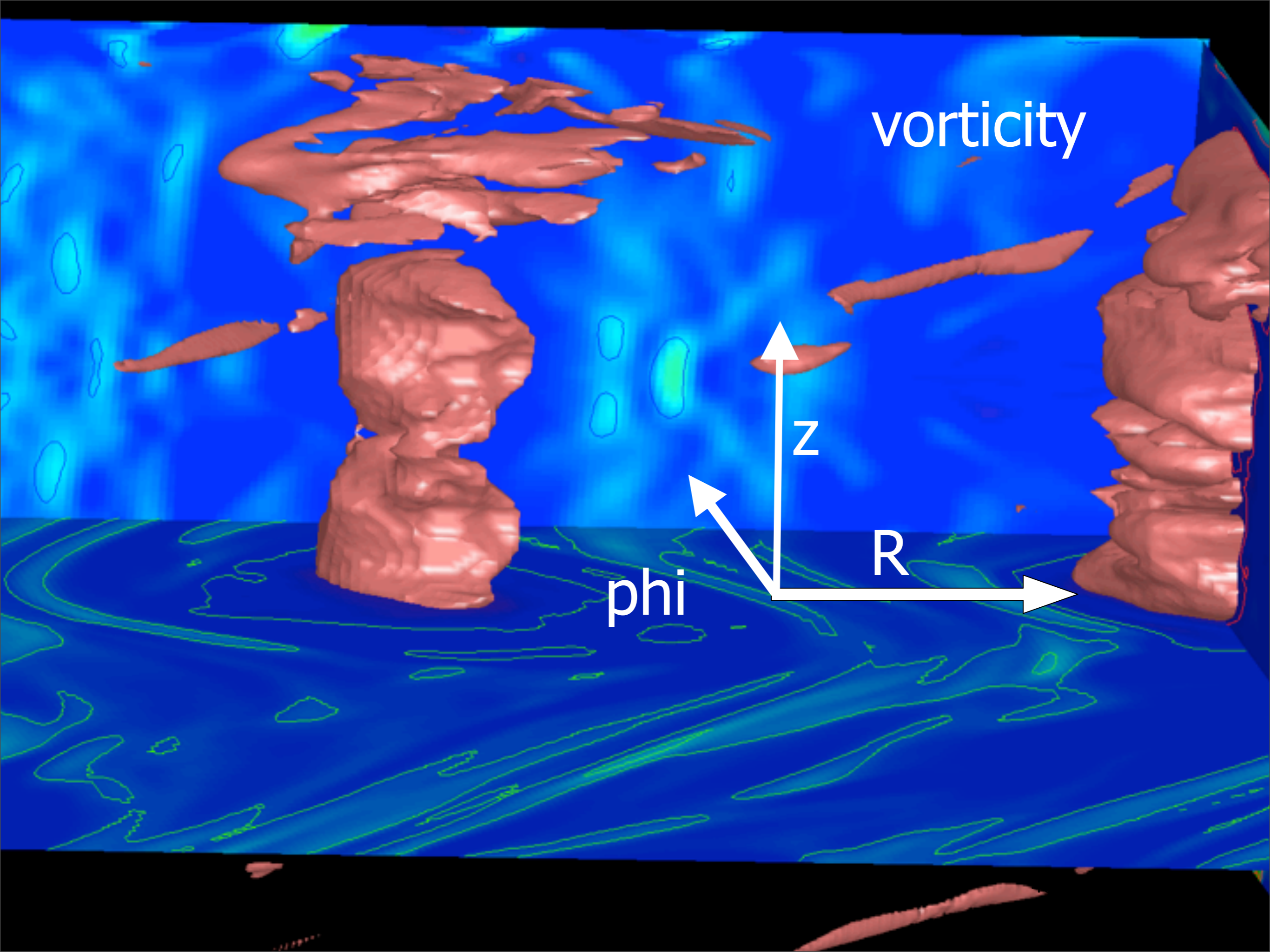


vorticity

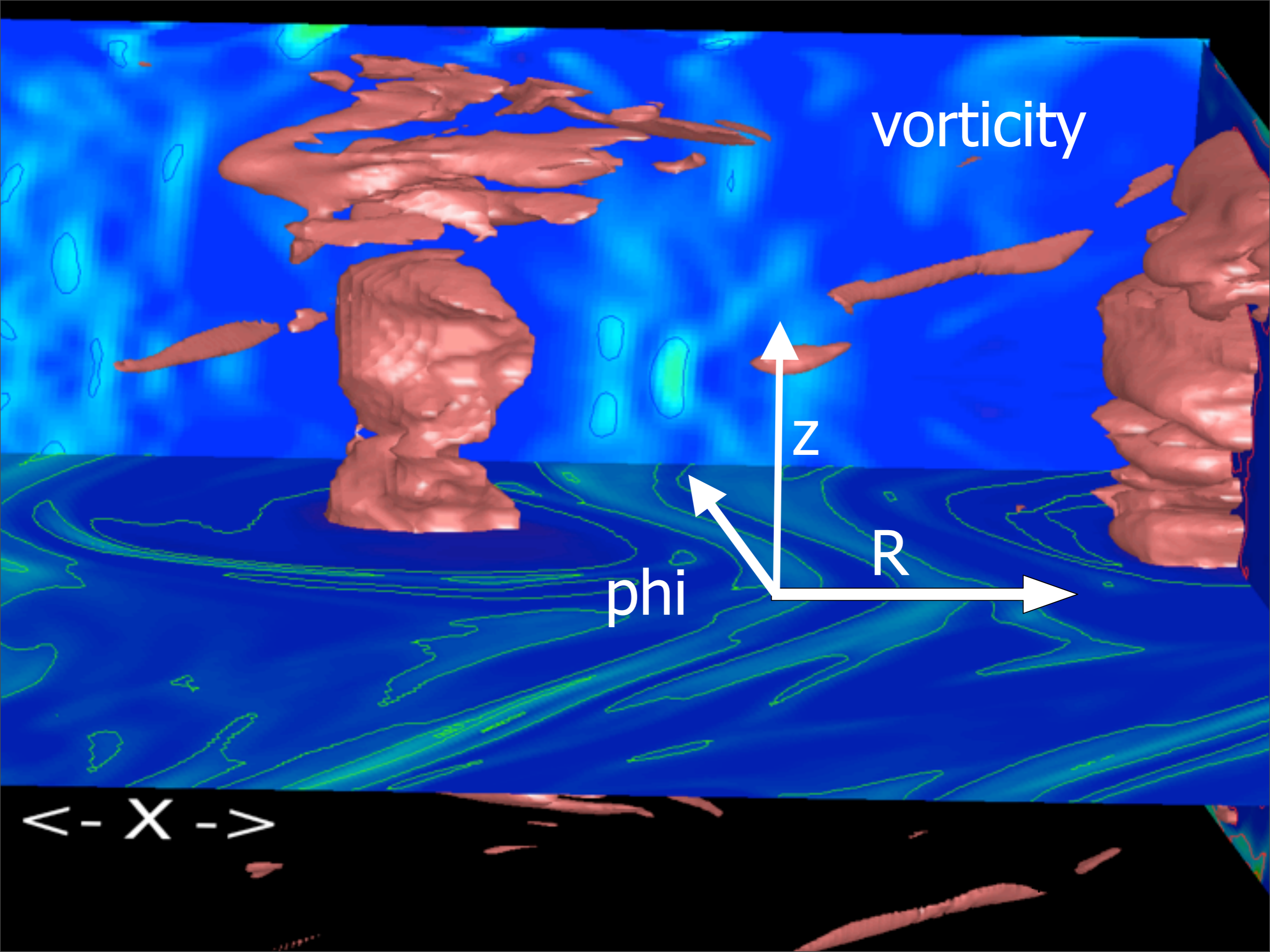
phi

z

R







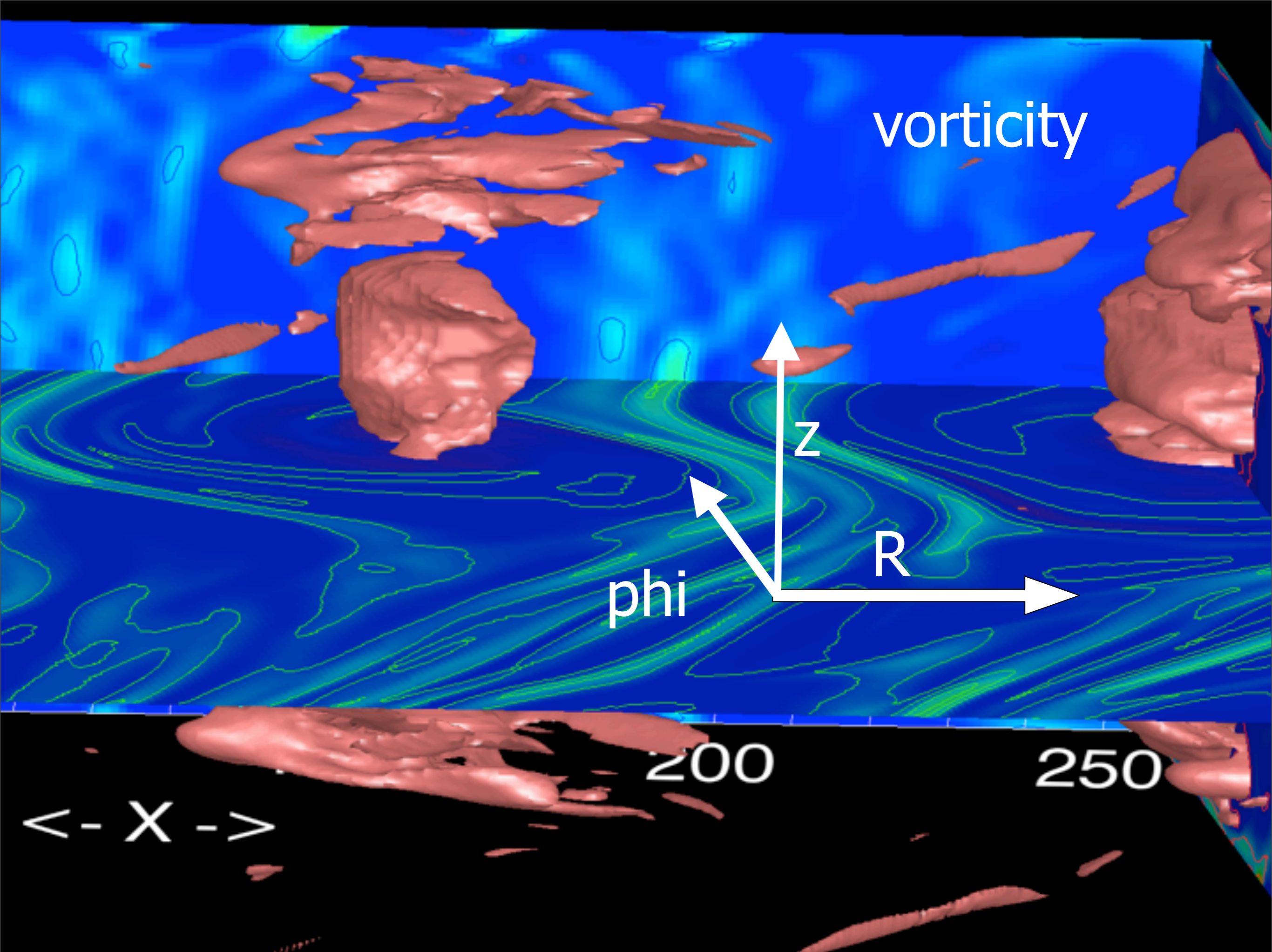
vorticity

$\phi$

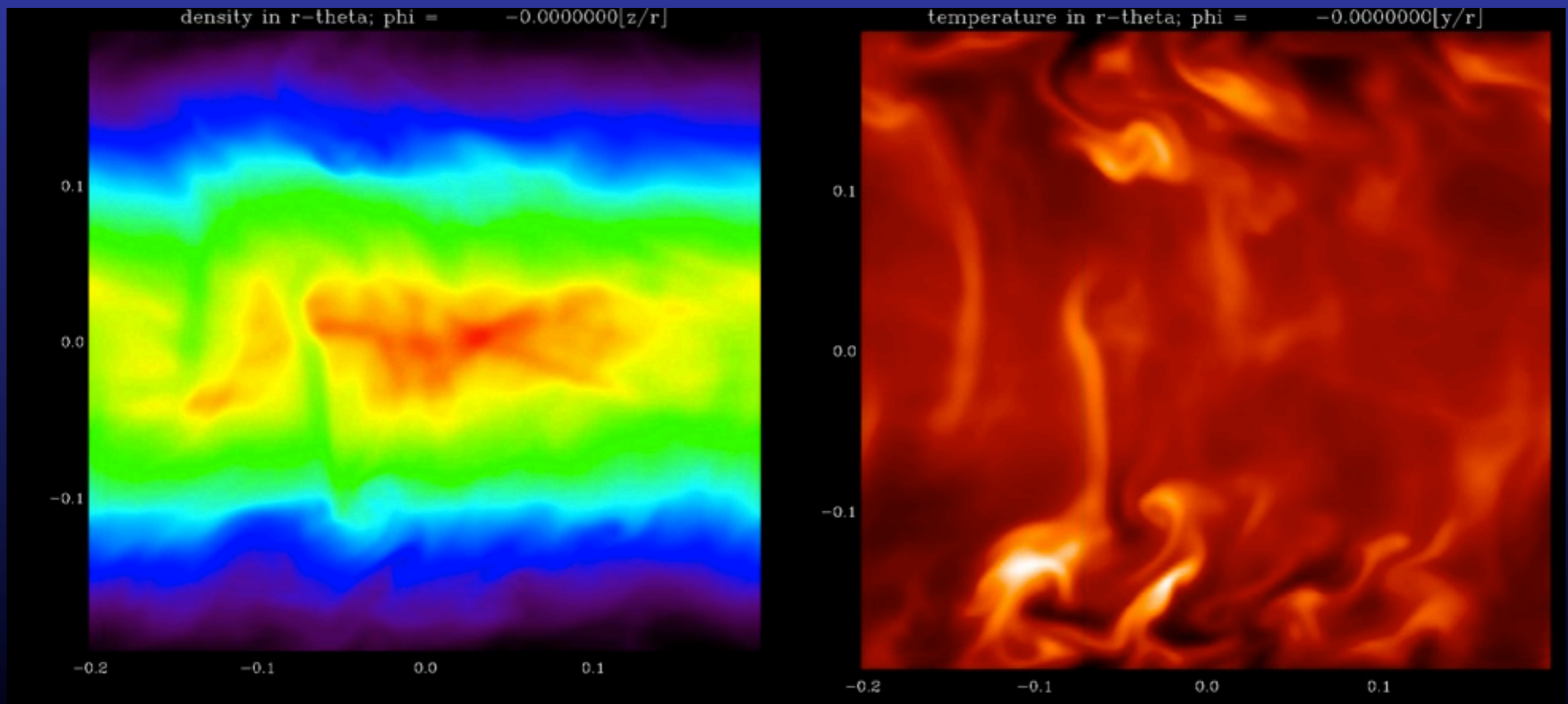
$z$

$R$

$\langle -X - \rangle$



# 2D Local (radial - vertical), Including thermal wind / vertical shear! How Come?



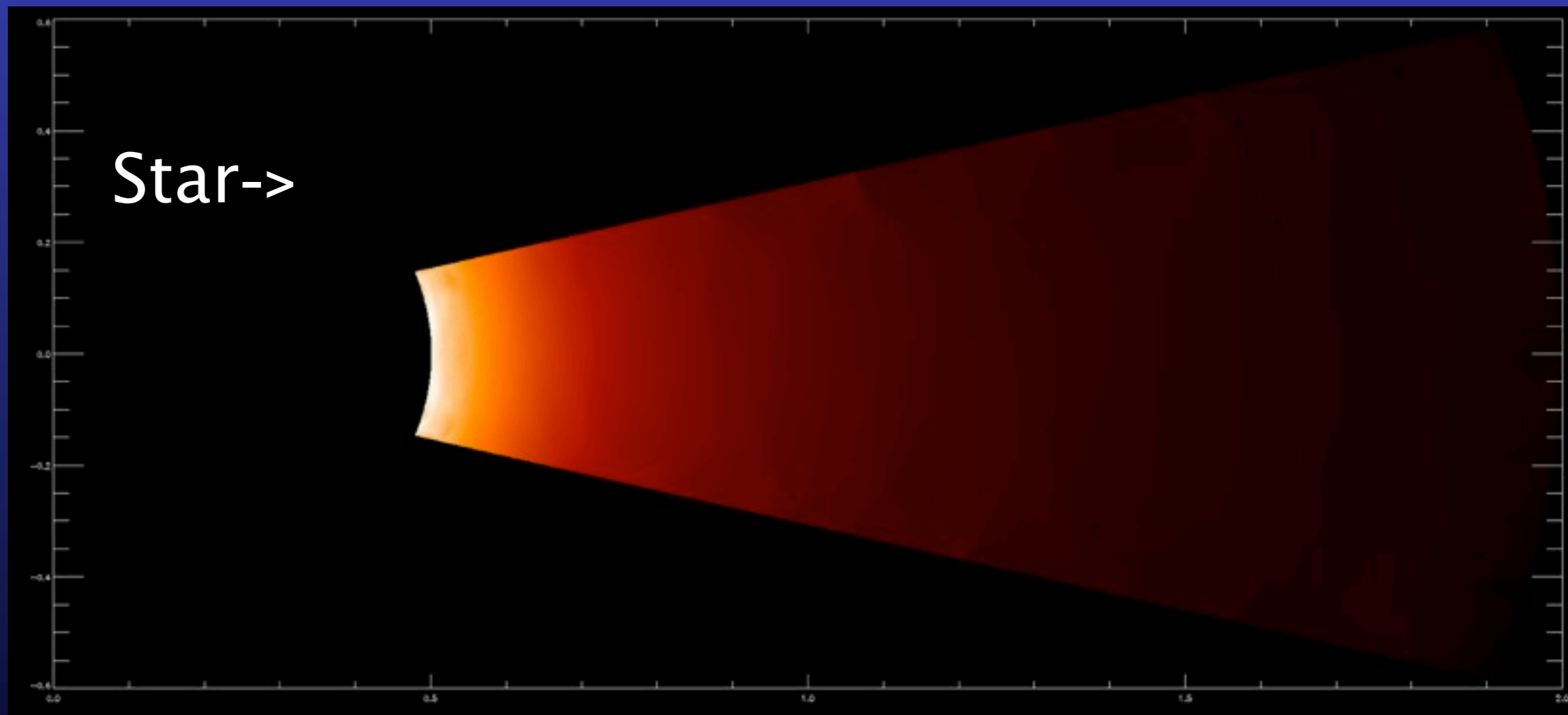
density

temperature pert.

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Tobert Klahr - Planet Form

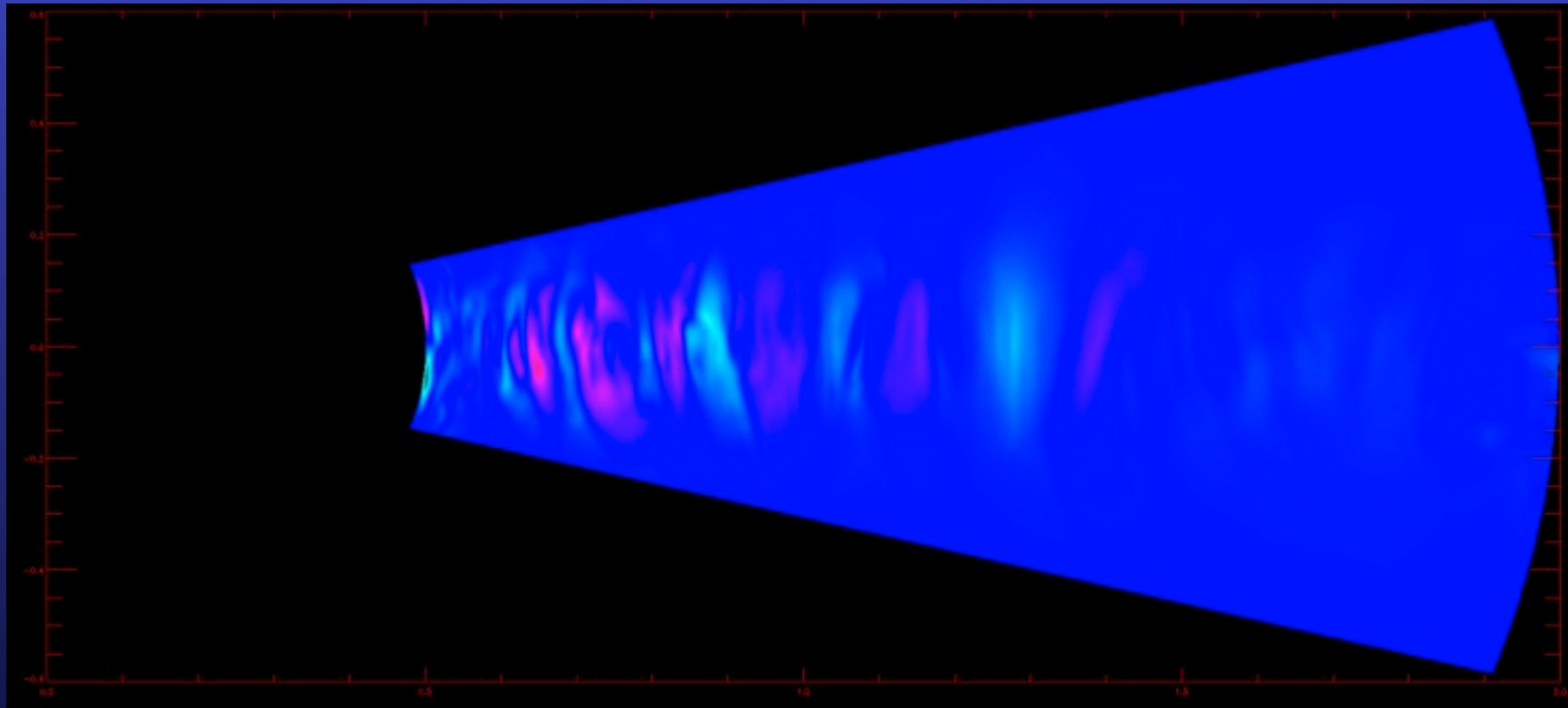
2D axisymmetric Pluto Simulation: Temperature  
due to irradiation from star and thermal  
relaxation  $\tau = 0.1$  (also works for flux limited  
diffusion in irradiated disks)



Thermal wind:

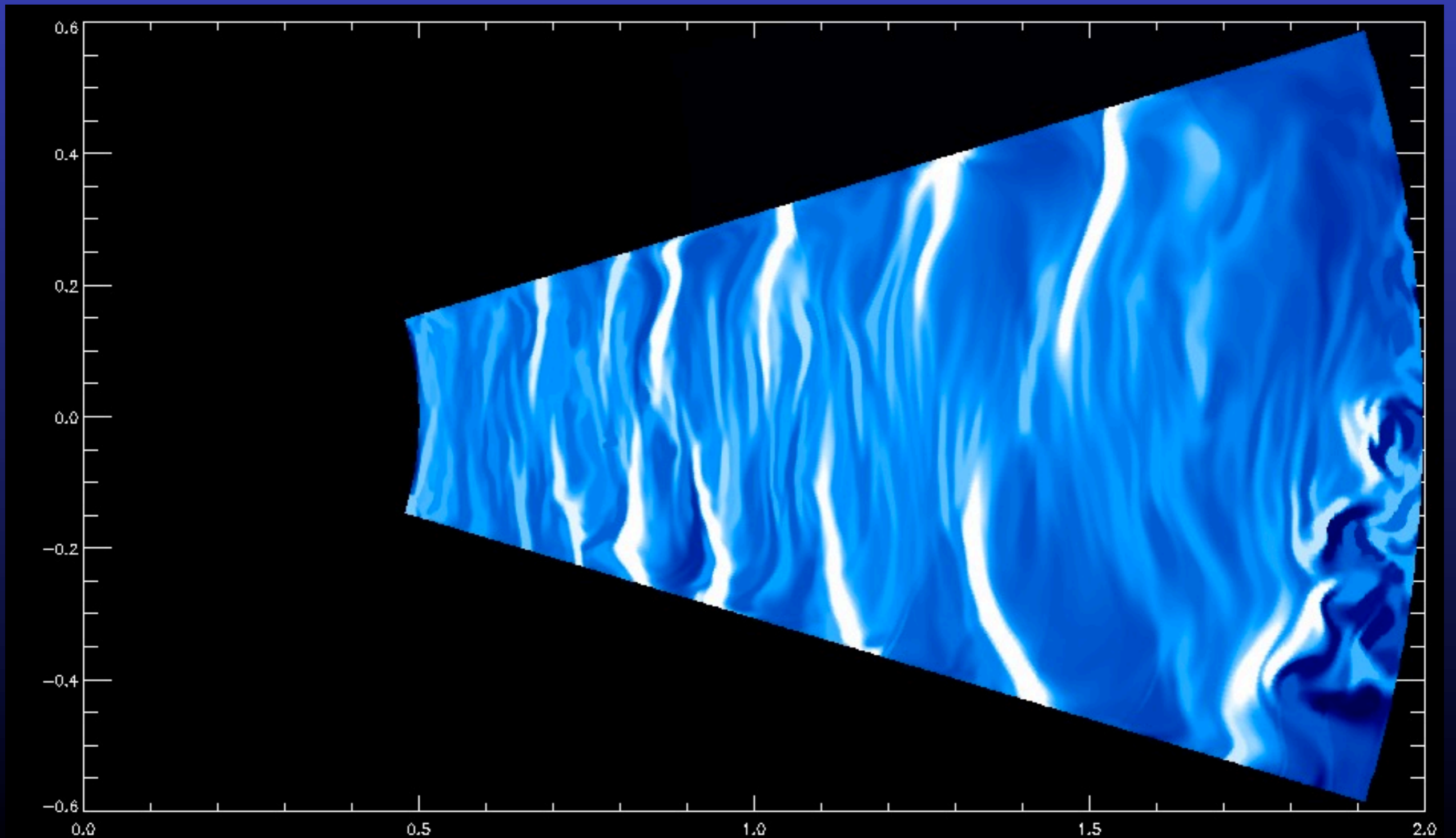
$$\Omega_K \left[ 1 + \frac{1}{2} \left( \frac{H}{R} \right)^2 \left( p + q + \frac{q}{2} \frac{Z^2}{H^2} \right) \right]$$

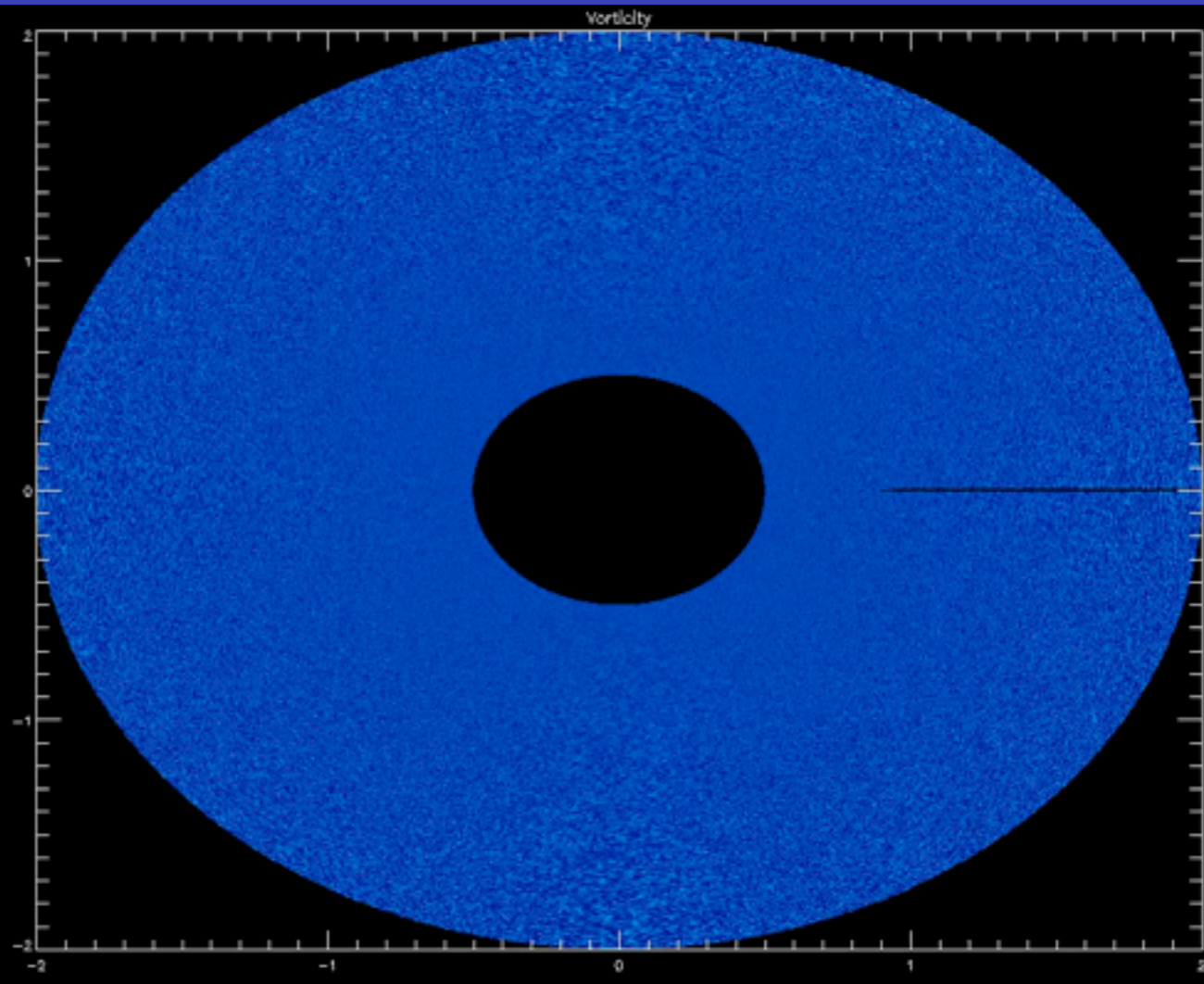
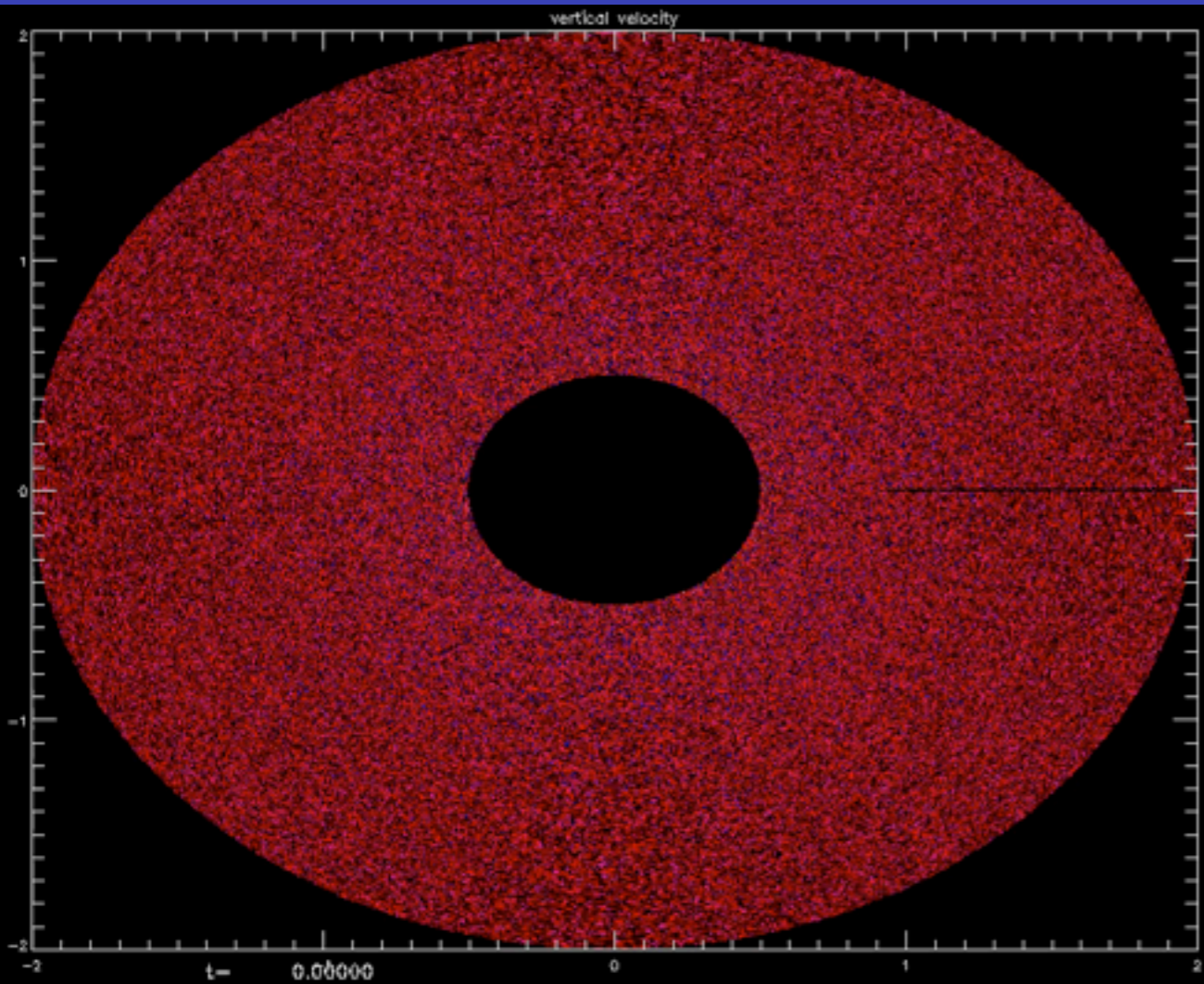
# 2D axisymmetric Pluto Simulation: Overstability due to thermal wind leads to convection like motion: Symmetric Instability

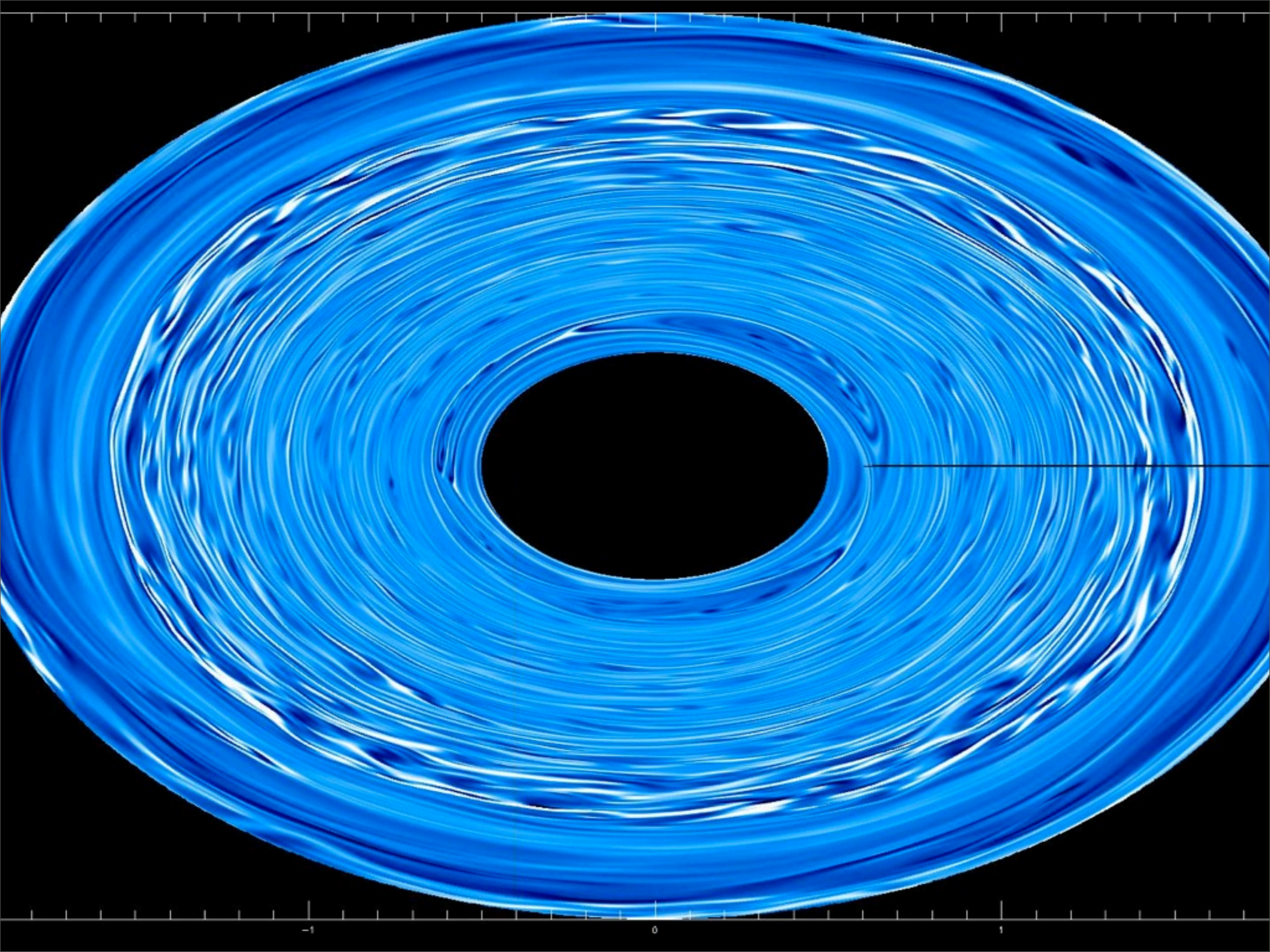


Modification of Solberg-Hoiland Criterion, including thermal relaxation:  
In collaboration with Alexander Hubbard

# 2D axisymmetric Pluto Simulation: Resulting vorticity perturbations





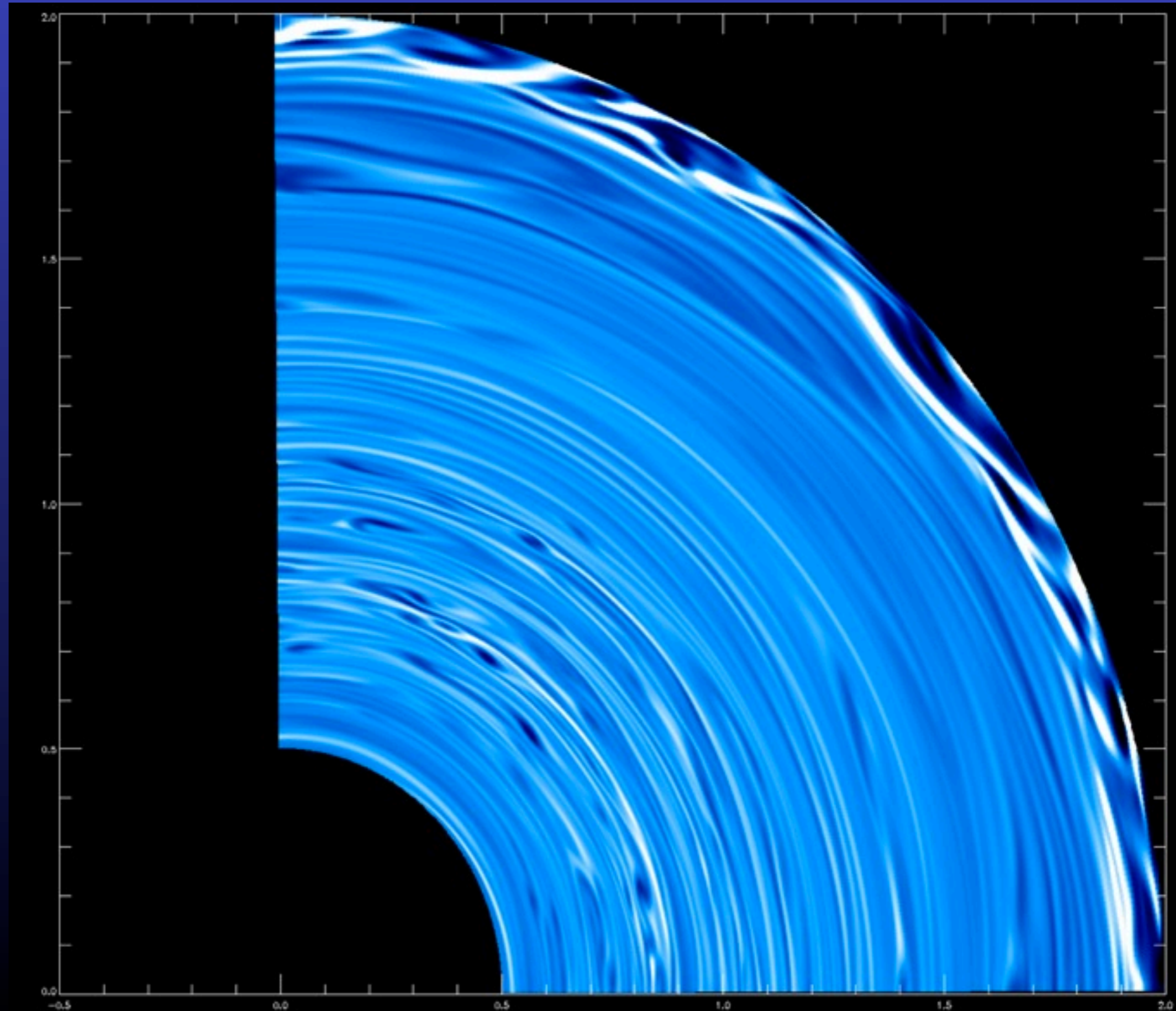




# Full 3D Pluto Simulation: Spontaneous Formation of Vortices from tiny perturbations

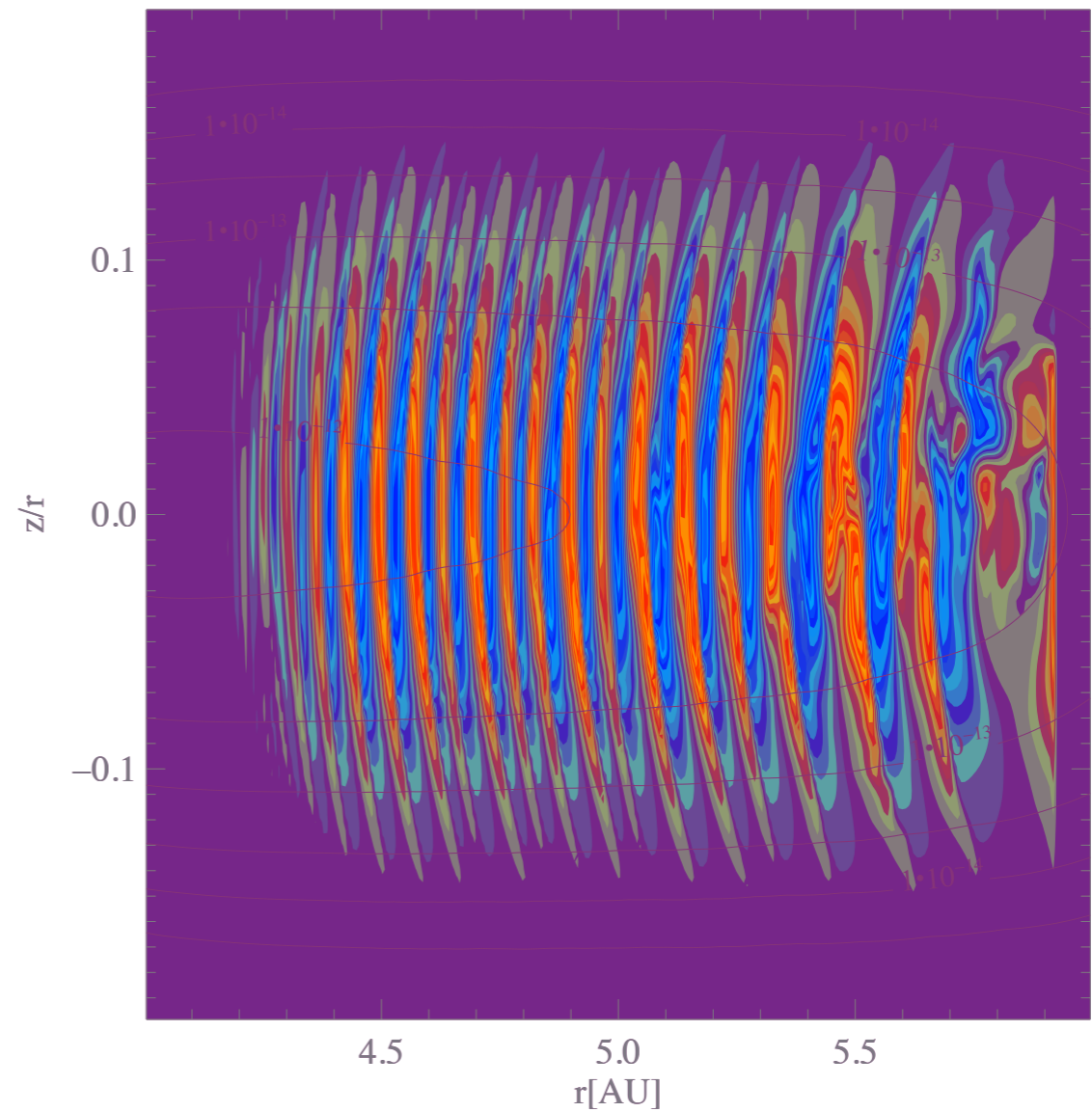
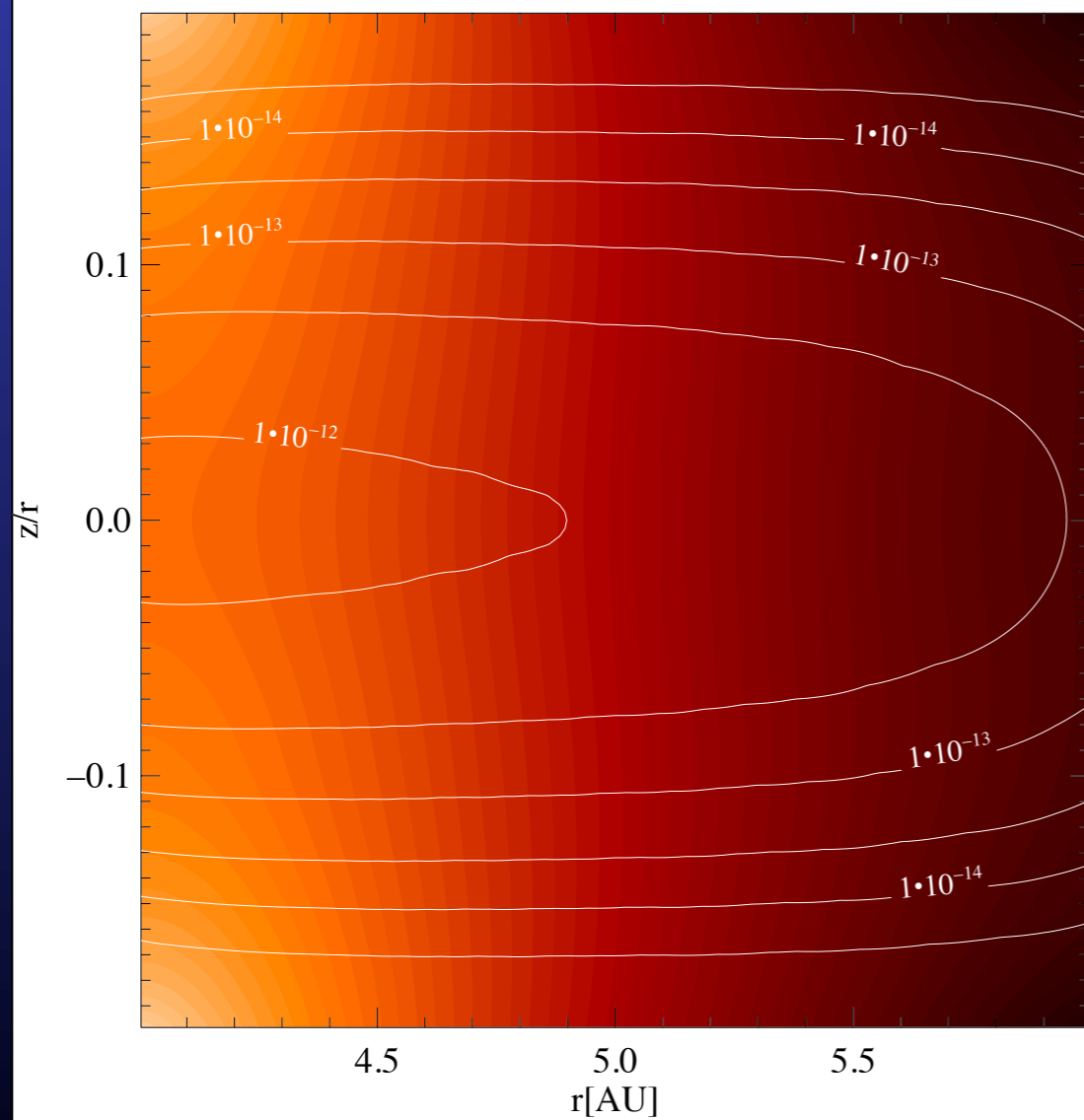
=> Baroclinic Disks  
Form Vortices  
abundantly and  
further amplify  
them further!

Mission  
accomplished.



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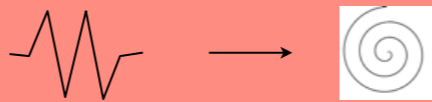
# 2D / 3D - Radiation Hydro of Irradiated disks in Disks with dust opacities and realistic parameters!



# Disk Weather

Interface  
MRI - Dead Zone

Papaloizou-Pringle  
Rossby Wave Inst.

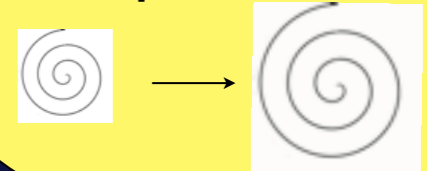


linear  
axisymmetric  
Vertical Shear Inst.:  
Solberg-Hoiland  
See Nelson, Gressel & Umurhan

add.: "Solberg-Hoiland  
Violations" for  $Nz^2 = 0$

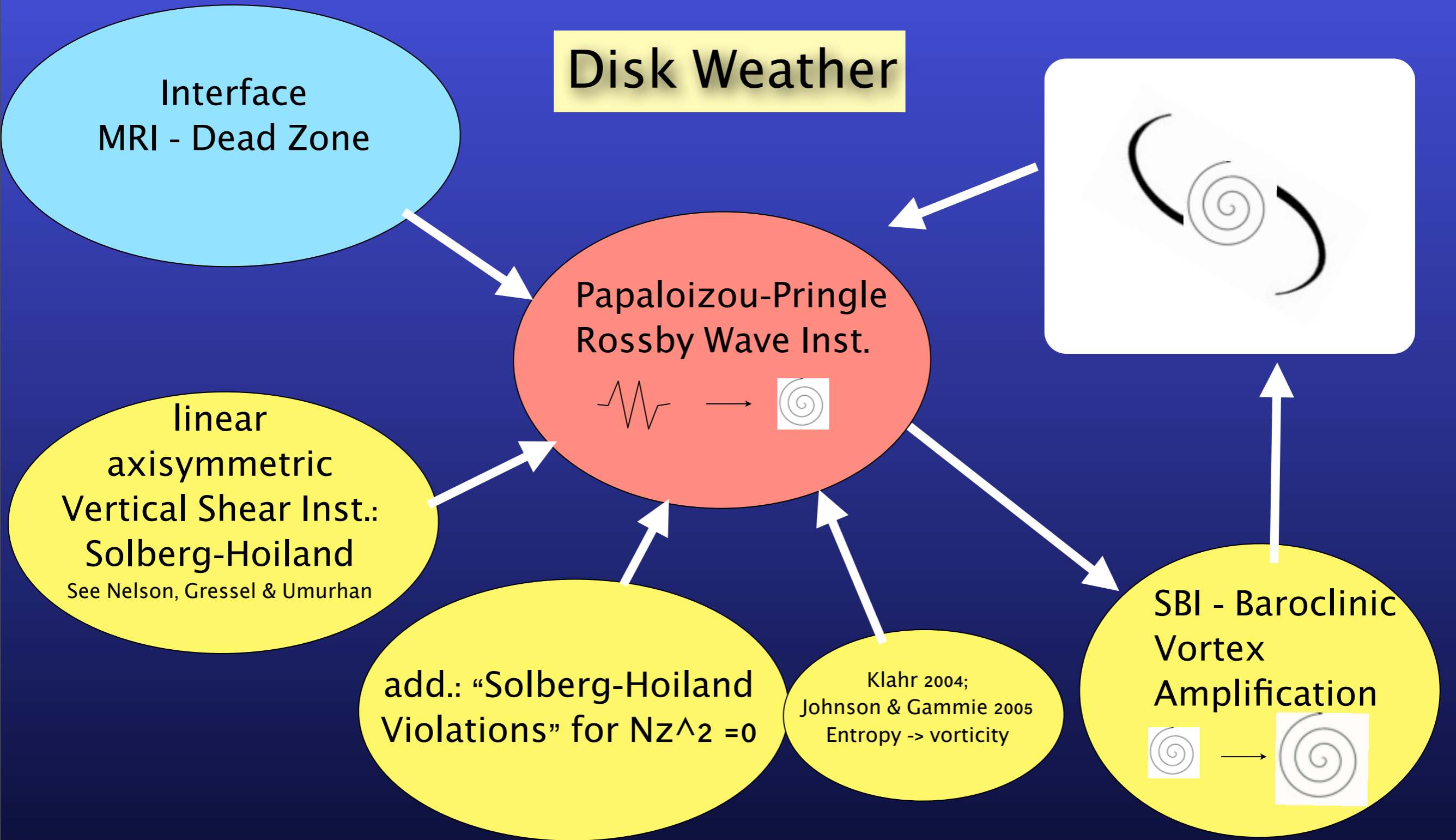
Klahr 2004;  
Johnson & Gammie 2005  
Entropy  $\rightarrow$  vorticity

SBI - Baroclinic  
Vortex  
Amplification



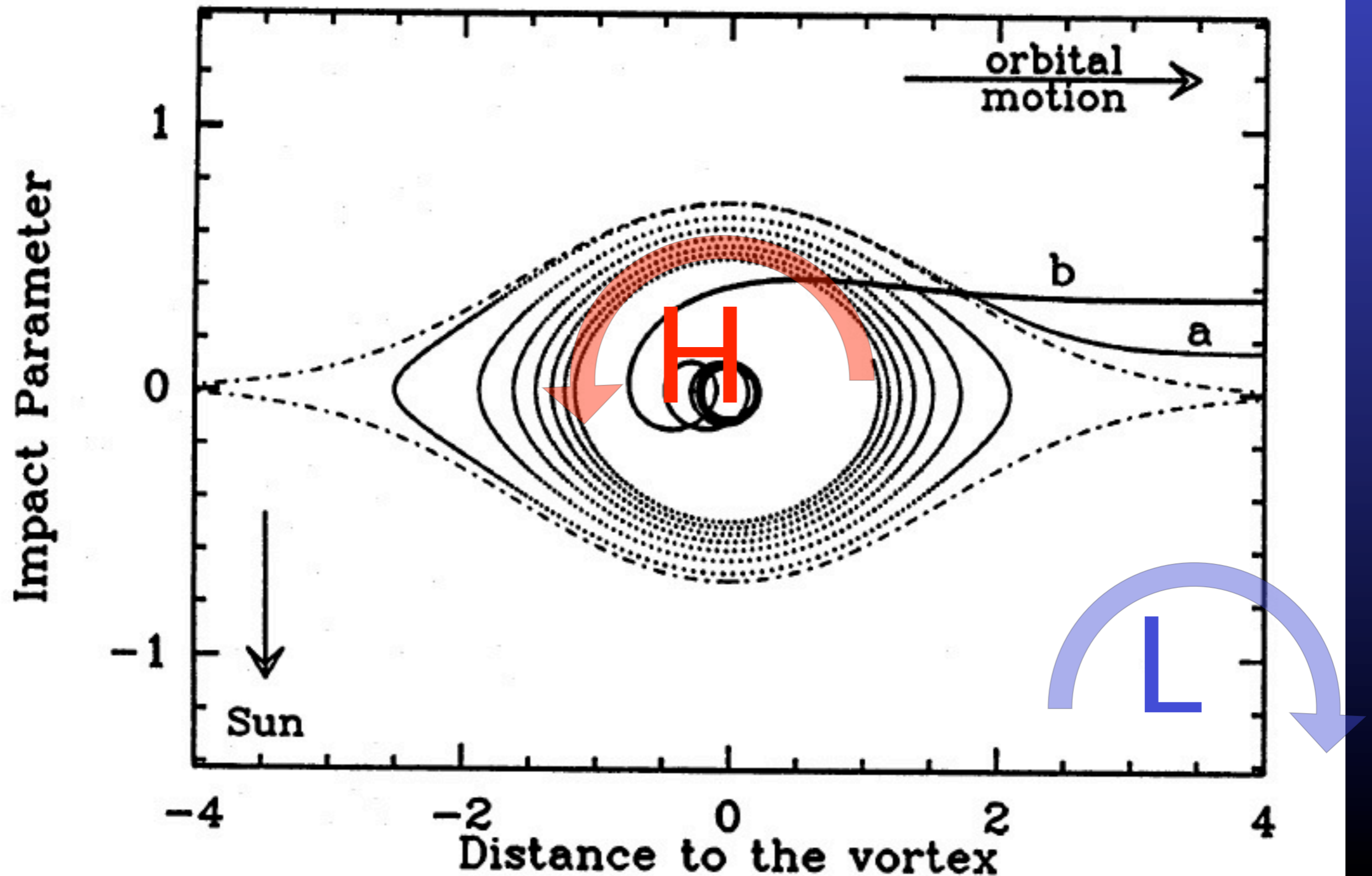
Radial Buoyancy :  $dS/dr < 0$  &  $dP/dr < 0$   
 $\rightarrow$  thermal wind  $dV_{\phi}/dz < 0$       plus: thermal relaxation

# Disk Weather

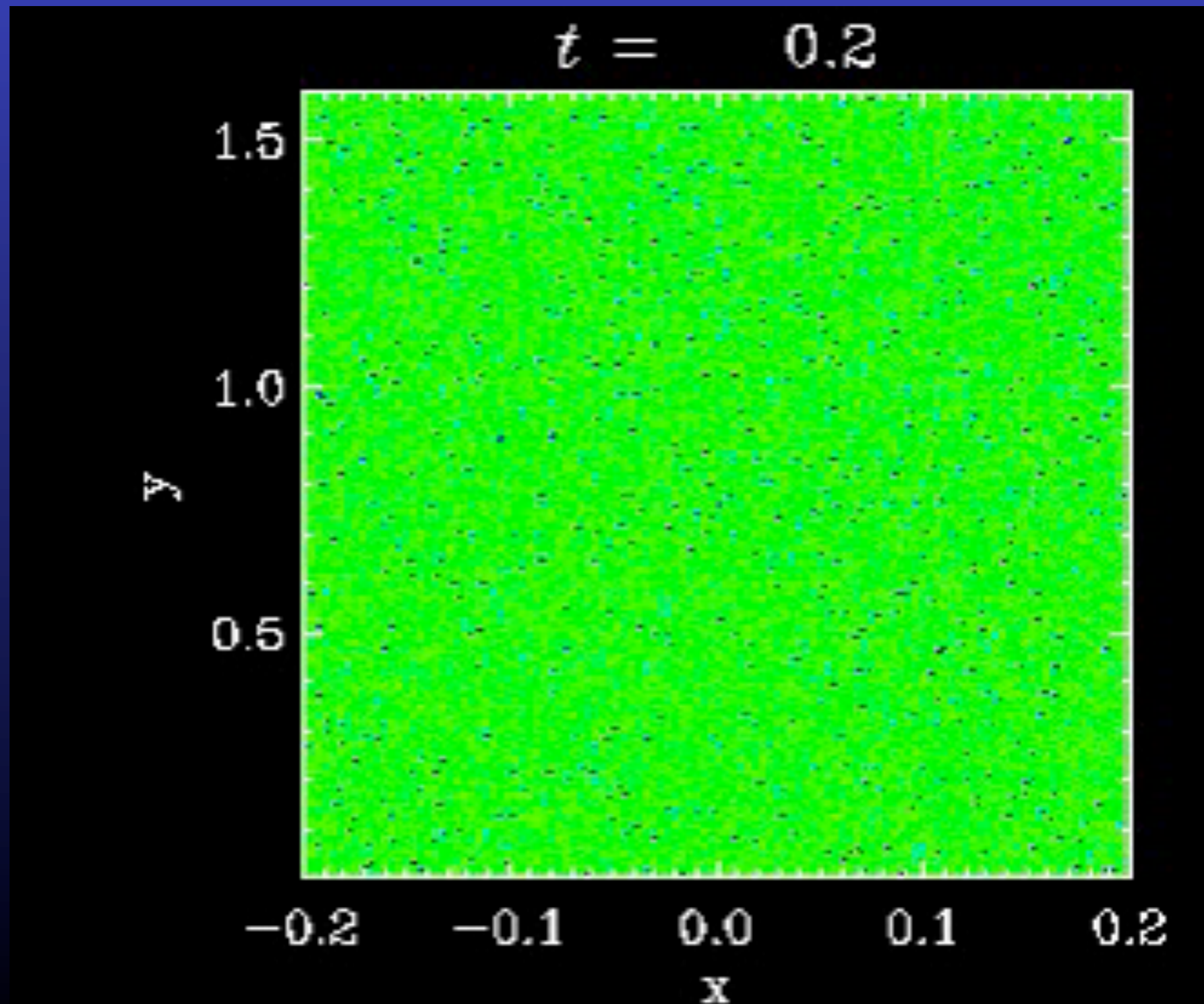


Radial Buoyancy :  $dS/dr < 0$  &  $dP/dr < 0$   
 $\rightarrow$  thermal wind  $dV_{\phi}/dz < 0$       plus: thermal relaxation

Small particles in pressure maxima e.g. a vortex  
e.g. Vortex is in balance between Coriolis forces and  
pressure = same for Zonal flow. Barge & Sommeria 1995

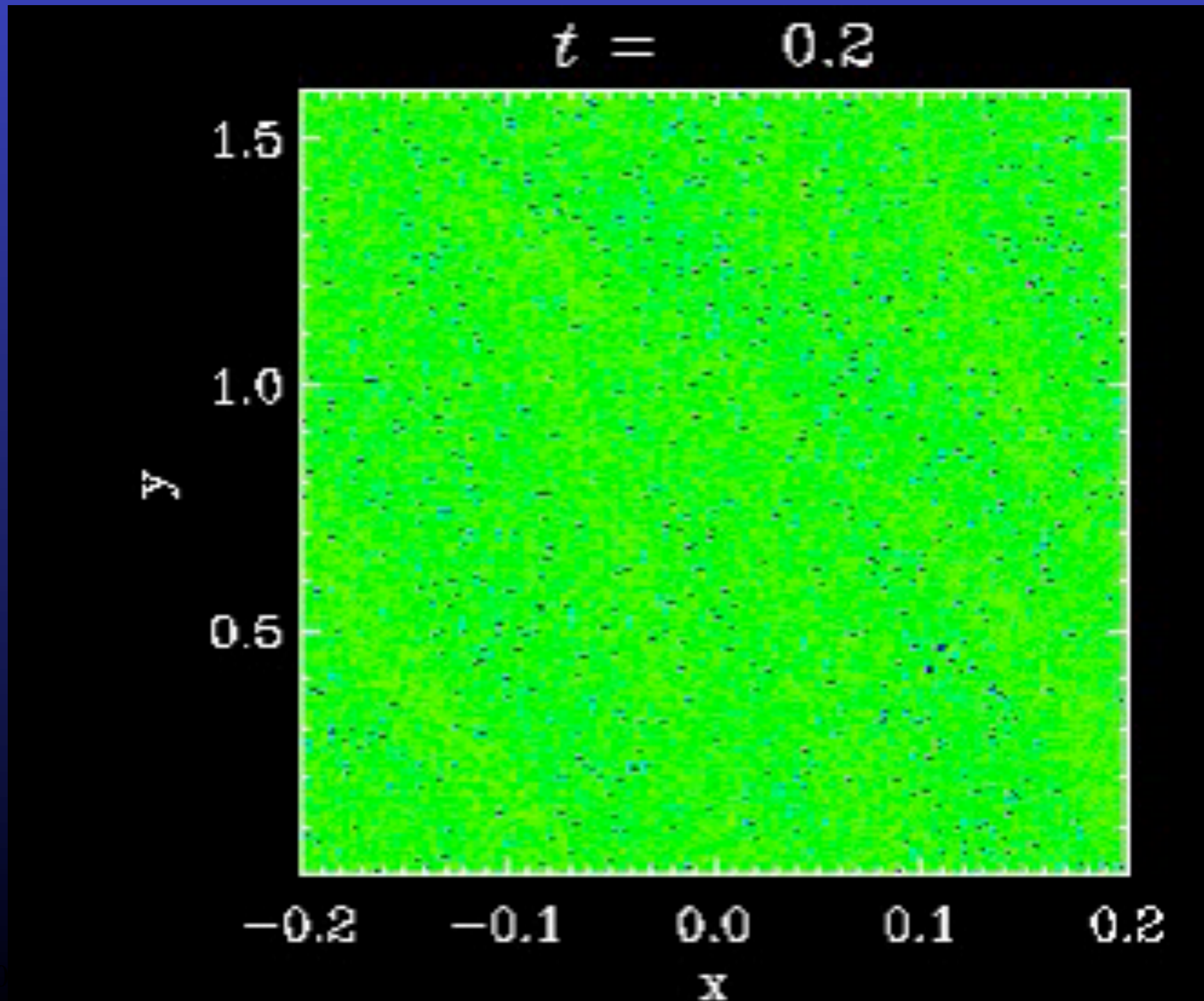


2D:  $St = 1$  particles (20cm)  
(white =  $\times 1000$ ) Natalie Raettig



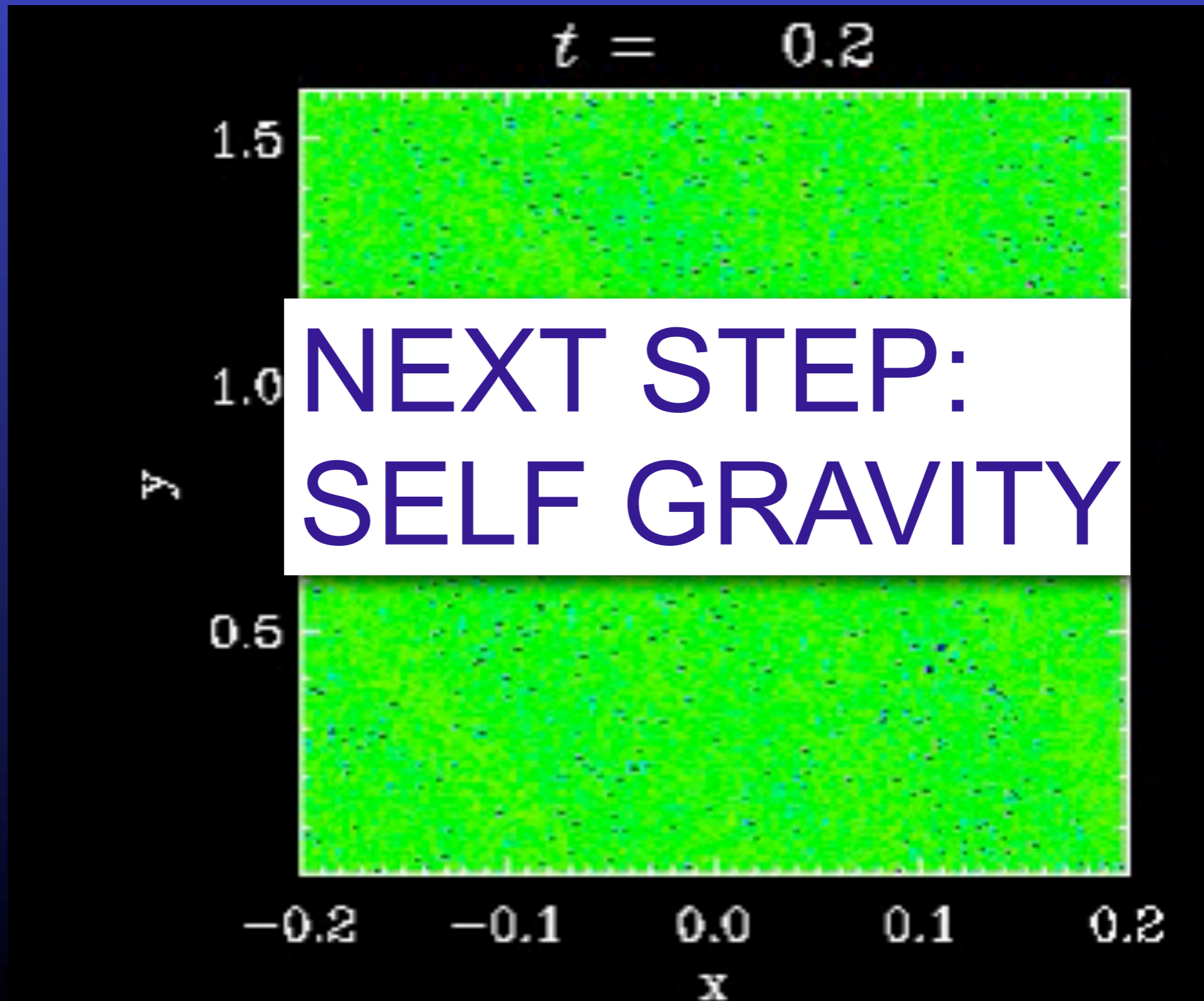
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$St = 0.05$  particles (few millimeter)  
(white =  $\times 1000$ ) Natalie Raettig



12/13/200

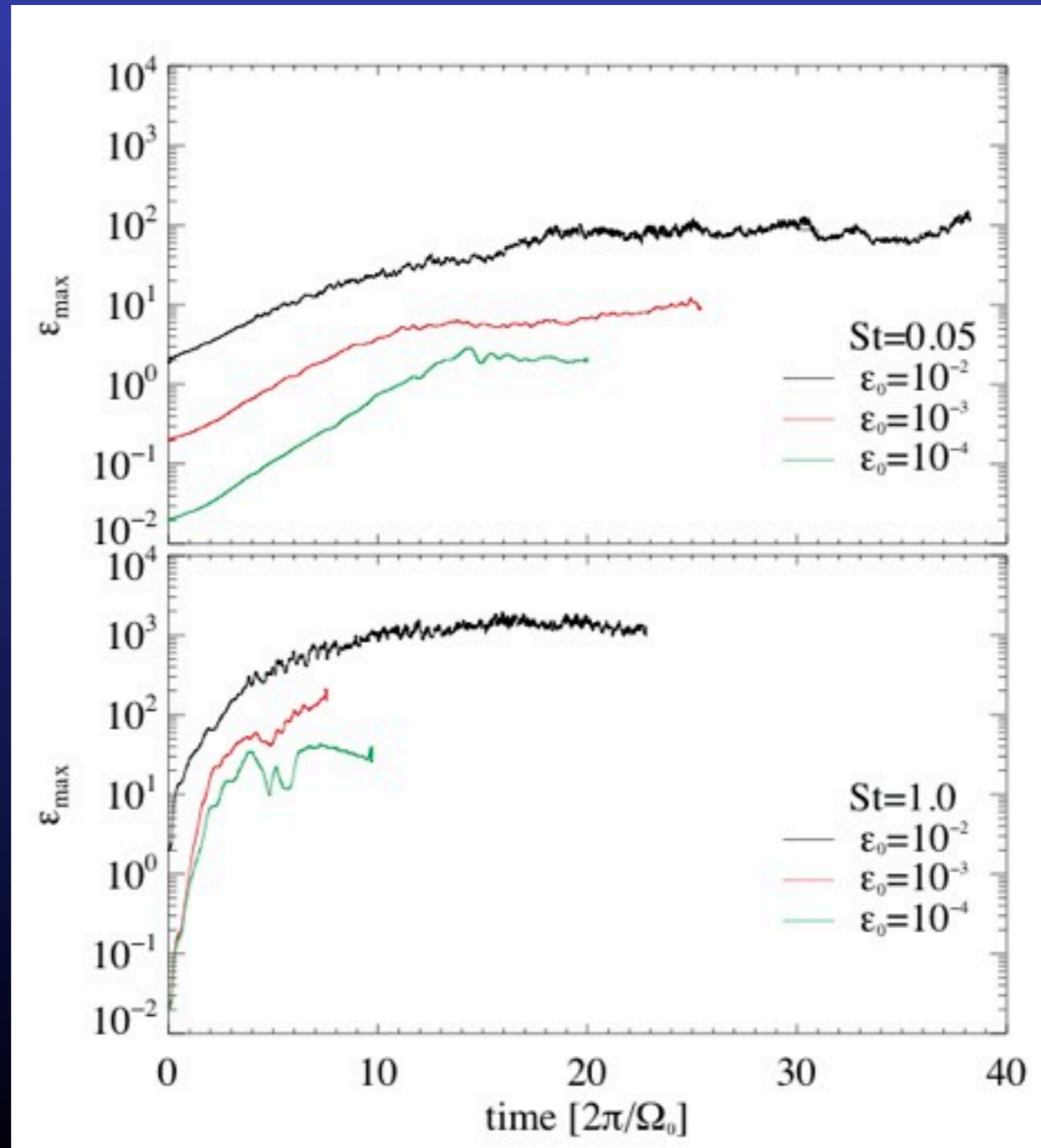
St = 0.05 particles (few millimeter)  
(white = x 1000) Natalie Raettig



12/13/2006



# Dust Concentration in 3D Vortices: Raettig and Klahr, in prep



# Conclusions:

- Disk turbulence can be magnetic in nature, but in resistive regions the entropy structure of the disk creates a thermal wind and eventually vortices.
- Irradiated disks are widely baroclinic and thus Vortices are continuously created.
- Vortices can concentrate  $St = 0.05$  dust at initial abundance of  $\epsilon_{ps} = 1E-4$  to the streaming instability and planetesimal formation.

